
第十二届全国脉冲星研讨会 (2023)

Hyperon star and hyperon-nucleon interaction research base on astronomical observation

Xiang-Dong Sun (孙向东)

Collaborators: Zhi-Qiang Miao (繆志强)

Bao-Yuan Sun (孙保元)

Ang Li (李昂)



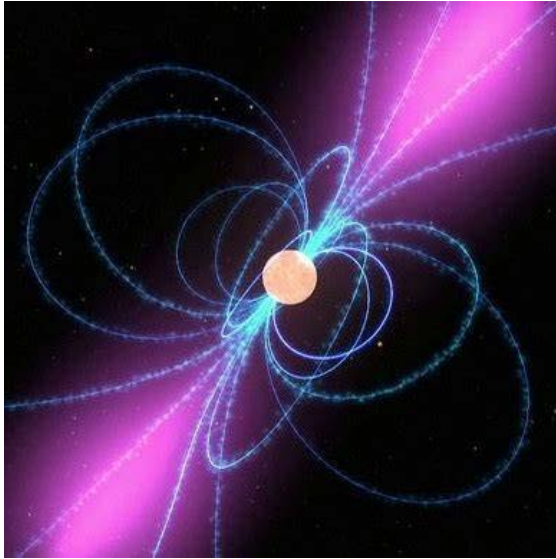
arXiv:2205.10631 ApJ

Outline

- **Background & motivation**
- **Theoretical framework**
- **Bayesian analysis**
- **Discussion**
- **Summary**



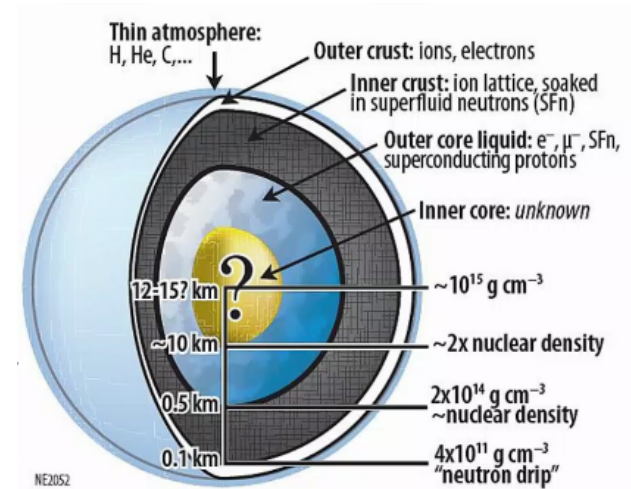
Background



credit: NASA

- 1932 Landau et al. proposed the existence of neutron stars

- 1967 Bell et al. discovered the first neutron star



(image credit: NASA, NICER Team)



- GW170817 opens the era of GW multi-messenger astronomy

● Inner composition?



Background

- Inner core component: Λ , Σ , Ξ , Δ , Ω ...



LHC(EUR)



JLab(US)



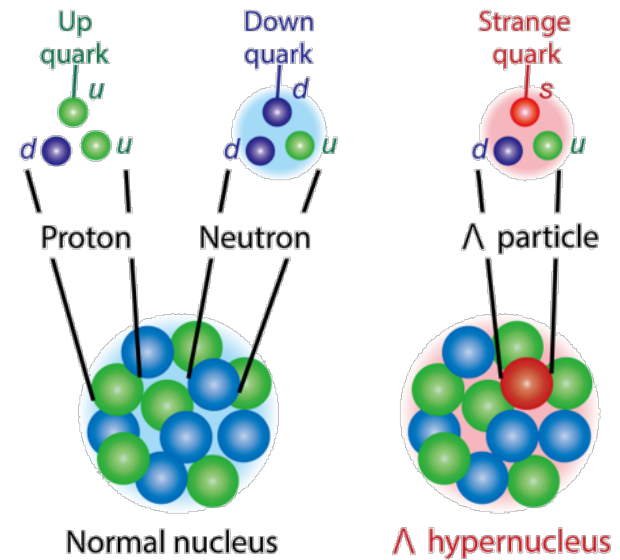
J-PARC(JPN)



FAIR(GER)



HIAF(CHN)



- Laboratory experiments: Extracting baryon-baryon interactions from hyperon-nucleon scattering, structure of the hypernucleus and reactions, (half-life period $\sim 10^{-10}$ s)
- Some available experimental data for single Λ separation energy

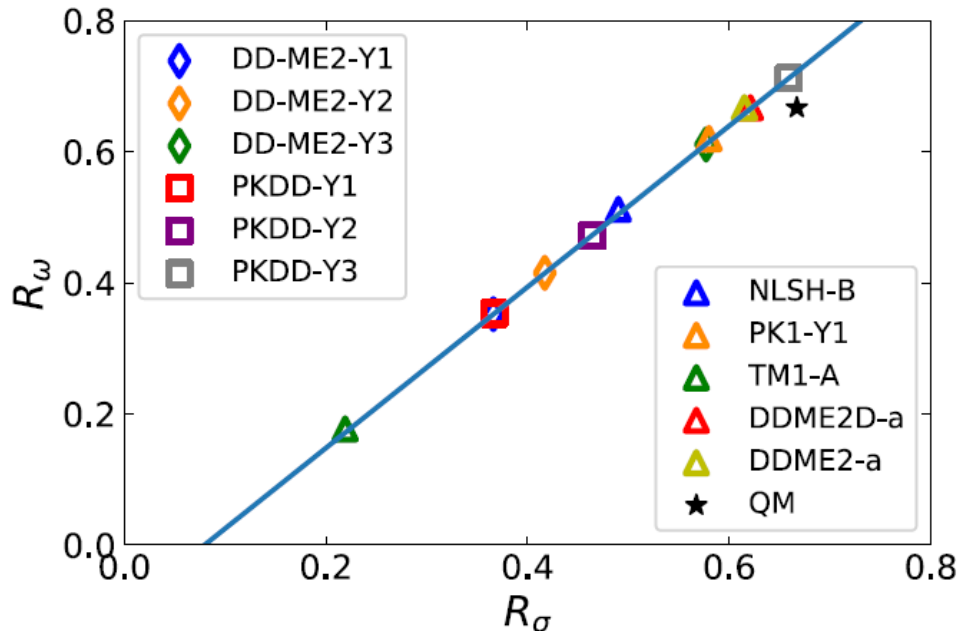
- The inner core of neutron stars might contain hyperons - **Hyperon stars**
- Difficulty : **Hyperon-nucleon interaction**, Model dependent.



Motivation

Meson-hyperon coupling constant: g_{mY}
 Meson-nucleon coupling constant: g_{mN}

$$R_{mY} = g_{mY}/g_{mN} \quad (m = \sigma, \omega, \rho, \delta)$$



Rong et al PHYSICAL REVIEW C 104, 054321 (2021)

Y. T. Rong et al. 2021 PRC
 X. S. Wang et al. 2013 CTP
 C. M. Keil et al. 2000 PRC

➤ **Motivation:** To combine the **NICER** mass, radii data of pulsars as well as **GW170817** data and **the strong correlation** between hyperon nucleon interaction (Rong et al.), to perform Bayesian analysis to Λ hyperon nucleon interaction, and to constrain Λ hyperon nucleon interaction parameter space as well as the properties of hyperon stars.



Theoretical model

► Model: Relativistic Mean Field (Baryon octet)

$$\mathcal{L}_{\text{free}}^B = \sum_B \bar{\psi}_B [i\gamma_\mu \partial^\mu - m_B] \psi_B;$$

$$\mathcal{L}_{\text{int}}^B = \sum_B \bar{\psi}_B [-g_{\sigma B} \sigma - g_{\omega B} \gamma_\mu \omega^\mu - g_{\rho B} \gamma_\mu \vec{\rho}^\mu \cdot \vec{\tau}] \psi_B;$$

$$\begin{aligned} \mathcal{L}_m = & -\frac{1}{2} m_\sigma^2 \sigma^2 + \frac{1}{2} m_\omega^2 \omega_\mu \omega^\mu + \frac{1}{2} m_\rho^2 \vec{\rho}_\mu \cdot \vec{\rho}^\mu \\ & + \frac{1}{2} \partial_\mu \sigma \partial^\mu \sigma - \frac{1}{4} \Omega_{\mu\nu} \Omega^{\mu\nu} - \frac{1}{4} \vec{R}_{\mu\nu} \cdot \vec{R}^{\mu\nu}; \end{aligned}$$

$$\begin{aligned} \mathcal{L}_{\text{NL}} = & -\frac{1}{3} g_2 \sigma^3 - \frac{1}{4} g_3 \sigma^4 + \frac{1}{4} c_3 (\omega_\mu \omega^\mu)^2 \\ & + \Lambda_\nu (g_{\omega B}^2 \omega_\mu \omega^\mu) (g_{\rho B}^2 \rho_\mu \rho^\mu); \end{aligned}$$

$$\begin{aligned} \mathcal{L}_{4f}^B = & -\frac{1}{2} \sum_B \alpha_S^{NB}(\rho) (\bar{\psi}_N \psi_N) (\bar{\psi}_B \psi_B) \\ & - \frac{1}{2} \sum_B \alpha_V^{NB}(\rho) (\bar{\psi}_N \gamma_\nu \psi_N) (\bar{\psi}_B \gamma_\nu \psi_B) \\ & - \frac{1}{2} \sum_B \alpha_{TS}^{NB}(\rho) (\bar{\psi}_N \vec{\tau} \psi_N) (\bar{\psi}_B \vec{\tau} \psi_B) \\ & - \frac{1}{2} \sum_B \alpha_{TV}^{NB}(\rho) (\bar{\psi}_N \vec{\tau} \gamma_\nu \psi_N) (\bar{\psi}_B \vec{\tau} \gamma_\nu \psi_B); \end{aligned}$$

$$\begin{aligned} \mathcal{L}_{\text{hot}}^B = & -\frac{1}{3} \beta_S^{NN} (\bar{\psi}_N \psi_N)^3 - \frac{1}{4} \gamma_S^{NN} (\bar{\psi}_N \psi_N)^4 \\ & - \frac{1}{4} \gamma_V^{NN} [(\bar{\psi}_N \gamma_\mu \psi_N)] [(\bar{\psi}_N \gamma^\mu \psi_N)]^2, \end{aligned}$$

► Nucleon-nucleon interaction:

17 effective RMF interactions:

Walecka, DDRMF, NLRMF, DD-PC, NL-PC

► hyperon-nucleon interaction:

$R_{\sigma\Lambda}$, $R_{\omega\Lambda}$: Relax SU(3) symmetry, [0 - 1]

$R_{\sigma\Sigma}$, $R_{\omega\Sigma}$, $R_{\rho\Sigma}$, $R_{\sigma\Xi}$, $R_{\omega\Xi}$, $R_{\rho\Xi}$:

- ✓ **Vector coupling**: SU(3) symmetry
- ✓ **Scalar coupling**: Potential depth

$$U_\Sigma = +34 \text{ MeV}$$

$$U_\Xi = -14 \text{ MeV}$$

$$R_{\sigma\Sigma} = 0.443 \quad \text{Fortin et al. 2017 PRC;}$$

$$R_{\sigma\Xi} = 0.302 \quad \text{Colucci \& Sedrakian 2013 PLB}$$



Bayesian Formula

- software: Python Bilby, PyMultiNest, Toast, Corner

$$P(\theta|D) = \frac{P(D|\theta)P(\theta)}{\int P(D|\theta)P(\theta)d\theta}$$

$P(\theta)$ is the prior probability of the parameter set θ

$P(D|\theta)$ is total likelihood function

$$\begin{array}{l}
 P(\theta) \\
 \left. \begin{array}{l}
 R_{\sigma\Lambda} \sim U[0, 1] \quad R_{\omega\Lambda} \sim U[0, 1]. \\
 \varepsilon_c \sim U[0.6 \times 10^{15}, 3 \times 10^{15}] \text{ g/cm}^3 \\
 \varepsilon_c \sim U[0.3 \times 10^{15}, 1 \times 10^{15}] \text{ g/cm}^3
 \end{array} \right\} \begin{array}{l}
 \text{NUCL} \\
 \text{NICER} \quad \begin{array}{l}
 \text{PSR J0740+6620} \\
 \text{PSR J0030+0451}
 \end{array}
 \end{array} \\
 \left. \begin{array}{l}
 M \sim U[1.18, 1.21] M_{\odot} \\
 q \sim U[0.5, 1]
 \end{array} \right\} \text{GW170817}
 \end{array}$$



Bayesian Formula

➤ **NICER:** PSR J0740+6620 and PSR J0030+0451

$$P_{\text{NICER}}(d_{\text{NICER}}|\theta) = \prod_j P_j(M(\theta), R(\theta))$$

Abbott et al. 2017 PRL
Abbott et al. 2019 PRX

we equate the individual likelihood P_j to the joint posterior density distribution of M and R

➤ **GW170817:**

$$P_{\text{GW}}(d_{\text{GW}}|\theta) = F(\Lambda_1(\theta; M_1), \Lambda_2(\theta; M_2), \mathcal{M}, q).$$

Abbott et al. 2017 PRL
Abbott et al. 2019 PRX
Hernandez et al. 2019 MN

$F(\cdot)$ is the interpolation function

M_1 and M_2 are related to the chirp mass \mathcal{M} and the mass ratio $q = M_2/M_1$

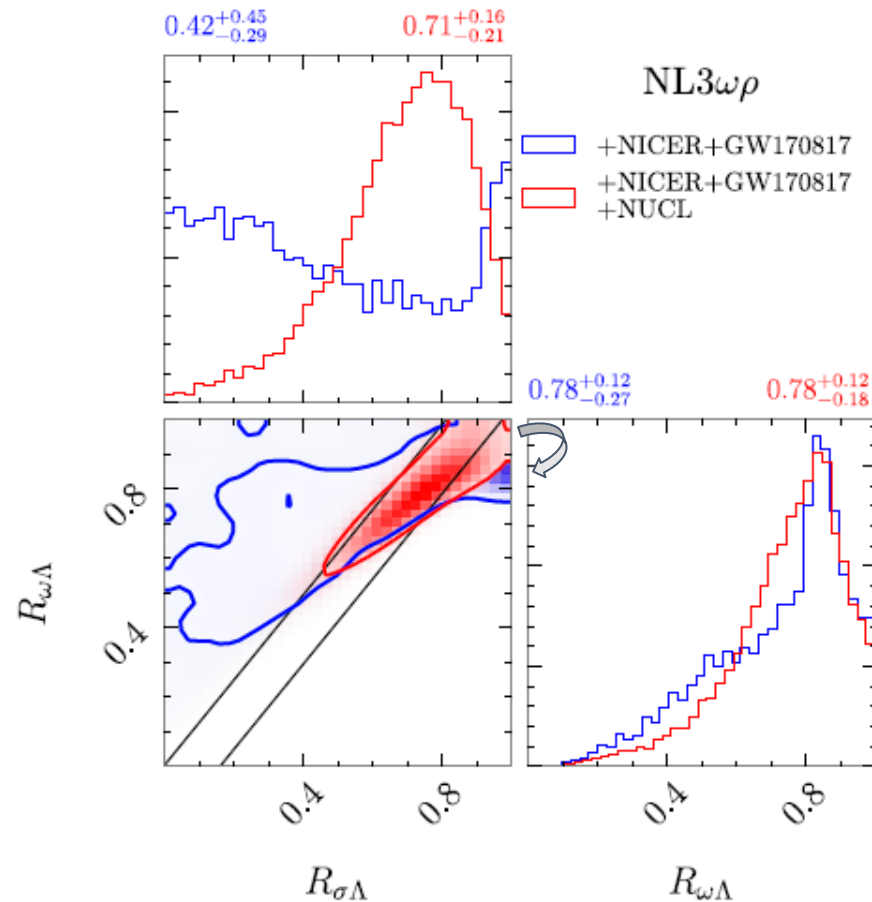
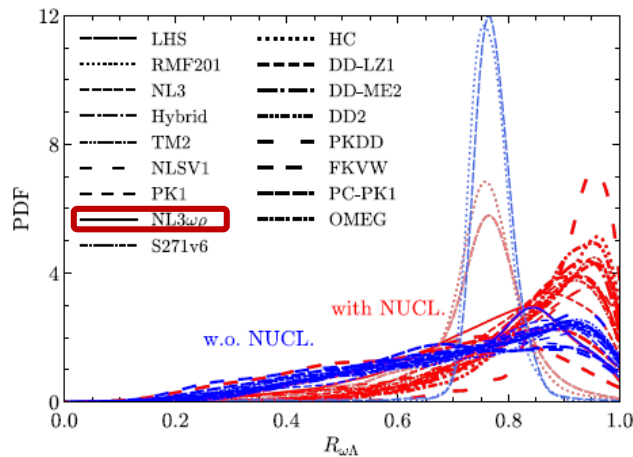
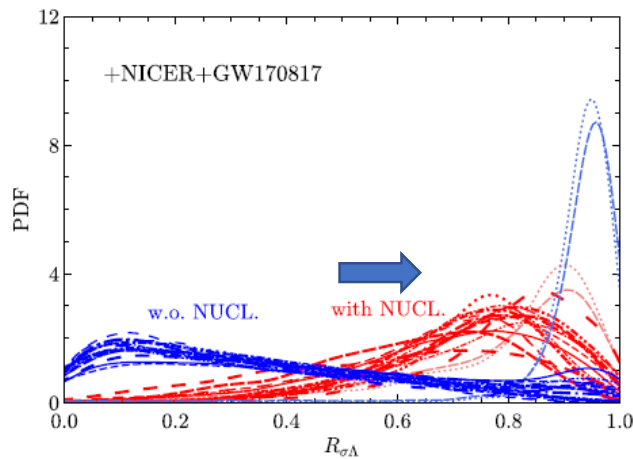
➤ **NUCL:** Single Λ separation energy constraint

$$P_{\text{NUCL}}(d_{\text{NUCL}}|\theta) = \exp \left[-\frac{1}{2} \frac{(R_{\sigma\Lambda} - \bar{R}_{\sigma\Lambda})^2}{\sigma_{R_{\sigma\Lambda}}^2} \right]$$

Rong et al. 2021 PRC



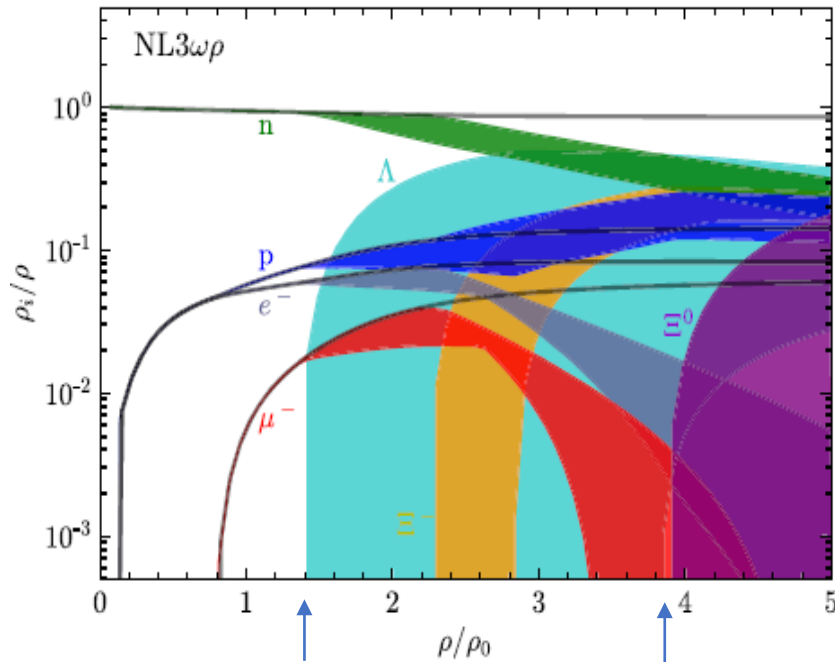
Λ hyperon-nucleon interaction



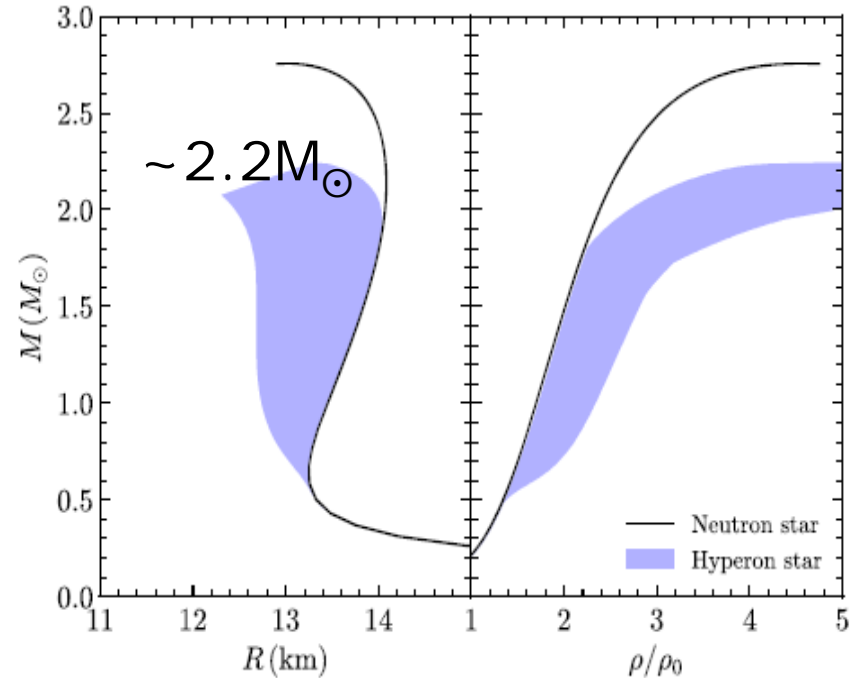
- ❑ Hypernuclei constraint **favors large values of $R_{\sigma\Lambda}$ and $R_{\omega\Lambda}$** and disfavors small values of both couplings
- ❑ The addition of astrophysical observational data on top of the laboratory $R_{\sigma\Lambda}$ - $R_{\omega\Lambda}$ correlation **rotates the linear correlation slightly** towards the direction of small values of $R_{\omega\Lambda}$



Properties of hyperon stars



**Threshold density of Λ hyperons:
 $\sim 1.4\text{--}3.8\rho_0$**

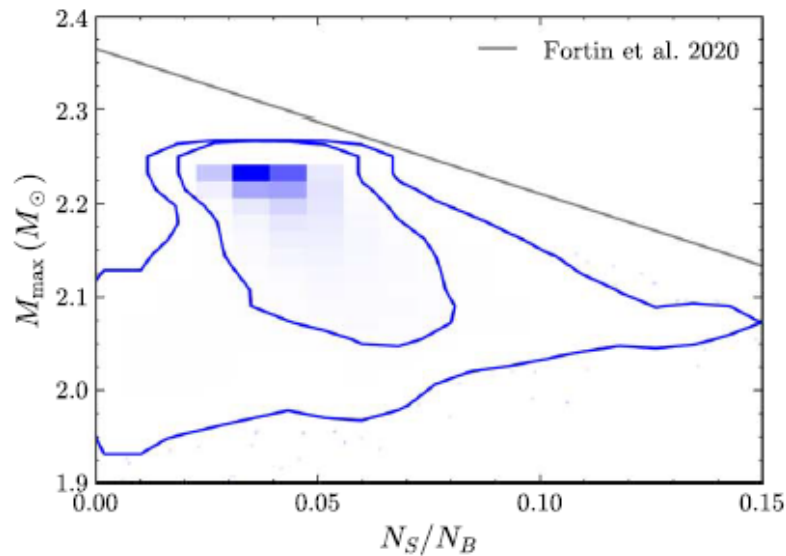
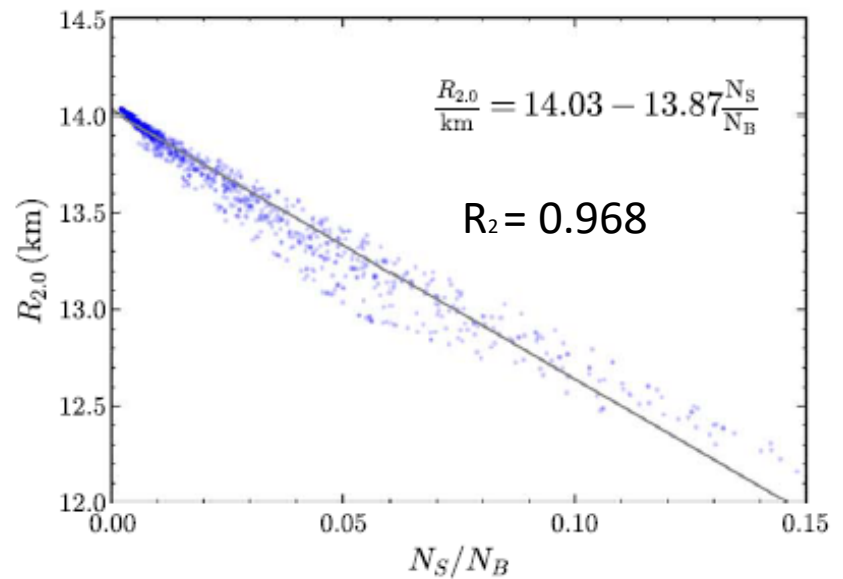
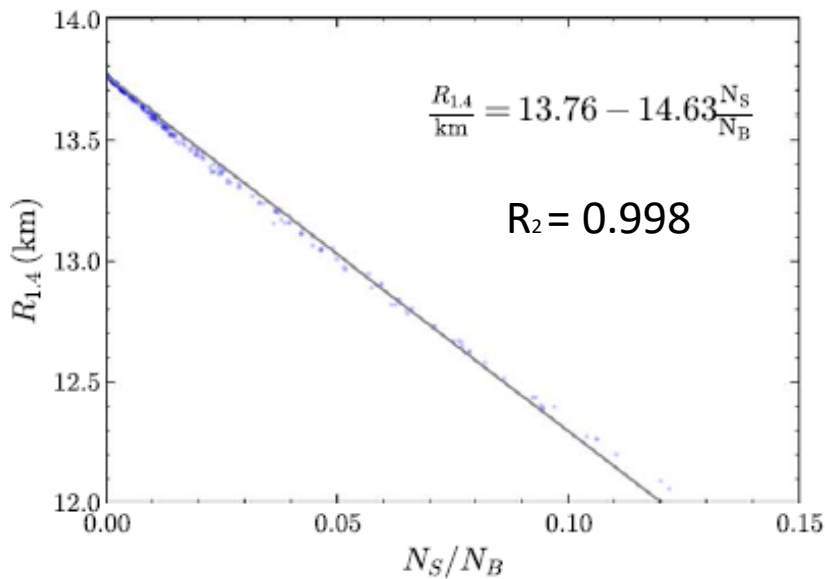


Due to hyperons, the maximum mass is lowered by $\sim 20\%$: $M_{\max} = 2.176^{+0.085}_{-0.202} M_{\odot}$ (68% credible interval).

From the referee “The present article **addresses a long-standing issue** in neutron star physics, namely the hyperon puzzle. The authors **incorporate new information from hypernuclei** calculations and treat the hyperon couplings in a more general way than what exists in the present literature. This is an **interesting work that can have important future implications.**”



Strangeness abundance



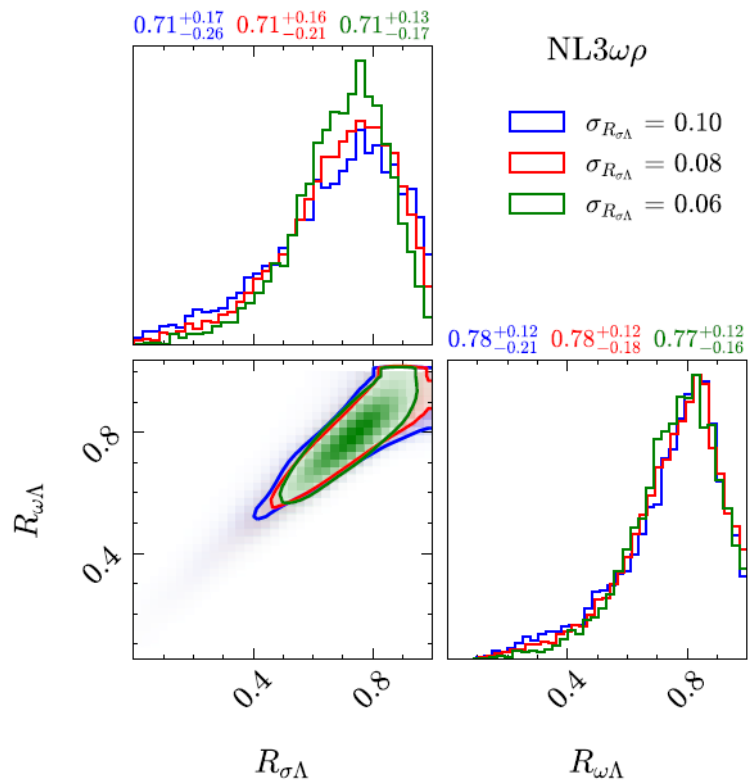
- Excellent **linear anticorrelations** between N_S/N_B and $R_{1.4}$ as well as $R_{2.0}$, with the determination coefficients of $R_2 = 0.998$ and $R_2 = 0.968$.
- There appears also an anticorrelation between N_S/N_B and M_{max} of hyperon stars.



Summary

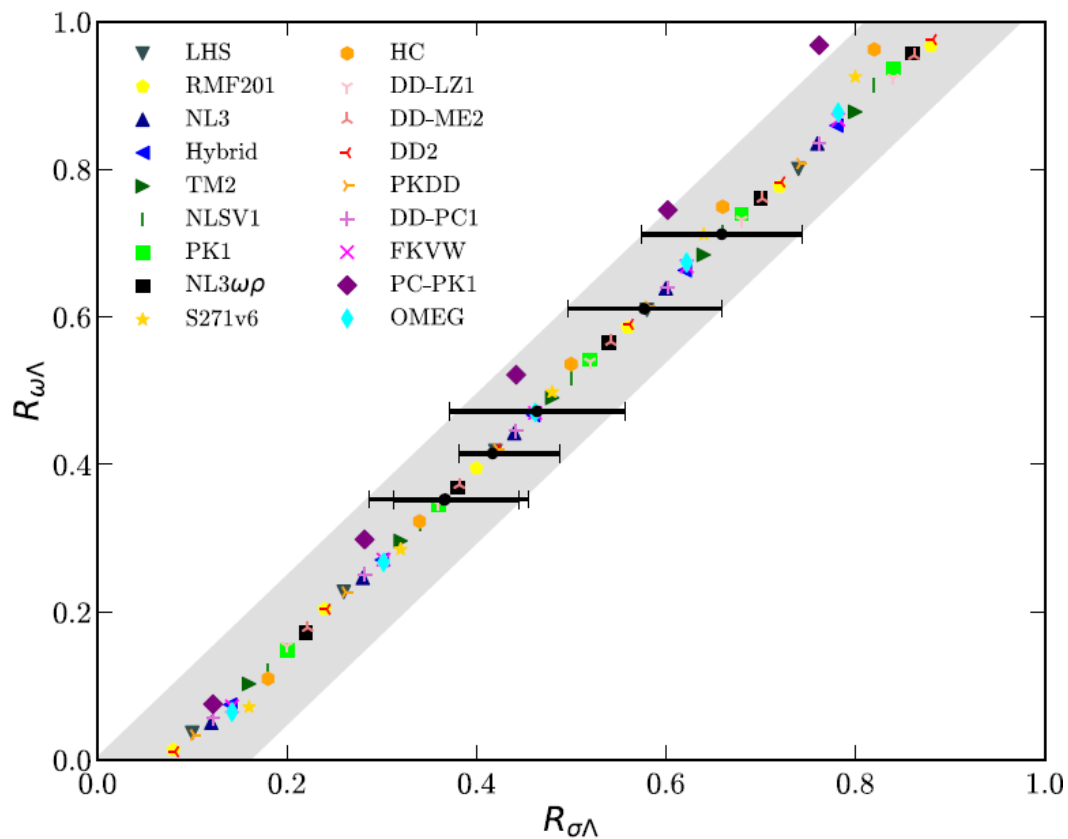
- ❑ We combine the **NICER & GW170817** data and the **strong correlations** of Λ hyperon-nucleon interactions to perform the Bayesian analysis with 17RMF models.
- ❑ We focus Λ hyperon, we **relax the SU(3) symmetry**, take accurately measured mass, radius (for PSR J0030+0451 and PSR J0740+6620), and tidal deformability (for GW170817) and laboratory constraint from the single Λ hypernuclei data.
- ❑ The laboratory constraint from the single Λ hypernuclei data determines the phenomenological interactions, **preventing a small R scalar coupling** through a strong positive correlation between the scalar and vector ones.
- ❑ We find the Λ hyperon **threshold can be as low as ~ 1.4 times** the nuclear saturation density, and the hyperon star EOS is moderately stiff ($M_{\max} = 2.2M_{\odot}$).
- ❑ The decrease of the radii of typical $1.4 M_{\odot}$ and $2.0 M_{\odot}$ stars can both **be linearly depicted as functions of the strangeness fraction** in hyperon stars.

Appendix



	$R_{\sigma\Lambda}$	$R_{\omega\Lambda}$	M_{\max}/M_{\odot}	ρ_c/ρ_0	$R_{2.0}/\text{km}$	$R_{1.4}/\text{km}$	$\Lambda_{1.4}$
$\sigma_{R_{\sigma\Lambda}} = 0.10$	$0.707^{+0.174}_{-0.261}$	$0.777^{+0.116}_{-0.209}$	$2.180^{+0.062}_{-0.228}$	$4.865^{+0.055}_{-0.565}$	$13.980^{+0.088}_{-1.533}$	$13.769^{+0.000}_{-1.086}$	$940.165^{+0.000}_{-448.040}$
$\sigma_{R_{\sigma\Lambda}} = 0.08$	$0.712^{+0.157}_{-0.215}$	$0.778^{+0.121}_{-0.183}$	$2.176^{+0.085}_{-0.202}$	$4.846^{+0.046}_{-0.501}$	$13.968^{+0.096}_{-1.512}$	$13.769^{+0.000}_{-1.084}$	$940.165^{+0.000}_{-443.756}$
$\sigma_{R_{\sigma\Lambda}} = 0.06$	$0.710^{+0.135}_{-0.173}$	$0.772^{+0.120}_{-0.159}$	$2.168^{+0.075}_{-0.186}$	$4.966^{+0.032}_{-0.475}$	$13.950^{+0.075}_{-1.423}$	$13.769^{+0.000}_{-1.064}$	$940.165^{+0.000}_{-430.676}$

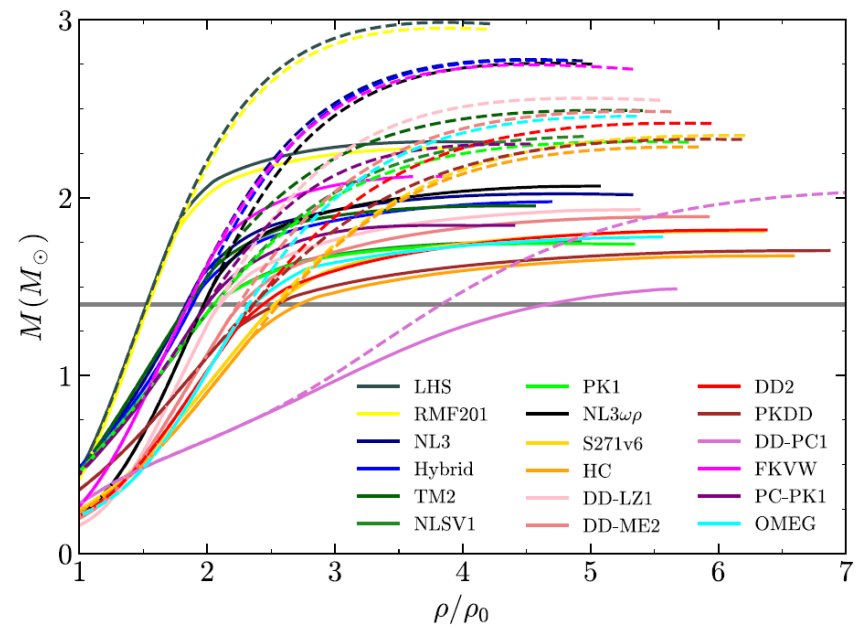
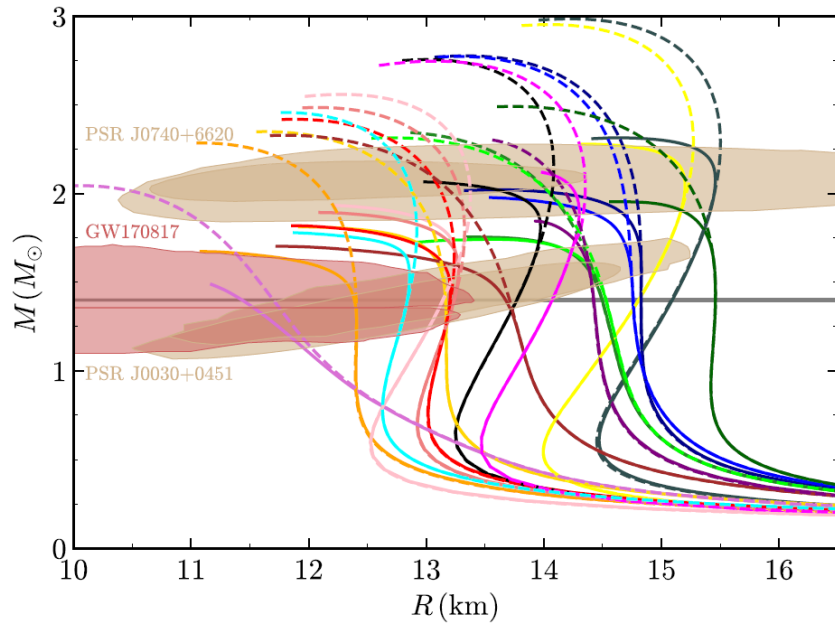
Appendix



$$R_{\omega\Lambda} \approx \frac{-g_{\sigma N} \sigma}{U_N - g_{\sigma N} \sigma} R_{\sigma\Lambda} + \frac{U_\Lambda}{U_N - g_{\sigma N} \sigma}$$

$$R_{\omega\Lambda} \approx \frac{-\alpha_S^{NN} \rho_S}{U_N - \alpha_S^{NN} \rho_S - \beta_S^{NN} \rho_S^2 - \gamma_S^{NN} \rho_S^3 - \gamma_V^{NN} \rho_V^3 - \alpha_S^{N\Lambda} \rho_S^\Lambda - \alpha_V^{N\Lambda} \rho_V^\Lambda} R_{\sigma\Lambda} + \frac{U_\Lambda}{U_N - \alpha_S^{NN} \rho_S - \beta_S^{NN} \rho_S^2 - \gamma_S^{NN} \rho_S^3 - \gamma_V^{NN} \rho_V^3 - \alpha_S^{N\Lambda} \rho_S^\Lambda - \alpha_V^{N\Lambda} \rho_V^\Lambda},$$

Appendix



Appendix

