2022 FPS 11 Aug.3-5, 2022, Xiangtan University



Heating and Cooling of Accreting Neutron Stars as a Probe of Neutron Star Interiors

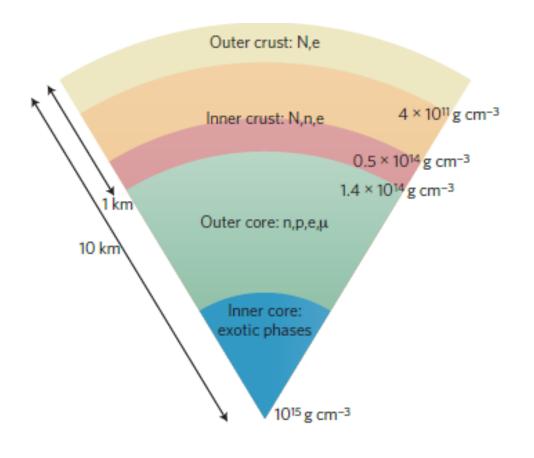
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School of Physical Science and Technology, Xinjiang University 2022.8.5

Today's Contents

- Background
 - Transiently accreting neutron stars
 - Heating and cooling of accreting neutron stars
- Elucidation of physics inside neutron stars from their cooling observations
 - EoS dependence of thermal evolution of accreting NS
 - Effect of neutrino heating on thermal evolution of accreting NS
- The mass and radius dependent of X-ray burst
- Conclusions & Future Perspective

Basic neutron star structure



Newton, 2013, Nat Phys

Typical value $M_{\rm NS} \sim 1.4 M_{\odot}$, $R_{\rm NS} \sim 10$ km

Core density >nuclear density

Possible exotic particles???

- No terrestrial experiments seem possible at such high densities and low temperature ($\ll 10^{10}$ K)
- Many equation of states (EoSs) for NS core matter have been developed

How to constrain EoS?

Many ways, but we consider <u>heating and</u> <u>cooling of transiently accreting neutron stars</u> <u>and X-ray burst as the tools.</u>

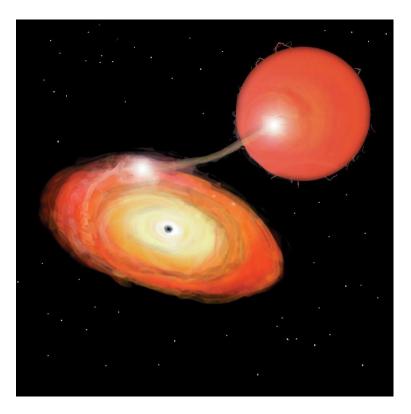
Transiently accreting neutron star

Low-mass x-ray binaries (LMXBs)

- $M_{\rm donar} \lesssim 1 M_{\odot}$
- Orbital period: minites~days
- Old system (t \gtrsim 1Gyr)
- Weak magnetic field $(B\sim 10^7 \text{to } 10^9 \text{ G})$
- Roche-Lobe Overflow

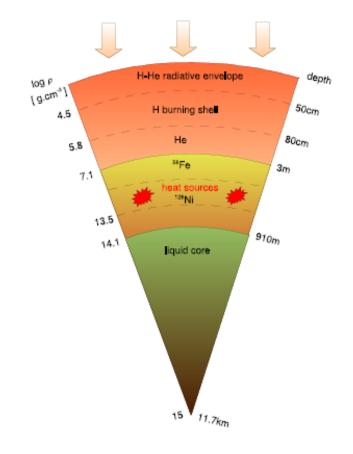
Soft x-ray transients (SXRTs)

- Active phase
 - accretion rate ~10⁻¹⁰-10⁻⁸ M_{\odot} yr⁻¹
 - $Lx \sim 10^{36-39} erg/s$
 - Weeks to months to years
- Quiescent phase (requires sensitive X-ray telescope)
 - little accretion
 - Lx~10³⁴ erg/s



(NASA website)

Accreting NSs as nuclear physics laboratories

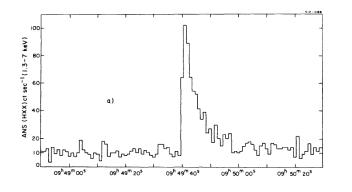


N Chamel, P Haensel . Living Rev. Relativity 2008

 compression/heating of the envelope

thermonuclear burning of light elements

Type I X-ray bursts and superbursts while accreting



compression/heating of the crust

e-captures, n-emission, pycnonuclear reactions

Cooling in quiescence

"Deep crustal heating": 1~2 MeV/nuc

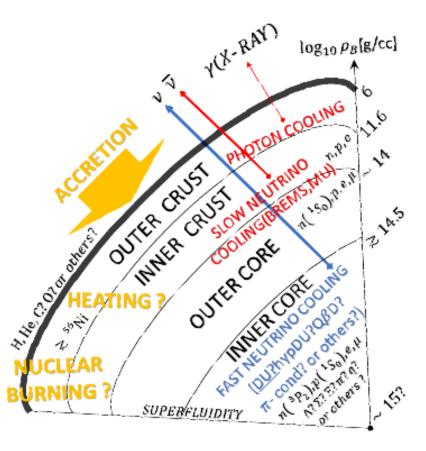
$$Q_i = 6.03 \dot{M}_{-10} \frac{q_i}{\text{MeV}} 10^{33} \text{ erg s}^{-1}$$

P. Haensel et al, 1990,2003,2008

Grindlay, 1976, ApJ

Energy Sources in Accreting NS

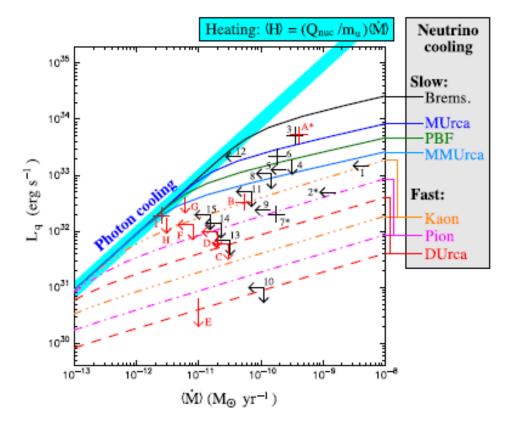
- Outer physics of NS
 - Release of gravitation
 - Nuclear burning 3α reaction→HCNO cycle (Wallance & Woosley 1981) →αp process→rp process (Woosley & Weaver 1981,84) →SnSbTe cycle (Schatz+2001)
- Inner physics of NS
 - Crust heating (Haensel & Zdnick 90,03,08)
 - Shallow heating (Deibel, 2016)
 - v emission (slow+Fast cooling)
 - EOS properties



Quiescent luminosity and burst phenomenon is related to not only outside but also inside physics of NS

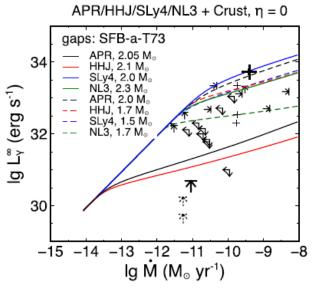
Luminosity in quiescence of soft X-ray transients

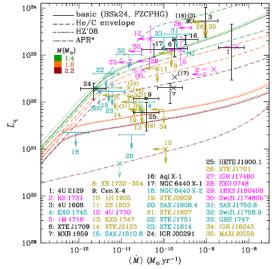
Steady state



Depend On:

- the rate of neutrino cooling
- superfluidity
- Equation of state
- Surface composition





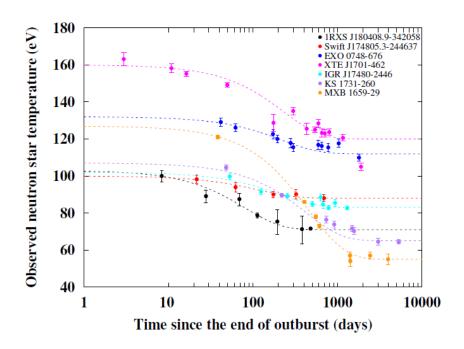
Wijnands et al, 2017

Han & steiner, 2017

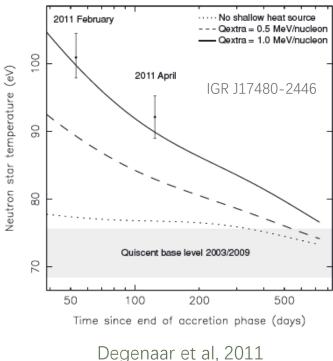
Potekhin, et al, 2020

Luminosity in quiescence of soft X-ray transients

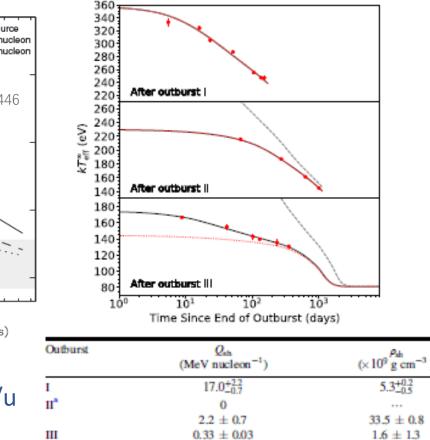
Crust cooling



Wijnands et al, 2017



- Magnitude: 0-17 MeV/u
- Depth: 10⁸-10¹⁰ g cm⁻²
- Physics origin: unknown

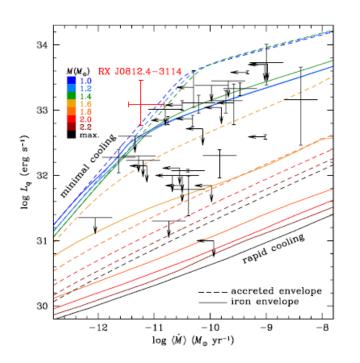


additional shallow heat source?

Parikh, et al, 2017

Luminosity in quiescence of soft X-ray transients

RX J0812.4-3114



• Low accretion rate: $5.8 \times 10^{-12} M_{\odot} \text{ yr}^{-1}$

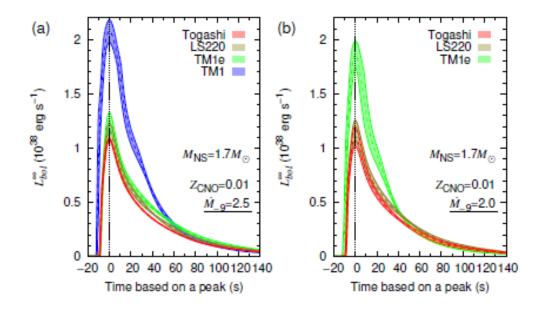
• High luminosity: $L_X \approx 1.6 \times 10^{33} \text{ erg s}^{-1}$

Possible explanation:

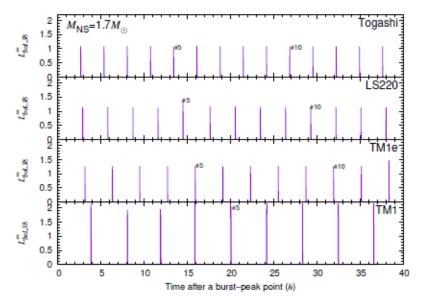
- RX Jo812.4-3114 contains a relatively low-mass NS with minimum cooling
- The system may be young enough that the NS has not fully cooled from the supernova explosion
- Additional heating?

Zhao, et al, 2019

Effects of equation of state on type-I x-ray burst

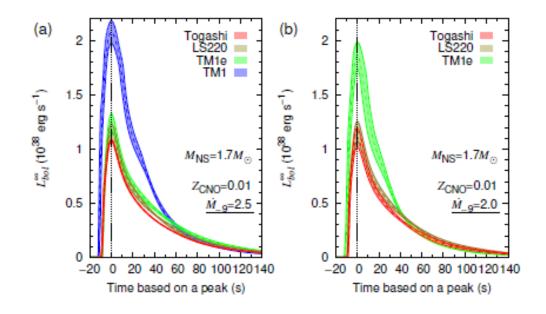




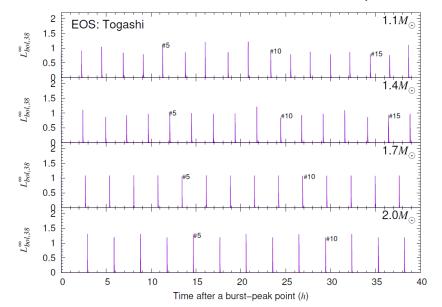


 Previous studies about multi-zone modeling cover with accreted layer, giving boundary conditions on NS crust, but don't consider some physics of interior NS • Dohi, et al. studied X-ray bursts using a general relativistic stellarevolution code with several EoSs

Effects of equation of state on type-I x-ray burst



Dohi,2021,ApJ



 Previous studies about multi-zone modeling cover with accreted layer, giving boundary conditions on NS crust, but don't consider some physics of interior NS • M increases, the time interval becomes larger, the peak luminosity becomes higher

Basic Equations

The general relativistic evolutionary equations:

$$\begin{split} \frac{\partial M_{tr}}{\partial r} &= 4\pi r^2 \rho \ ,\\ \frac{\partial P}{\partial r} &= -\frac{GM_{tr}\rho}{r^2} \left(1 + \frac{P}{\rho c^2}\right) \left(1 + \frac{4\pi r^3 P}{M_{tr}c^2}\right) V^2 \ ,\\ \frac{\partial (L_r e^{2\phi/c^2})}{\partial M_r} &= e^{2\phi/c^2} \left(\varepsilon_{\rm n} - \varepsilon_{\nu} + \varepsilon_{\rm g}\right) \ ,\\ \frac{\partial \ln T}{\partial \ln P} &= \nabla_{\rm rad} \ ,\\ \frac{\partial \phi}{\partial M_{tr}} &= \frac{G(M_{tr} + 4\pi r^3 P/c^2)}{4\pi r^4 \rho} V^2 \ , \end{split}$$

where

$$\frac{\partial M_{tr}}{\partial M_r} = \frac{\rho}{\rho_0} V^{-1}, \quad V \equiv \left(1 - \frac{2GM_{tr}}{c^2 r}\right)^{-1/2}.$$

Fujimoto et al, 1984; Thorne et al, 1977;

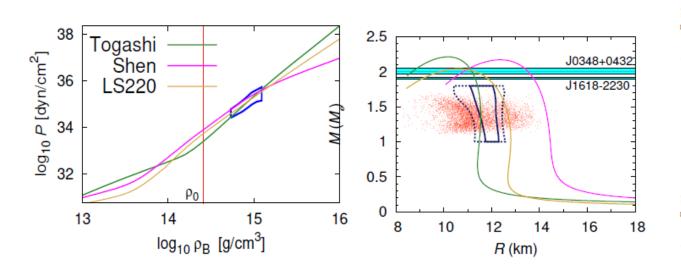
Physics input:

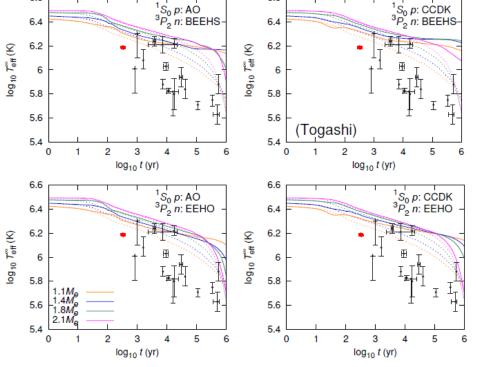
- Equation of state
- Neutrino emission
 - Standard cooling
 - Modified URCA
 - Bremsstrahlung process
 - fast cooling
 - Direct URCA
 - Pion condensation process

• ...

- PBF
- heating
 - deep crustal heating
 - shallow heating ...

EoS dependence of the cooling curves





6.6

6.6

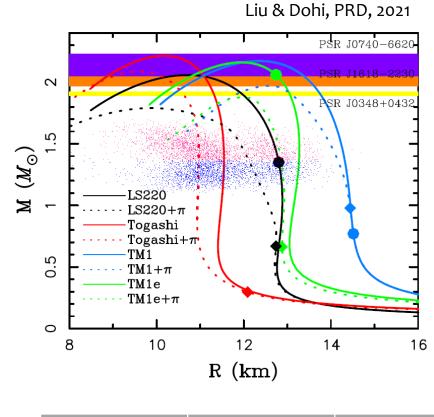
Dohi, Hashimoto, 2019

- ▶ $2M_{\odot}$ NS observation
- ➢ GW170817
- Low mass X-ray binary observations

DU process of Togashi EoS is prohibited

Any other fast cooling processes work in NS? (exotic particles such as hyperons, pions, kaons, and quarks)

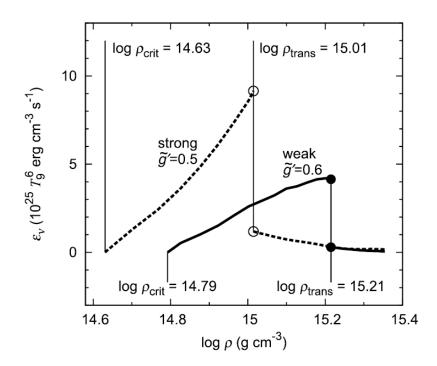
EoS+pion condensation



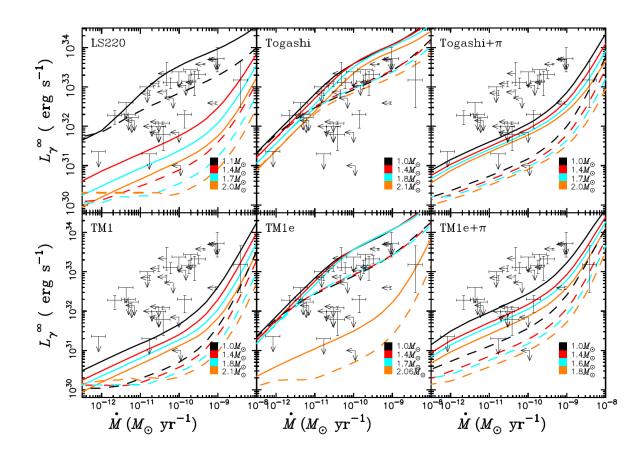
| EoS | Togashi | Togashi+π |
|--------------------------|---------|-----------|
| $M_{\rm max}(M_{\odot})$ | 2.21 | 2.09 |

Effects of pion condensation

- Soften the EoSs
- Strong pion neutrino emission



Quiescent Luminosity of accreting NS

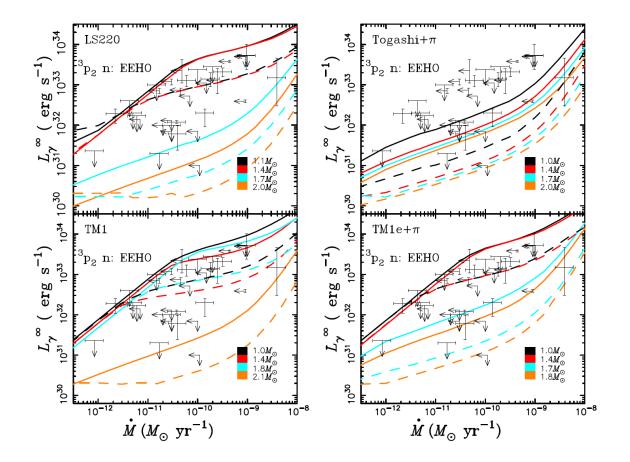


Liu & Dohi, PRD, 2021

Heating curves with TM1,Togashi+ π , TM1e+ π located too low, superfluidity is needed

| EOS | LS220 | Togashi | TM1 | TM1e | $LS220 + \pi$ | Togashi $+\pi$ | $TM1 + \pi$ | TM1e + π |
|--------------------|-------|---------|------|------|---------------|----------------|-------------|--------------|
| M_{DU}/M_{\odot} | 1.35 | | 0.77 | 2.06 | 0.67 | 0.30 | 0.98 | 0.66 |

Quiescent Luminosity of accreting NS

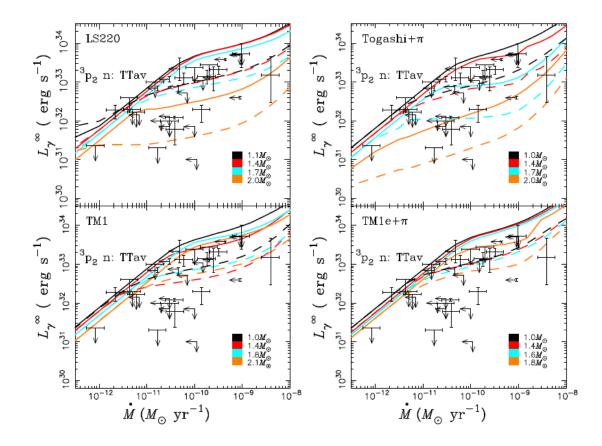


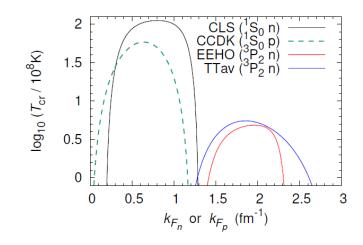
Liu & Dohi, PRD, 2021

Neutron ${}^{1}S_{0}$: CLS Proton ${}^{1}S_{0}$: CCDK

EEHO for Neutron ${}^{3}P_{2}$ is too weak to explain some hot observations

Quiescent Luminosity of accreting NS





Strong neutron ${}^{3}P_{2}$ superfluidity TTav can fit cooling observations

Crustal heating by neutrinos from the surface of accreting neutron stars

Charged Pion Production from infalling matter

T

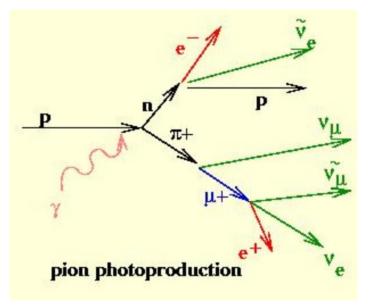
Kinetic energy:

$$= m_0 c^2 \left(\frac{1}{\sqrt{1 - R_{\rm S}/R}} - 1\right)$$

Bildsten, et al., 1992

| Model | R_{14} (km) | T_{14} (MeV) | R_{20} (km) | <i>T</i> ₂₀ (MeV) |
|--------------|---------------|----------------|---------------|------------------------------|
| FSU2 (soft) | 12.89 | 200.5 | 12.03 | 377.2 |
| FSU2 (stiff) | 14.10 | 178.0 | 12.95 | 334.2 |
| HLPS (soft) | 9.95 | 289.5 | 9.68 | 565.2 |
| HLPS (stiff) | 13.59 | 186.8 | 14.14 | 291.6 |

Fattoyev, et al, 2018



If $T_{\text{kin}} \gtrsim 290$ MeV (pion production threshold), it could occur near the surface due to free-fall material

Crustal heating by neutrinos from the surface of accreting neutron stars

• After pion production

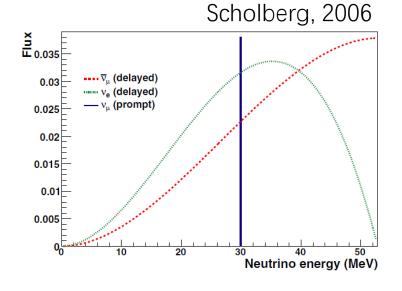
•
$$\pi^0 \rightarrow \gamma + \gamma$$

• $\pi^+ \rightarrow \mu^+ + \nu_{\mu}, \mu^+ \rightarrow e^+ + \nu_e + \bar{\nu}_{\mu}$
• $\pi^- \rightarrow \mu^- + \bar{\nu}, \mu^- \rightarrow e^- + \bar{\nu}, \mu^- + \bar{\nu}_{\mu}$

•
$$\pi^- \rightarrow \mu^- + \bar{\nu}_\mu, \mu^- \rightarrow e^- + \bar{\nu}_e + \nu_\mu$$

$$E_{\nu_{\mu}} = 29.8$$
 MeV, $E_{\nu_{e}} = 33.3$ MeV, $E_{\overline{\nu}_{\mu}} = 37.7$ MeV from neutrino energy spectrum

- Then, assuming the probability 50% that neutrinos move into the crust , the total carrying energy is:
- For neutrino transport, all neutrinos are assumed to eventually deposit in the crust

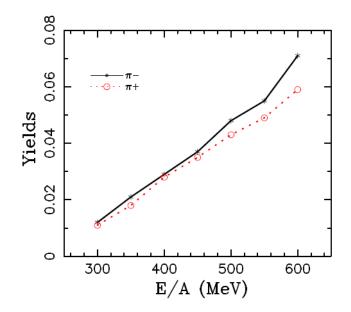


$$\begin{split} q_\nu &\approx 0.5 \left(E_{\nu_\mu} + E_{\nu_e} + E_{\overline{\nu}_\mu} \right) N_{\pi^+} \\ = &50.4 \ \text{MeV} N_{\pi^+} \end{split}$$

Does the pion production affect quiescent luminosity of transiently accreting neutron stars?

Estimation of q_{ν}

- The pion production obtained from the IBUU transport model
- We consider α -Fe collision



| E/A (MeV) | $N_{\pi^{-}}$ | N_{π^+} | $q_{\nu}(\text{MeV})$ |
|-----------|---------------|-------------|-----------------------|
| 300 | 0.012 | 0.011 | 1.16 |
| 350 | 0.021 | 0.018 | 1.97 |
| 400 | 0.029 | 0.028 | 2.87 |
| 450 | 0.037 | 0.035 | 3.63 |
| 500 | 0.048 | 0.043 | 4.59 |
| 550 | 0.055 | 0.049 | 5.24 |
| 600 | 0.071 | 0.059 | 6.55 |

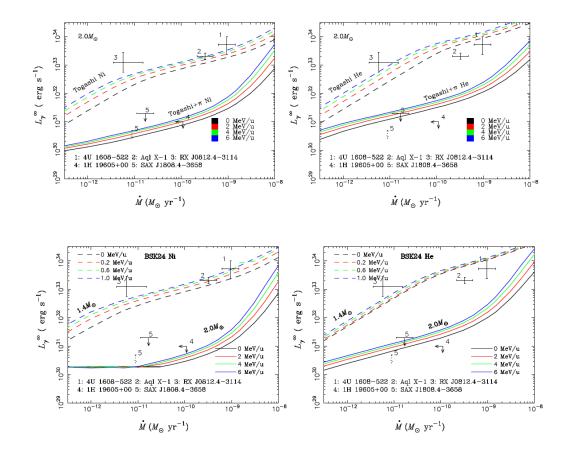
$$\begin{aligned} q_{\nu} &\approx 0.5 \left(E_{\nu_{\mu}} + E_{\nu_{e}} + E_{\overline{\nu}_{\mu}} \right) \left(N_{\pi^{+}} + N_{\pi^{-}} \right) \\ &= 50.4 \text{ MeV} \left(N_{\pi^{+}} + N_{\pi^{-}} \right) \end{aligned}$$

- Decelerated by radiative pressure
- ➢ free-fall condition
- Electromagnetic acceleration

Quiescent luminosities of SXTs with neutrino heating

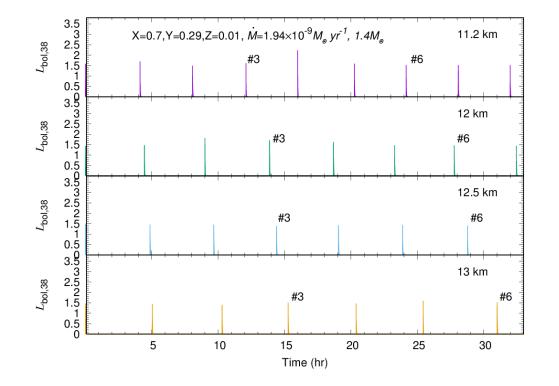
| Model | R_{14} (km) | T_{14} (MeV) | R_{20} (km) | T_{20} (MeV) |
|-----------------|---------------|----------------|---------------|----------------|
| Togashi | 11.55 | 232.1 | 11.17 | 427.2 |
| Togashi + π | 10.97 | 249.7 | 10.33 | 493.8 |
| BSK24 | 12.54 | 207.2 | 12.27 | 363.5 |

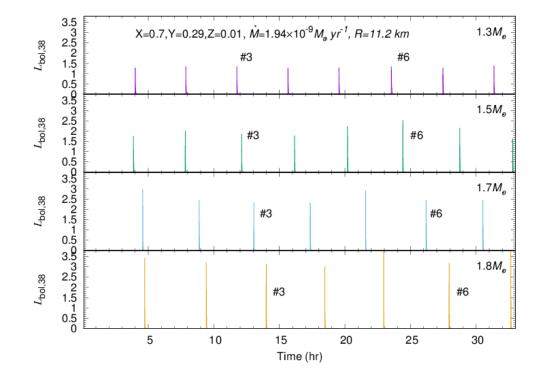
- $q_{\nu} = 0 1$ MeV/u for $1.4 M_{\odot}$
- $q_{\nu} = 0 6$ MeV/u for 2.0 M_{\odot}
- SAX J1808.4-3658 (low accretion rate, high luminosity)can be explained with a $2M_{\odot}$ neutron star
- The observation on cold sources can still be explained with neutrino heating



Liu, et al, 2021, PRD

Mass and radius dependence of Type I x-ray burst

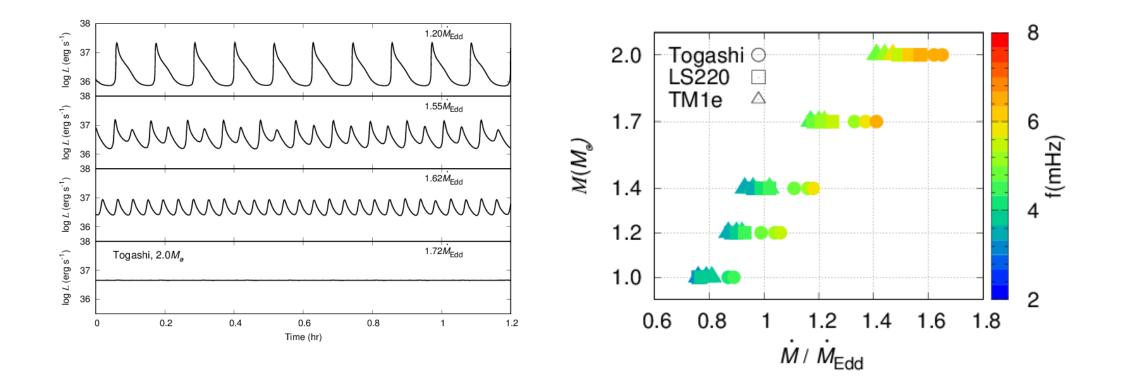




Large radius models get higher luminosity

Higher mass models get higher luminosity

EoS dependence of mHz Qpo



Summary & future work plan

- ✓ Studying the cooling and heating of accreting neutron stars in quiescence
 - EoSs dependence
 - Direct URCA threshold
 - superfluid
 - Effect of neutrino heating
 - The observation on hot source such as SAX J1808.4-3658can be explained with a $2M_{\odot}$ neutron star under the condition that neutrino heating is considered
- ✓ Studying the mass and radius dependence of type I x-ray burst

Future work plan:

- Studying the possible heating mechanism in the crust of accreting neutron stars
- Studying the EoS dependence of crust cooling
- Studying the EoS dependence of type I x-ray burst and mHz QPO