

# Fast/Future Pulsar Symposium 9

Xiamen University, Xiamen, Fujian, China



## Triaxially-deformed Freely-precessing Neutron Stars

Continuous electromagnetic and gravitational radiation

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arXiv:2007.02528 (accepted by MNRAS)

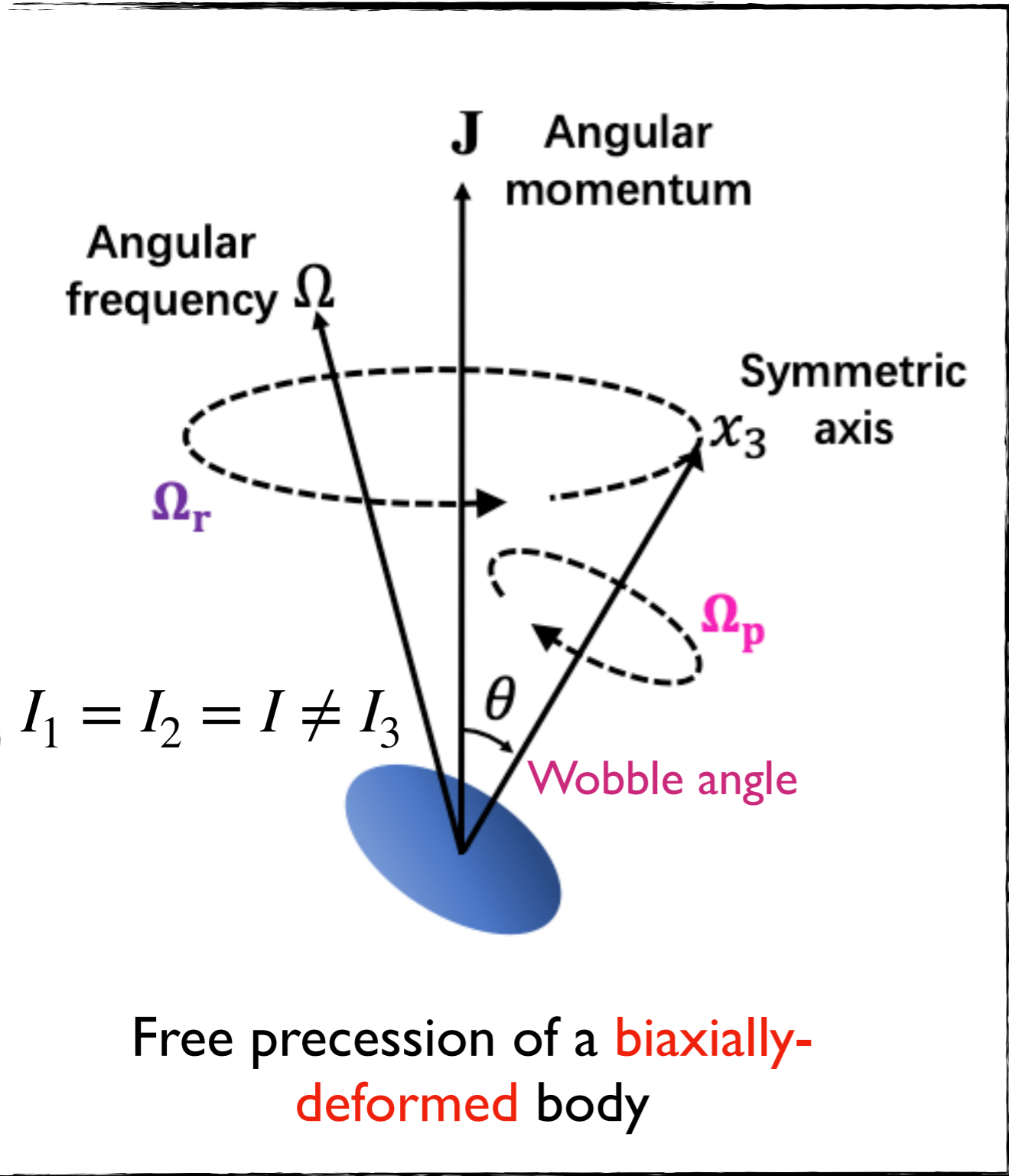
August 29, 2020

# Outline

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- Freely-precessing neutron stars (NSs)
- Continuous electromagnetic and gravitational radiation
- Indications for NS structures
- Summary

# What is free precession?



Two superimposed rotations:

$$\Omega = \Omega_r \mathbf{n}_J + \Omega_p \mathbf{n}_{x_3}$$

Rotation  
around  $\mathbf{J}$

Retrograde  
motion around  $x_3$

Free precessional angular frequency:

$$\Omega_p = \epsilon \cos \theta \Omega_r$$

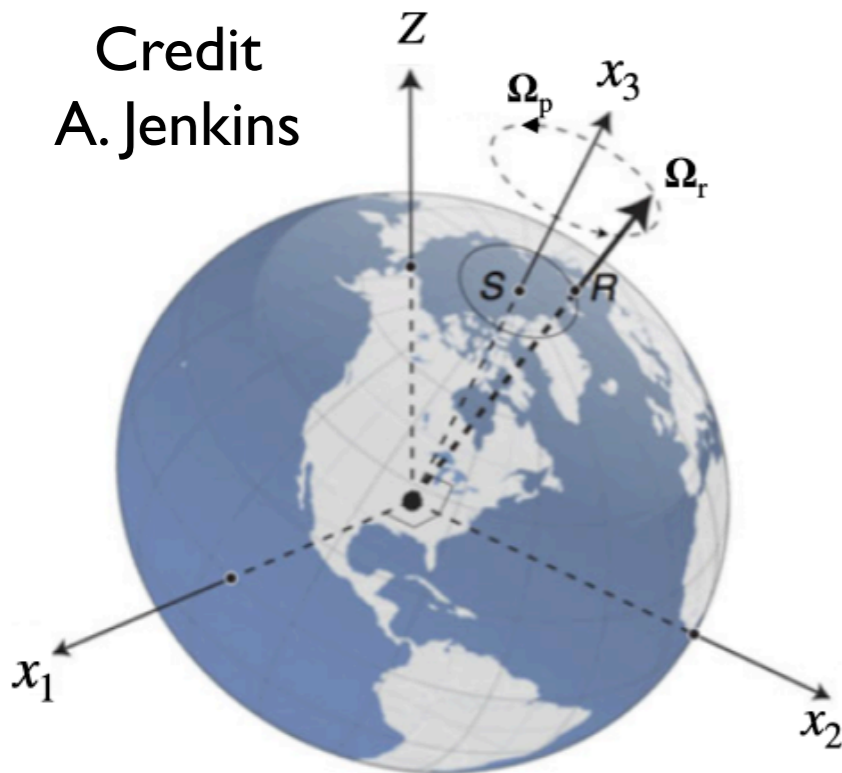
The oblateness  $\epsilon = \frac{I_3 - I}{I_3}$

# Precession of elastically-deformed bodies

Property	Earth	Neutron star
Moment of inertia: Solid crust	90%	< 5%
Moment of inertia: Liquid core	10%	> 95%
Rigidity parameter	0.7	$10^{-5}$ ; jelly
Magnetic field	Unimportant	Maybe
Free precession observed?	Yes, 14 month 'Chandler wobble'	Handful of candidates

Credit  
D. I. Jones

Credit  
A. Jenkins

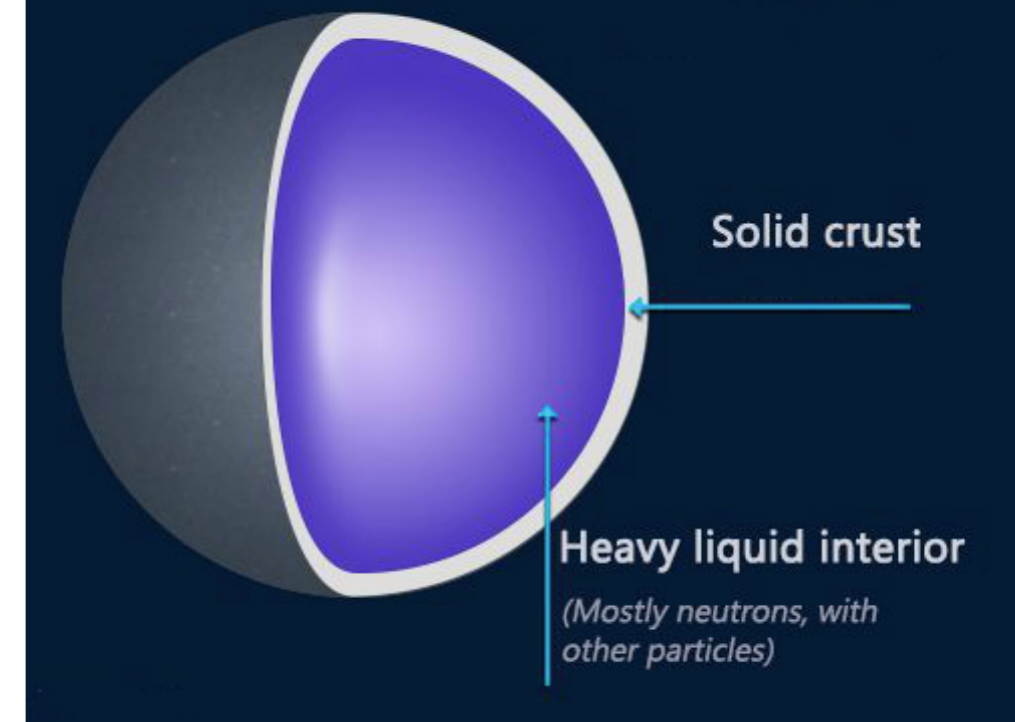


Rigidity parameter

$$\epsilon_{\text{elast}} = \eta \epsilon_{\Omega}$$

$$\text{Centrifugal deformation} \approx \frac{\Omega_r^2 R^3}{GM}$$

Conventional NS structure



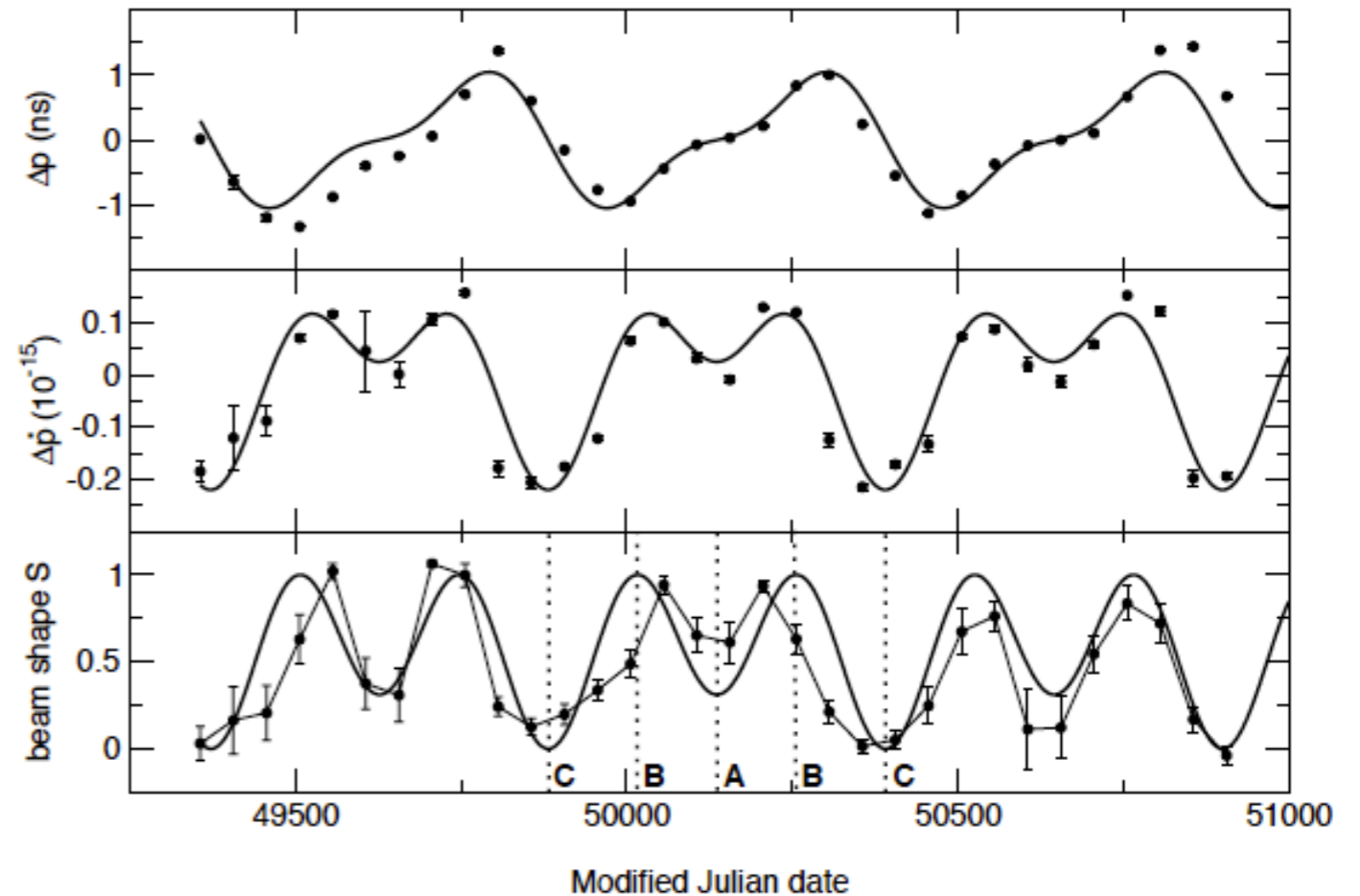
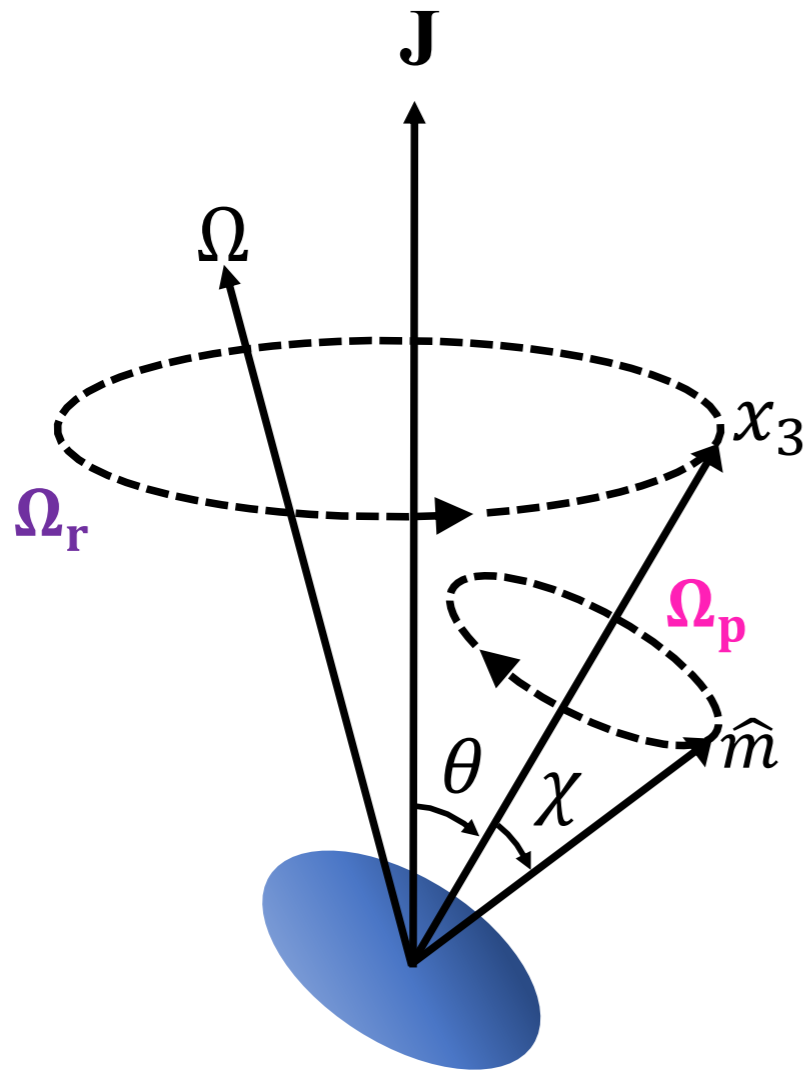
$$\frac{GM}{Rc^2} \sim 10^{-10}, \quad \epsilon_{\text{elast}} \sim \frac{1}{300}$$

Effect: latitude modulation

$$\frac{GM}{Rc^2} \sim 0.2, \quad \epsilon_{\text{elast}} < 10^{-7}$$

If exist, how to observe?

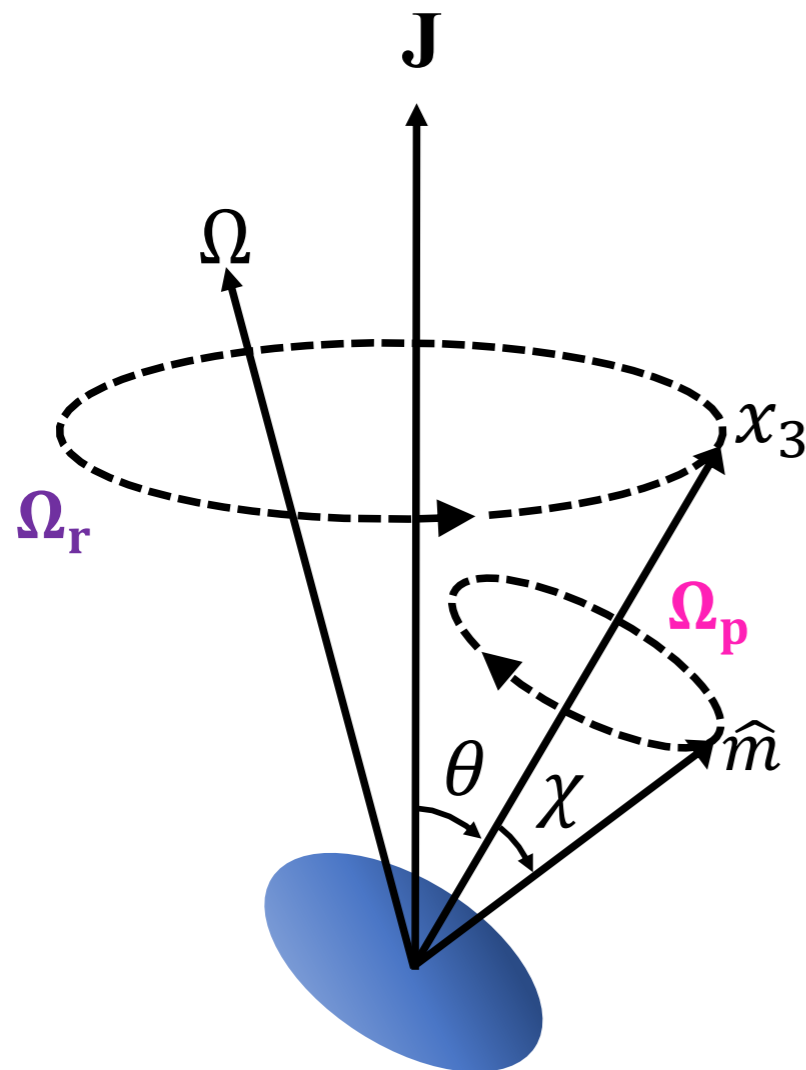
# Observational effects I: electromagnetic radiation



Timing and beam shape data for PSR B1828 –11  
(most possible evidence for free precession)

1. **Timing residuals:** the magnetic dipole undergoes two superimposed motion
2. **Pulse profile modulation:** the line of sight sweeps different region of emission cone

# Observational effects 2: continuous gravitational waves



Oblique rotator

$$h \sim \frac{2G}{c^4 r} \ddot{Q}$$

Strain of GWs | Distance to the NS | Mass quadrupole

- Emit at two harmonics, ( f, 2f)
- For small wobble angle, lower harmonic dominant

$$h \sim \frac{GI_0}{rc^4} f^2 \epsilon \theta$$

- Detectability limited by **oblateness**
- Searches are on going



# Our questions and motivation

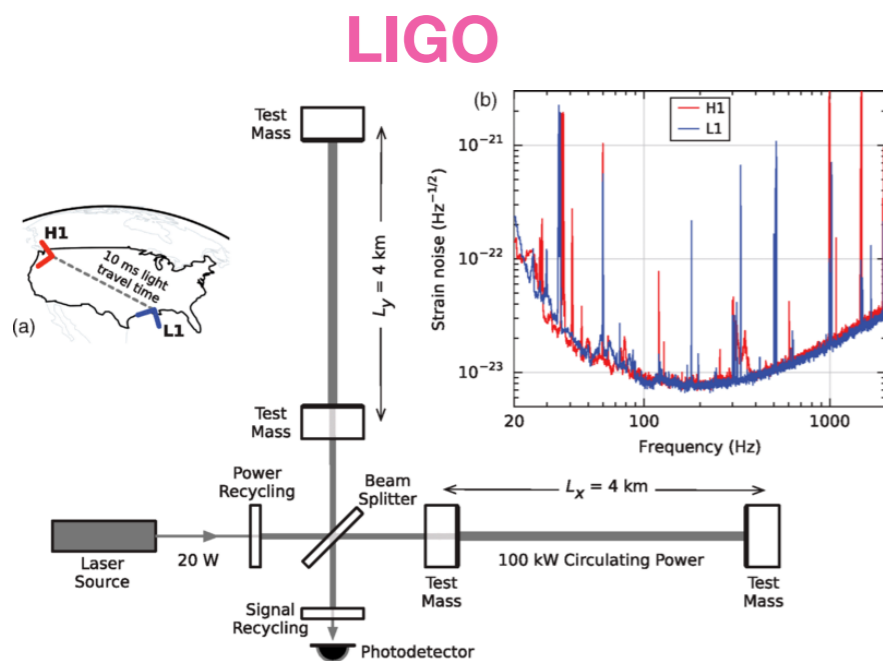
## 1. Necessary to be biaxially-deformed?

No, change of elastic field, accretion process, magnetic pressure...



Extend to fully triaxial case

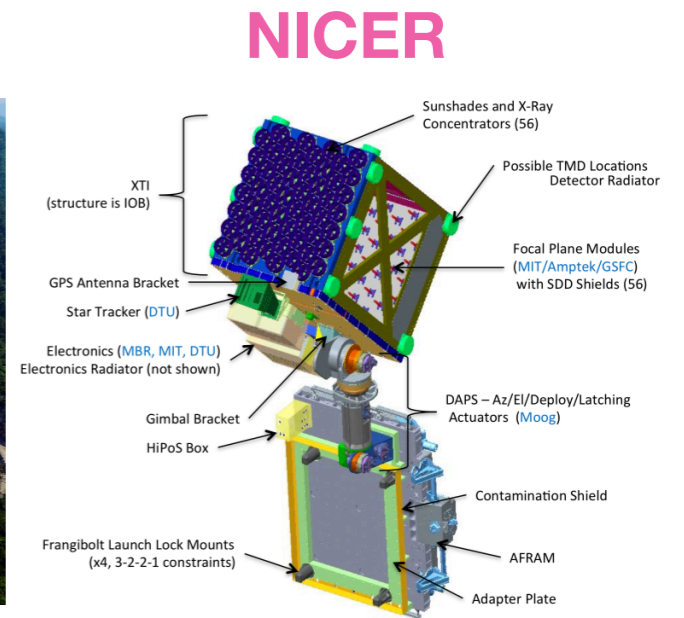
## 2. What Information from multi-messenger observation of precessing NS?



GW observation



FAST



Radio/X-ray timing

# Triaxially-deformed freely-precessing NSs

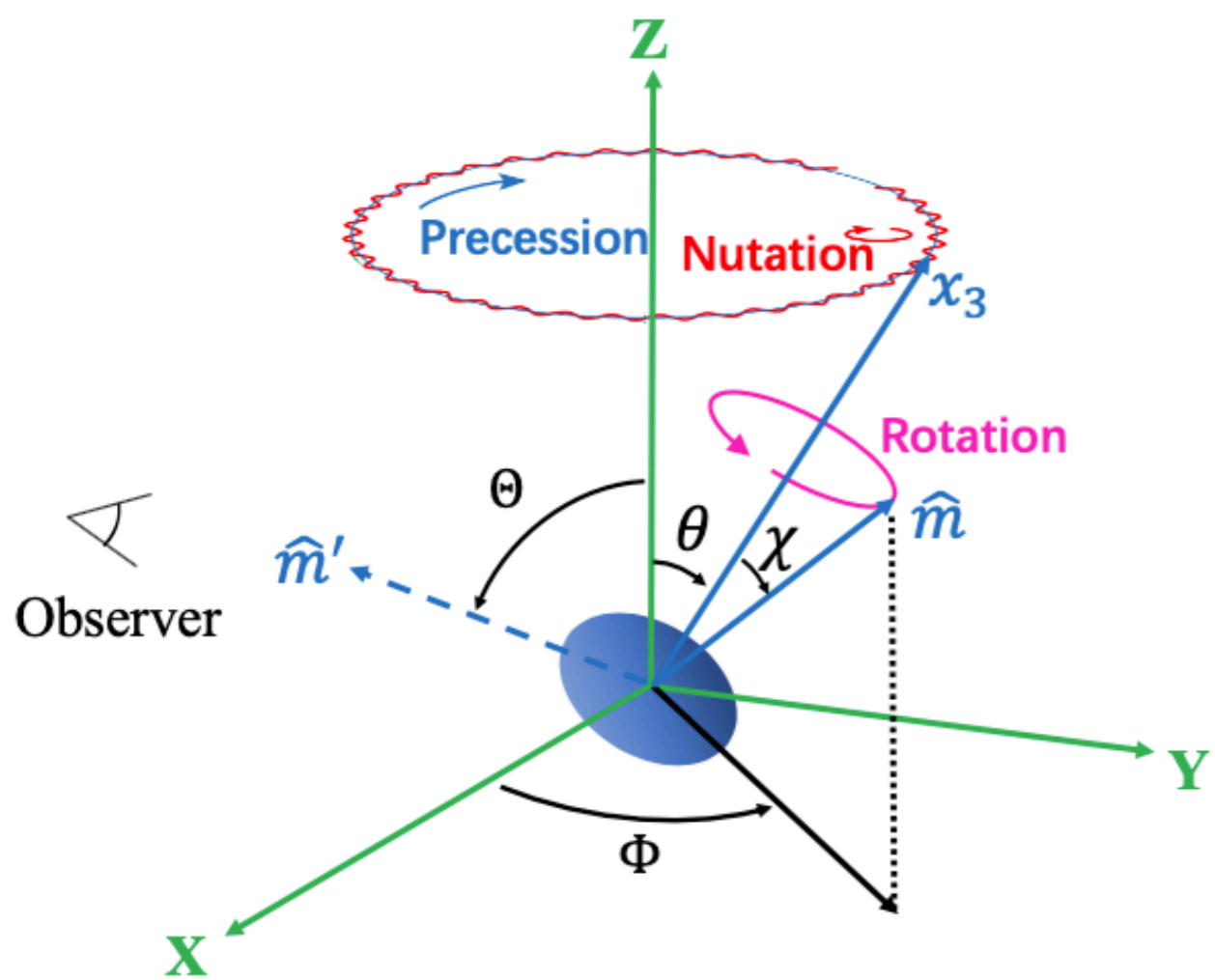
Rotation:  $\psi$

Nutation:  $\theta$

Precession:  $\phi$

Magnetic dipole:  $\hat{m}$

Magnetic inclination:  $\chi$



- Time evolution of the configuration

Analytical & Numerical

- Parameterized description of NS:

$$\epsilon \equiv \frac{I_3 - I_1}{I_3} \quad \delta \equiv \frac{I_2 - I_1}{I_3 - I_2} \quad \gamma \equiv \tan \theta_{\min}$$

Oblateness      Nonaxisymmetry      Wobble

**closely linked to the structure of NSs !**



# Estimation of $\epsilon$ , $\gamma$ and $\theta$

**Oblateness**  $\epsilon_{\text{elast}} \approx 4.9 \times 10^{-8} \left( \frac{V_c/V}{0.1} \right) \left( \frac{f_{\text{rot}}}{100\text{Hz}} \right)^2 \left( \frac{\mu}{10^{30} \text{ erg cm}^{-3}} \right) R_6^7 M_{1.4}^{-3}$

EoS dependent parameters:

Thickness  
Of the crust

Shear modulus  
of the crust

Mass - Radius

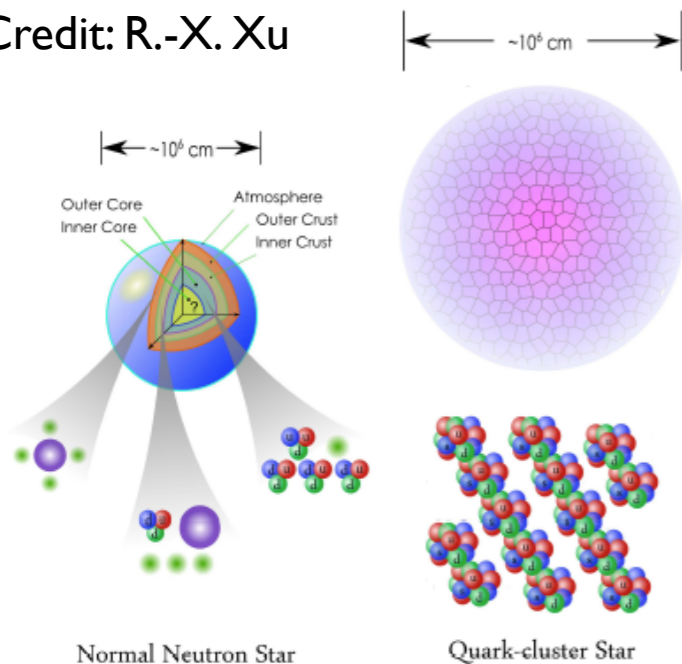
## Possible maximum oblateness

$$\epsilon_{\text{max}} \approx 10^{-6} \left( \frac{V_c/V}{0.1} \right) \left( \frac{\mu}{10^{29} \text{ erg cm}^{-3}} \right) \left( \frac{\sigma_{\text{break}}}{0.1} \right)$$

Breaking strain, highly uncertain

Recent lattice study: 0.04-0.1

Credit: R.-X. Xu



Conventional NS

$$\epsilon_{\text{max}} \sim 10^{-6}$$

Solid quark star

$$\epsilon_{\text{max}} \sim 10^{-4}$$

## Nonaxisymmetry

$$\delta \equiv \frac{I_2 - I_1}{I_3 - I_2} \quad \text{Possible to be any value}$$

## Wobble angle

$$\theta_{\text{max}} \approx 0.45 \left( \frac{100\text{Hz}}{f_r} \right)^2 \left( \frac{\sigma_{\text{break}}}{10^{-3}} \right) M_{1.4} R_6^{-3}$$

# Modulated timing signals

- Modulated timing signal of pulsar

$$\Delta\Phi = F(\psi, \theta, \phi)$$

Phase residual  
due to precession

=

Function of NS  
configuration

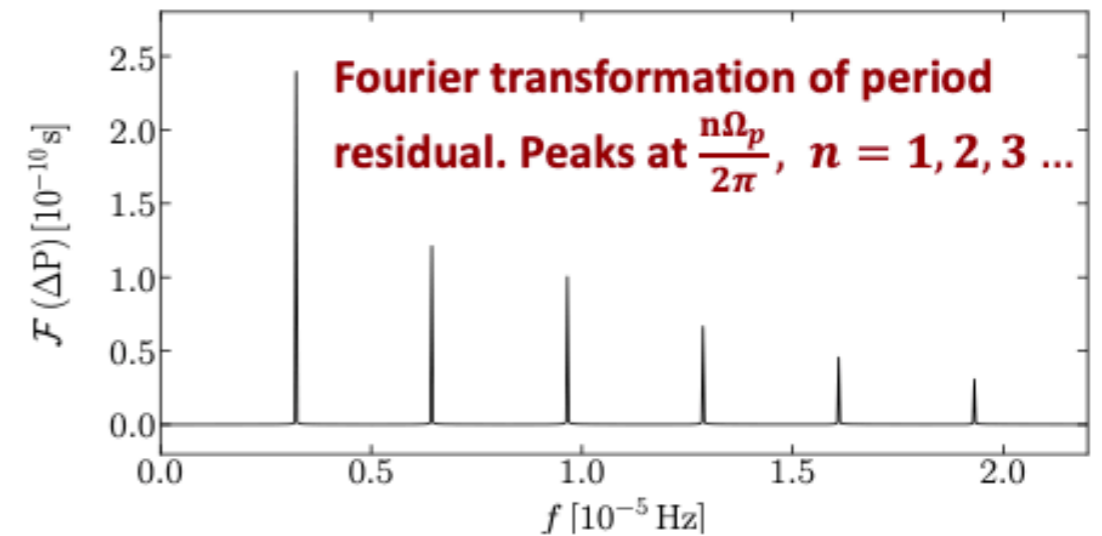
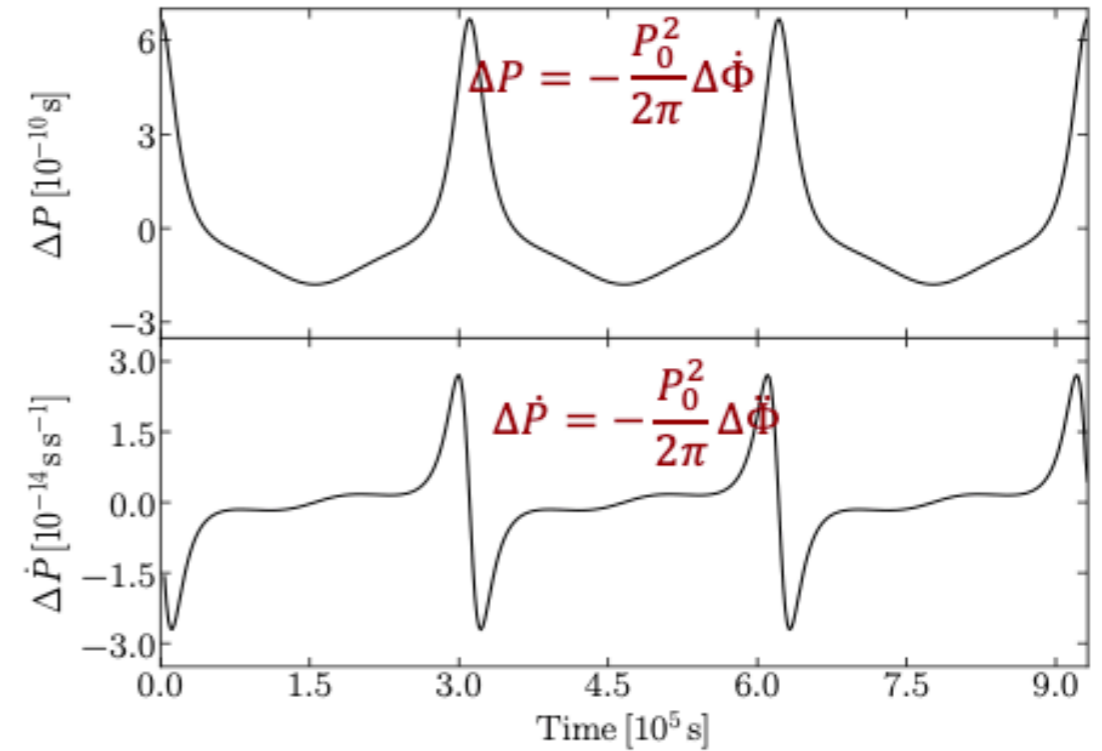
- Small wobble angle case

$$\Delta P \approx \frac{P_0^2}{2\pi} \Omega_p \gamma (\delta + 1) \cot \chi \cos(\Omega_p t) + \frac{P_0^2}{4\pi} \Omega_p \gamma^2 (1 + 2 \cot^2 \chi) \cos(2\Omega_p t)$$

1.  $\Omega_p$  is the free precession angular frequency

2. To second order expansion:  $\Omega_p$  and  $2\Omega_p$

## Large wobble angle case

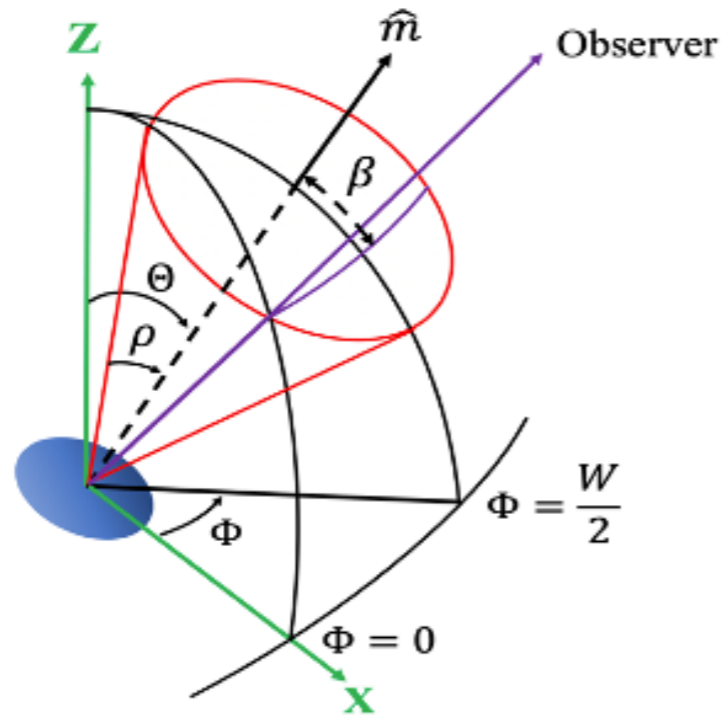


$$P_0 = 10 \text{ ms} \quad P_{\text{fp}} = 3.1 \times 10^5 \text{ s}$$

$$\epsilon = 4.9 \times 10^{-8} \quad \delta = 0.1 \quad \theta \in (0.79, 0.84)$$

# Modulated pulse signals

- Simple Cone model



- Modulated pulse widths

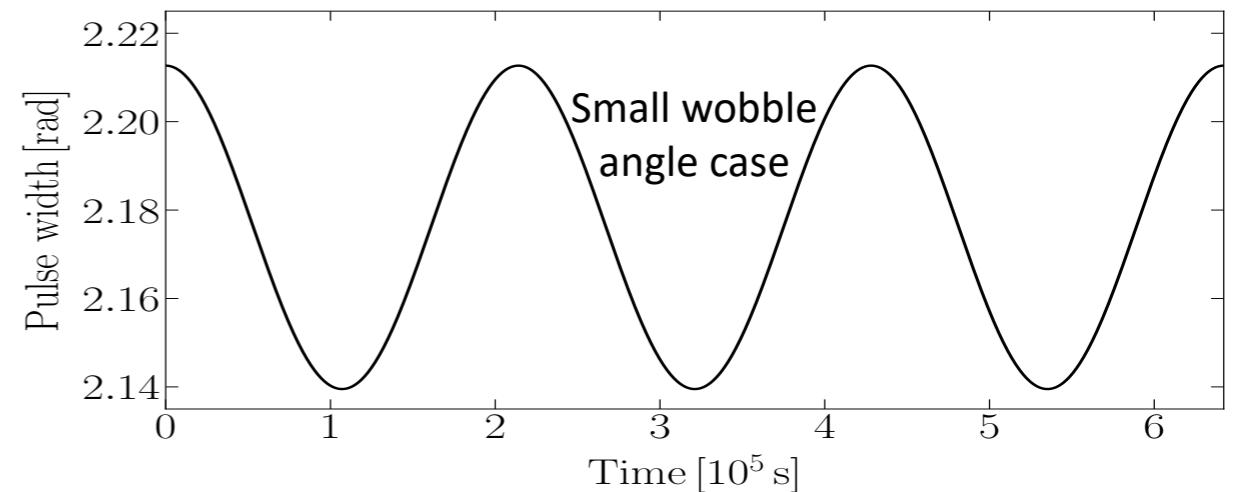
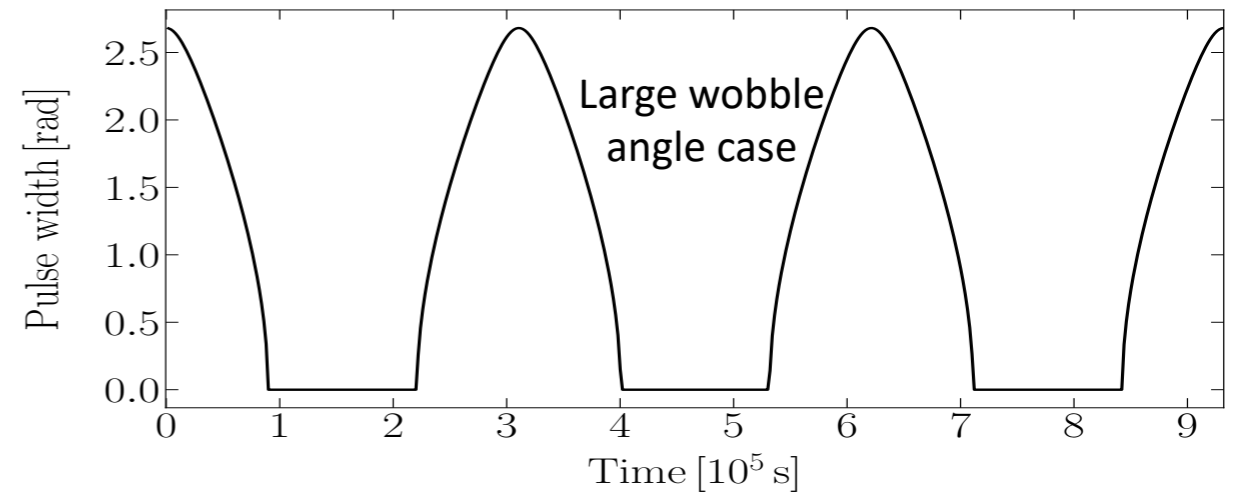
$$W = G(\Theta, \rho), \quad \Theta = H(\psi, \theta, \phi, \chi, \iota)$$



Pulse-width  
due to precession

=

Function of NS  
configuration, opening  
angle of emission  
cone, line of sight

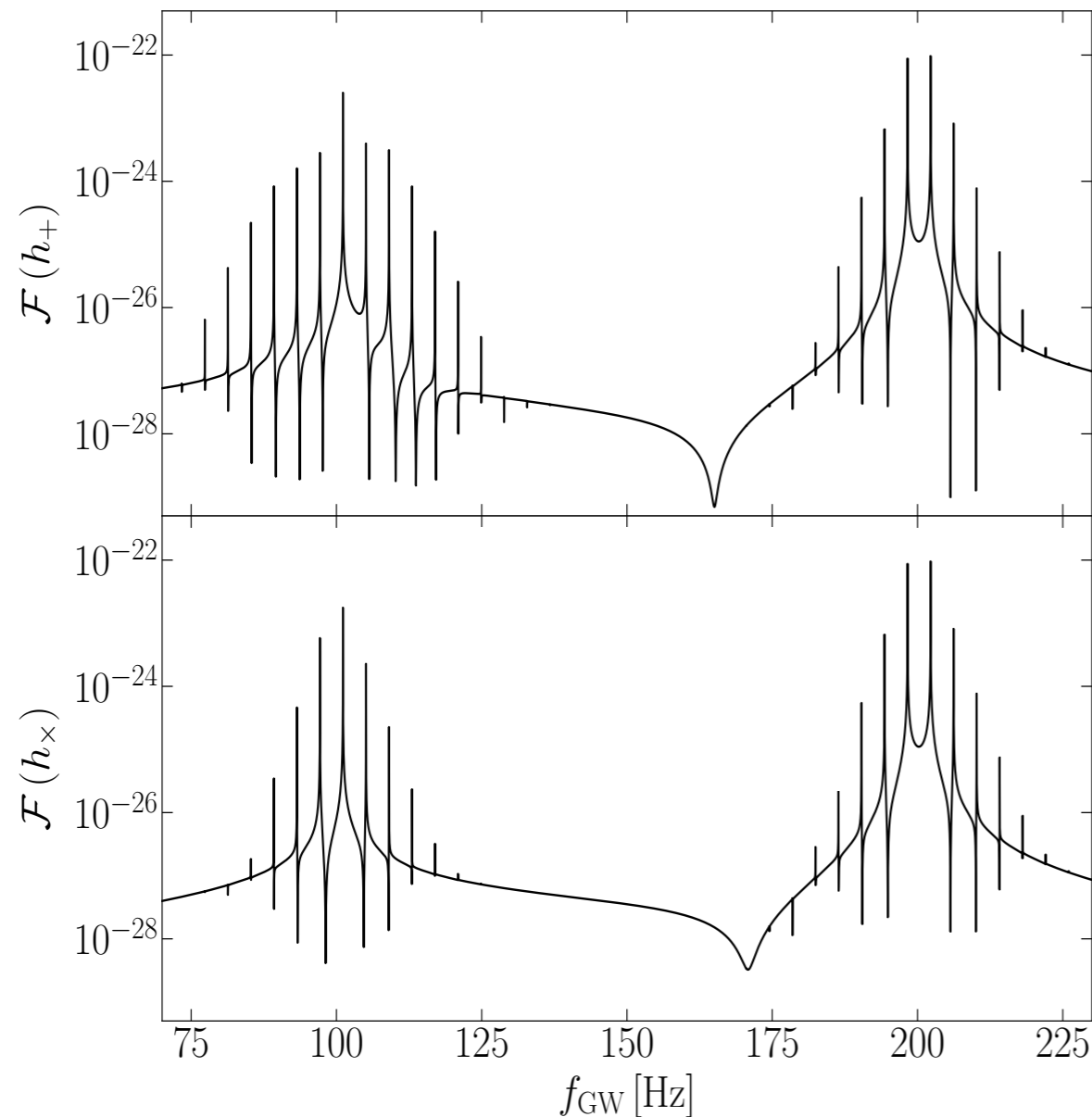


1. Large wobble angle: may lose the emission during precession

2. Small wobble angle: modulation is much weaker

# Continuous gravitational waves

- Generic waveform



$$h_{ij}^{\text{TT}} = \frac{2G}{c^4 r} \frac{d^2 I_{ij}}{dt^2} = -\frac{2G}{rc^4} \mathcal{R}_{ik} \mathcal{R}_{jl} A_{kl}$$

- Small  $\gamma$  and  $\delta$  case (to second order expansion)

First order lines at  $\Omega_r + \Omega_p$  and  $2\Omega_r$

$$A_{\times}^1 = 1.0 \times 10^{-28} \gamma \sin \iota \left( \frac{\epsilon}{4.9 \times 10^{-8}} \right) \left( \frac{f_r}{100 \text{ Hz}} \right)^2 \left( \frac{10 \text{ kpc}}{r} \right)$$

$$A_{\times}^2 = 2.1 \times 10^{-28} \delta \cos \iota \left( \frac{\epsilon}{4.9 \times 10^{-8}} \right) \left( \frac{f_r}{100 \text{ Hz}} \right)^2 \left( \frac{10 \text{ kpc}}{r} \right)$$

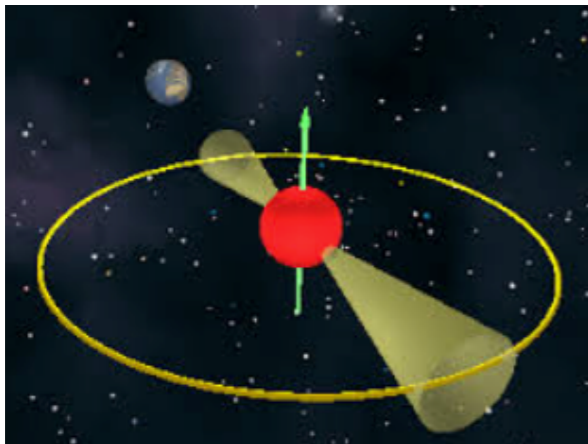
Second order lines at

$$2 \left( \Omega_r + \Omega_p \right), \Omega_r - \Omega_p, \Omega_r - 3\Omega_p, \text{ and } 2 \left( \Omega_r - \Omega_p \right)$$

# Multimessenger observation: Extraction of physical parameters

First order

Second order

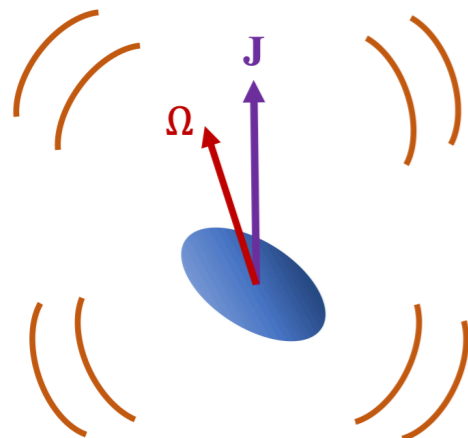


EM

$$\Delta P_1 \sim \gamma \cot \chi P_0^2 \Omega_p$$

$$\Delta P_2 \sim \gamma \delta \cot \chi P_0^2 \Omega_p$$

$$\Delta P_3 \sim \gamma^2 (1 + 2 \cot^2 \chi) P_0^2 \Omega_p$$



GW

$$A_{\times}^1 \sim \frac{I_0 \gamma \epsilon f_r^2 \sin \iota}{r}$$

$$A_{\times}^1 \sim \frac{I_0 \delta \epsilon f_r^2 \cos \iota}{r}$$

Dependent on

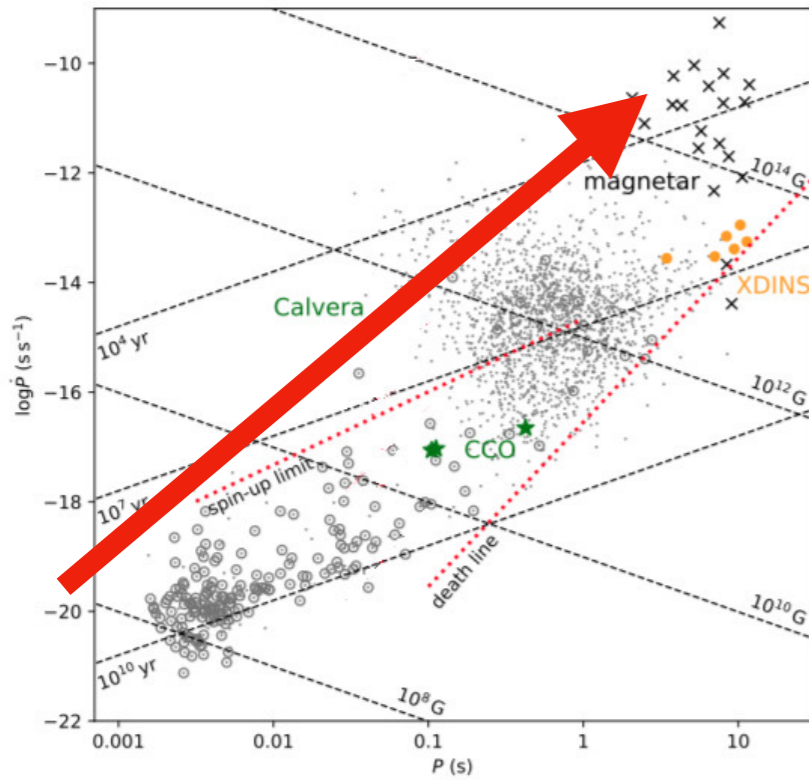
$\gamma^2$ ,  $\delta^2$ , and  $\gamma\delta$

Multi-messenger observation: extract  $\epsilon$ ,  $\delta$ ,  $\gamma$ ,  $\chi$ ,  $\iota$ , and  $\Omega_p$



# Future possible work

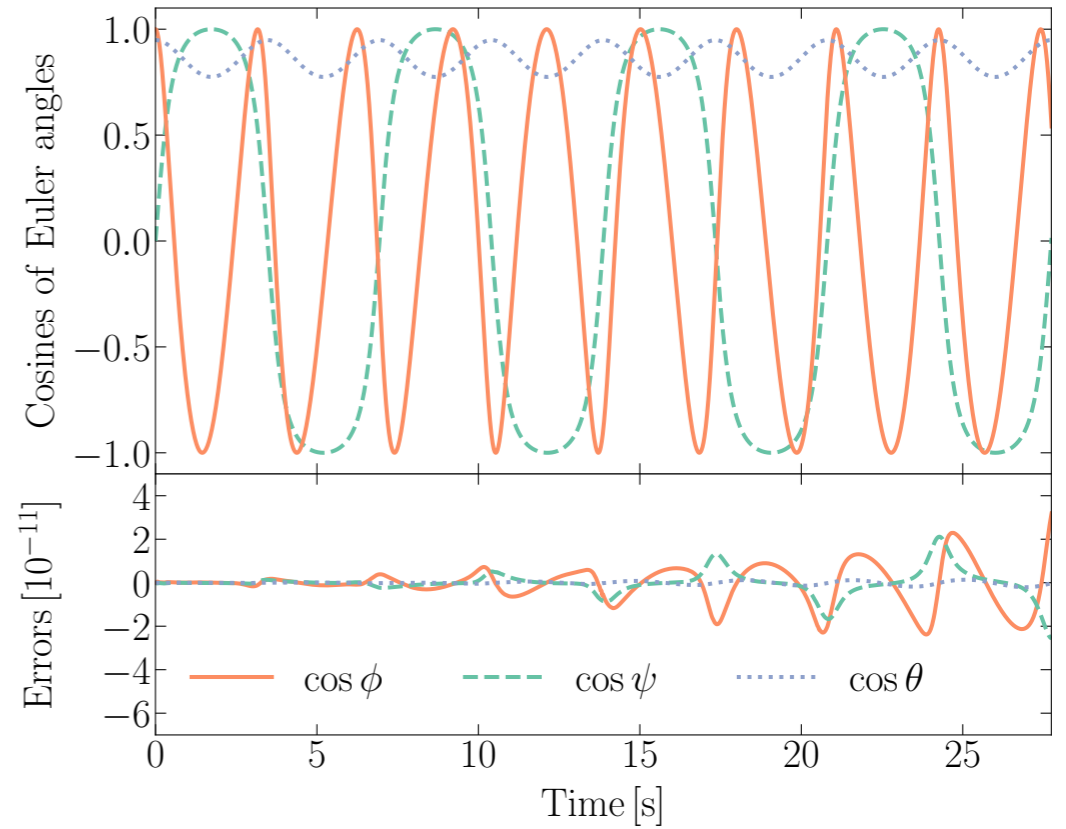
Credit: Yunyang Li



Numerical approach using quaternions

$$I_1/I_3 = 1/3$$

$$I_2/I_3 = 2/3$$



$$\Delta\Phi \sim \frac{1}{\pi} \cot\chi \frac{\theta}{\epsilon^2} \frac{P_0}{\tau_e}$$

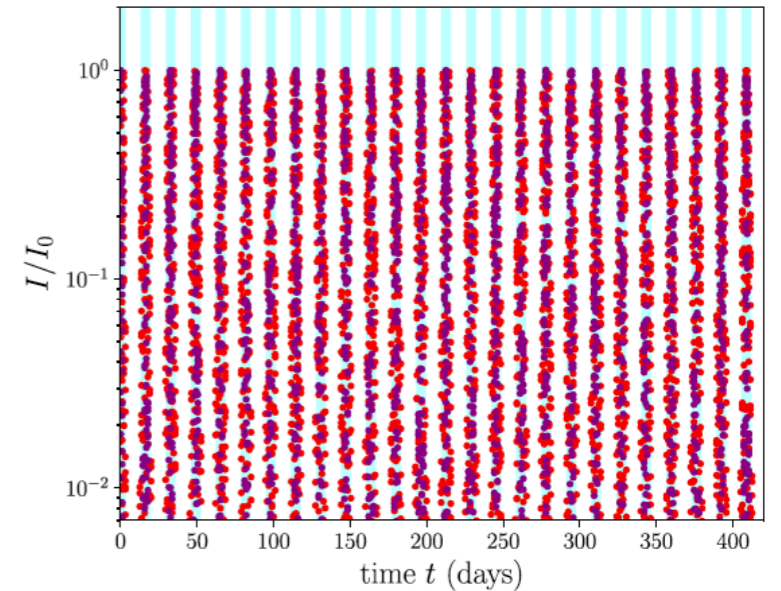
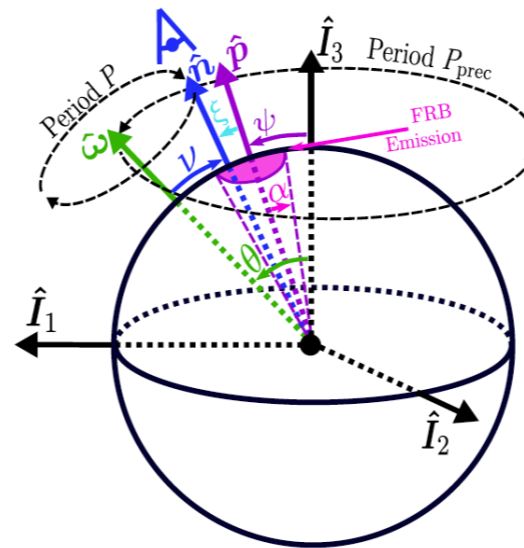
Spin down time scale

Millisecond pulsars: **unimportant**

Young pulsars: **greatly amplify** residuals

Magnetar: **must consider** and

$$\epsilon_{\text{mag}} = \beta \frac{R^4 B_\star^2}{GM^2} = 1.9 \times 10^{-6} \beta \left( \frac{B_\star}{10^{15} \text{G}} \right)^2 \frac{R_6^4}{M_{1.4}^2}$$



Precession of magnetar to explain  $\sim 16$  day periodicity of FRB 180916.J0158+65 (Zanazzi & Lai, 2020)



# Summary

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- Triaxially-deformed NSs: **new features**
- Multi-messenger observation: **valuable information on equation of state**
- Radio/X-ray timing and GWs searches: **on going**

**Thanks!**