Bow Shock Emission Model For the White Dwarf Pulsar AR Scorpii

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Collaboration: Jin-Jun Geng, Bing Zhang

Geng, Zhang & Huang, 2016, ApJL, 831, L10 (arxiv:1609.02508)

Outline

- **1. AR Scorpii --- introduction**
- **2. Bow shock emission model**
- **3. Conclusions**

A binary system: white dwarf + M star (116 pc) Orbital period: 3.56 h, nearly circular orbit

Marsh et al. 2016, Nat

A binary system: white dwarf + M star (116 pc) Orbital period: 3.56 h, nearly circular orbit

1.97 minutes period

High speed multi-band measurements

Marsh et al. 2016, Nat

A binary system: white dwarf + M star (116 pc) Orbital period: 3.56 h, nearly circular orbit

1.97 minutes period

Orbital phase (cycle)

Fourier amplitudes vs. frequency

Beat phase (cycles)

A binary system: white dwarf + M star (116 pc) Orbital period: 3.56 h, nearly circular orbit

1.97 minutes period

Beat frequency:
$$
U_B = U_S - U_O
$$

\n $P_B = 1.97 \text{ min}$

\n $P_O = 3.56 \text{ h}$

\n W_{max}

WD spin period: $P_s = 1.95$ min

A binary system: white dwarf + M star (116 pc) Orbital period: 3.56 h, nearly circular orbit

1.97 minutes period WD spin: $P_s = 1.95$ min

Takata, Yang & Cheng, 2017: Emission from relativistic electrons trapped by closed magnetic field lines of the magnetic WD

Fig. 1.— Schematic view of the AR Sco system and the coordinate in the study. The observer is located within the plane made by the spin-axis (z-axis) and x-axis. The spin

Fig. 10.— Spectrum of ARScorpii. The observational data (filled circles and thick d

ig. 6.— Light curve in 0.1-1eV energy bands as a function of the inclination angle (α) ie viewing angle (ζ) . The companion star is located between the WD and observer (in

Takata J., Yang H., Cheng K.S. 2017, ApJ

Our model:**bow shock emission**

Picture taken from the website: http://news.52shehua.com/article_img-6499.html

K. S. Cheng et al.: Bow shock in PSR 1259-63 / LS 2883 system

Kong, Cheng & Huang, 2012, ApJ also see: Kong, Yu, Huang, Cheng, 2011, MNRAS

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The WD as a pulsar

$$
M_{\rm WD} = 0.8 M_{\odot} \R_{\rm WD} = 7 \times 10^8 \text{ cm} \R_{\rm mash\, et\, al.\, 2016}
$$

$$
P = 1.95 \text{ min} \quad \dot{P} = 3.9 \times 10^{-13} \text{s s}^{-1}
$$

$$
\frac{1}{\sqrt{2\pi}}
$$

Spin down luminosity: $\dot{E}_{\rm rot}$ = $-4\pi^2 I \dot{P} P^{-3}$ $\simeq 3.0 \times 10^{33} \rm erg\; s^{-1}$ $\dot{E}_{\rm mag} \simeq \frac{(2\pi)^4 B_p^2 R_{\rm WD}^6}{6c^3 P^4}$ For dipole and wind outflow, magnetic spin down power:

$$
B_p \simeq \left(\frac{3M_{\text{WD}}c^3}{5\pi^2 R_{\text{WD}}^4} P \dot{P}\right)^{1/2} = 7.1 \times 10^8 \text{ G} \left(\frac{P}{1.95 \text{ minutes}} \frac{\dot{P}}{3.9 \times 10^{-13} \text{s s}^{-1}}\right)^{1/2}
$$

 $R_{\rm lc} = cP/2\pi = 5.6\,\times\,10^{11}\,\rm cm$ > WD-MD distance: $d \, \sim \, 7.6\,\times\,10^{10} \rm cm$

The MD is in the magnetosphere of the WD.

The WD as a pulsar (Zhang & Gil 2005)

$$
B_p \simeq \left(\frac{3M_{\rm WD}c^3}{5\pi^2 R_{\rm WD}^4} P \dot{P}\right)^{1/2} = 7.1 \times 10^8 \,\mathrm{G} \left(\frac{P}{1.95 \text{ minutes}} \frac{\dot{P}}{3.9 \times 10^{-13} \mathrm{s} \text{ s}^{-1}}\right)^{1/2}
$$

$$
R_{\rm lc} = cP/2\pi = 5.6 \times 10^{11} \,\rm cm
$$

Polar cap opening angle:

$$
\theta_{\text{open}} = \left(\frac{R_{\text{WD}}}{R_{\text{lc}}}\right)^{1/2} = 2^{\circ} \left(\frac{P}{1.95 \text{ minutes}}\right)^{-1/2}
$$

$$
\left(\begin{array}{c}\n\bullet \\
\bullet \\
\bullet\n\end{array}\right)
$$

$$
M_{\text{WD}} = 0.8M_{\odot}
$$

\n
$$
R_{\text{WD}} = 7 \times 10^8 \text{ cm}
$$

\n
$$
P = 1.95 \text{ min}
$$

\n
$$
\dot{P} = 3.9 \times 10^{-13} \text{s s}^{-1}
$$

Polar cap radius:

$$
R_{\text{pc}} = R_{\text{WD}} \left(\frac{R_{\text{WD}}}{R_{\text{lc}}} \right)^{1/2} = 2.5 \times 10^7 \left(\frac{P}{1.95 \text{ minutes}} \right)^{-1/2} \text{cm}
$$

Maximum potential drop across the cap:

$$
\Phi_{\text{max}} = \frac{2\pi^2 B_p R_{\text{WD}}^3}{c^2 P^2} \simeq 3.9 \times 10^{11} \text{statV}
$$

$$
\times \left(\frac{P}{1.95 \text{ minutes}}\right)^{-3/2} \left(\frac{\dot{P}}{3.9 \times 10^{-13} \text{s s}^{-1}}\right)^{1/2}
$$

Can accelerate electrons to:

$$
\gamma_e = q_e \Phi_{\text{max}}/m_e c^2 \simeq 2.3 \times 10^8
$$

The WD as a pulsar: polarization observations

2,500 2,000 Total counts **Total counts:** 1,500 1,000 500 Total polarized counts 3,000 **Total polarized counts:** 2,000 1,000 Ω 50 Linear polarization (%) 40 30 **Linear polarization:** 20 10 **(up to 40%)** $\overline{0}$ 250 Position angle (°) 200 150 **Polarization angle:** 100 50 \overline{O} Circular polarization (%) 10 5 **Circular polarization:** \circ **(up to n%)** -5 -10

Buckley et al: Polarization behavior is similar to Crab pulsar.

0.610

Buckley et al. 2017, Nat Astron.

0.625

0.620

BJD - 2457462.0

 0.615

The WD as a pulsar

nature astronomy

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Polarimetric evidence of a white dwarf pulsar in the binary system AR Scorpii

D. A. H. Buckley^{1*}, P. J. Meintjes², S. B. Potter¹, T. R. Marsh³ and B. T. Gänsicke³

The variable star AR Scorpii (AR Sco) was recently discovered to pulse in brightness every 1.97 min from ultraviolet wavelengths into the radio regime. The system is composed of a cool, low-mass star in a tight, 3.55-hour orbit with a more massive white dwarf. Here we report new optical observations of AR Sco that show strong linear polarization (up to 40%) that varies strongly and periodically on both the spin period of the white dwarf and the beat period between the spin and orbital period, as well as low-level (up to a few per cent) circular polarization. These observations support the notion that, similar to neutronstar pulsars, the pulsed luminosity of AR Sco is powered by the spin-down of the rapidly rotating white dwarf that is highly magnetized (up to 500 MG). The morphology of the modulated linear polarization is similar to that seen in the Crab pulsar, albeit with a more complex waveform owing to the presence of two periodic signals of similar frequency. Magnetic interactions between the two component stars, coupled with synchrotron radiation from the white dwarf, power the observed polarized and non-polarized emission. AR Sco is therefore the first example of a white dwarf pulsar.

 $\frac{5}{6}$ 1,500 $\frac{1}{2}$ 1,000 anor Ē 3 2.00 30 20 10 $\widehat{\in}$ 200 150 100

Polarization behavior is similar to Crab pulsar. The WD is highly magnetized: ~ 5x10⁸ G.

"**AR Sco is therefore the first example of a white dwarf pulsar.**"

Buckley et al. 2017, Nat Astron.

Two peaks, with similar amplitude: a nearly perpendicular rotator

The peaks are wide (duty cycle ~50%):

$$
\theta_{\rm A} \sim 90^{\circ}/2 = 45^{\circ}
$$

$$
r_{\rm A} = 5.8 \times 10^{10} \text{ cm}
$$

 $\bar{B} = (B_A + B_B)/2 \simeq 1200$ G

Synchrotron emission from the bow shock

$$
\bar{B}=(B_{\rm A}+B_{\rm B})/2\simeq 1200~{\rm G}
$$

Electron distribution:

$$
\frac{dn_e}{d\gamma_e} = \begin{cases} C\gamma_e^{-p}, & \gamma_m \leq \gamma_e \leq \gamma_c, \\ C\gamma_c\gamma_e^{-p-1}, & \gamma_c < \gamma_e \leq \gamma_{\text{max}}, \end{cases}
$$

Characteristic Lorentz factors:

$$
\gamma_m = 45
$$

\n
$$
\gamma_c = 73 (P/1.95 \text{ min})^{-1}
$$

\n
$$
\gamma_{\text{max}} = (6\pi q_e / \sigma_T \bar{B})^{1/2} = 3.4 \times 10^6
$$

\n
$$
F_{\nu, \text{peak}} = \frac{\sqrt{3} q_e^3 \bar{B}}{m_e c^2} \frac{n_e V}{4\pi D_L^2}
$$

\n
$$
D_L = 116 \text{ pc}
$$

Synchrotron emission from the bow shock

$$
\bar{B} = (B_{A} + B_{B})/2 \simeq 1200 \text{ G}
$$
\n
$$
\frac{dn_{e}}{d\gamma_{e}} = \begin{cases}\nC\gamma_{e}^{-p}, & \gamma_{m} \leq \gamma_{e} \leq \gamma_{c}, \\
C\gamma_{c}\gamma_{e}^{-p-1}, & \gamma_{c} < \gamma_{e} \leq \gamma_{max},\n\end{cases}
$$
\n
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\gamma_{m} = 45
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\gamma_{max} = (6\pi q_{e}/\sigma_{T} \bar{B})^{1/2} = 3.4 \times 10^{6}
$$
\n
$$
F_{\nu, \text{peak}} = \frac{\sqrt{3} q_{e}^{3} \bar{B}}{m_{e} c^{2}} \frac{n_{e} V}{4 \pi D_{L}^{2}}
$$
\n
$$
p \sim 2.4
$$
\n
$$
n_{e} = 3.5 \times 10^{8} \text{ cm}^{-3}
$$

Origin of the electrons

$$
n_e = 3.5 \times 10^8 \text{ cm}^{-3}
$$

The Goldreich & Julian charge density:

$$
n_{\rm GJ} = \frac{\Omega \cdot B_{\rm B}}{2\pi q_e c} = 1.1 \text{ cm}^{-3}
$$

The electrons should come from the M-star:

10¹⁰
\n10¹¹
\n10¹¹
\n10¹¹
\n10¹²
\n
$$
\frac{1}{\mu}
$$

\n10¹³
\n $\frac{1}{\mu}$
\n10¹⁴
\n10¹⁵
\n10¹⁶
\n10¹⁷
\n10¹⁸
\n10¹⁹
\n $\alpha \nu^{1/3}$
\n ν_{nl}
\

$$
\dot{M} = 4\pi R_{\rm MD}^2 v n_e m_p / \zeta
$$

$$
= 4.1 \times 10^{-11} (\zeta/0.2)^{-1} M_{\odot} \, \text{yr}^{-1}
$$

Note that typical M-star wind is:

$$
10^{-15} \text{--} 10^{-10} M_{\odot} \text{ yr}^{-1}
$$

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Another similar bow shock explanation: A Precessing White Dwarf Synchronar?

Katz 2017, ApJ

Conclusions

AR Scorpii:

- **(1) Pulsed emission is observed (radio to UV).**
- **(2) Spectrum is non-thermal.**
- **(3) The WD is a pulsar. A bow shock can be generated due to the interaction between the WD pulsar and the M-dwarf companion.**

Thank You! Geng, Zhang & Huang, 2016, ApJL, 831, L10 arxiv:1609.02508