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北京大學科維理天文與天體物理研究所



X-Ray Pulse Profile of Scalarized Neutron Stars

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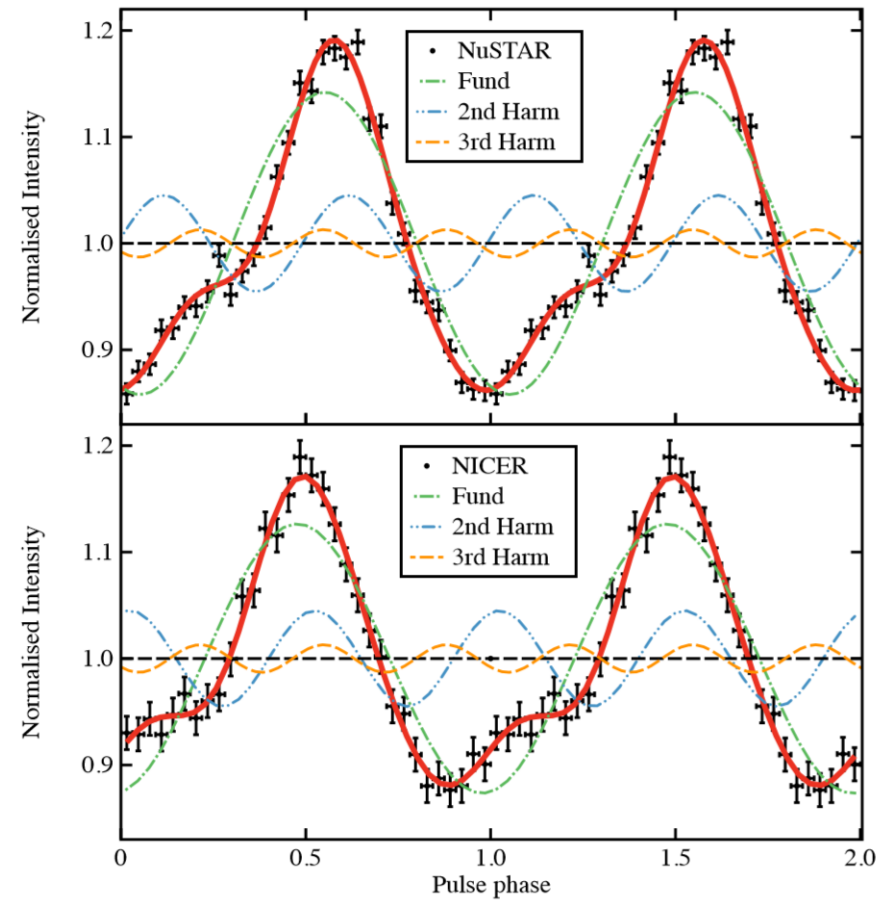
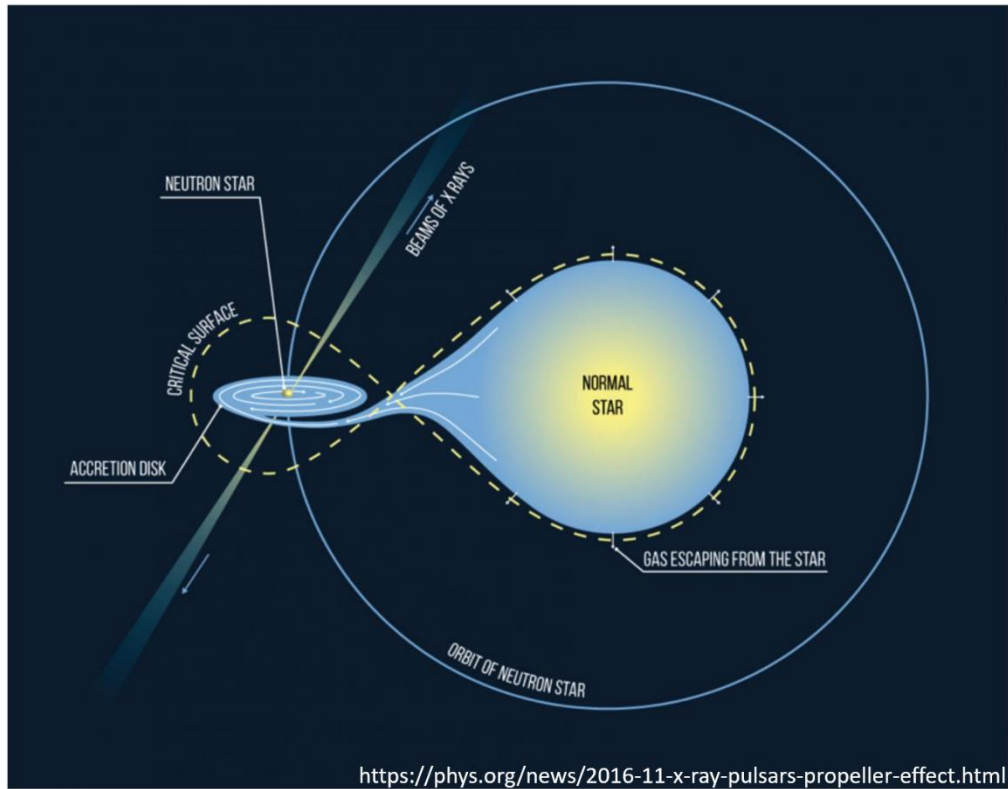
Based on Rui Xu, Yong Gao, and Lijing Shao, arXiv:2007.10080

For **FAST/Future Pulsar Symposium 9** at Xiamen University, Xiamen
August 28-30, 2020

Outline

- Observed flux from X-ray pulsars
- Spacetime around scalarized neutron stars (NSs)
- Pulse profile of scalarized NSs

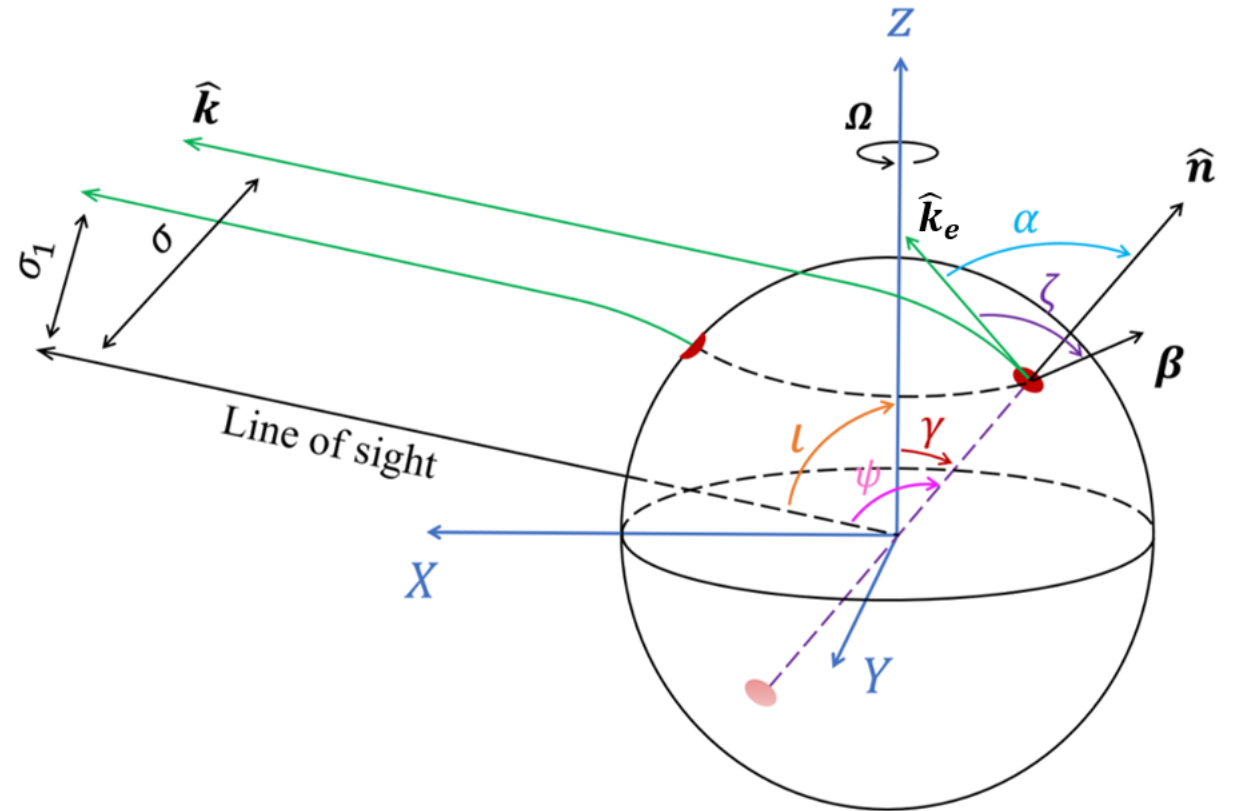
X-ray pulsar



A. Sanna, et al. , A&A 617, L8 (2018)

Hot spot emission model

- ❖ A pair of small hot spots
- ❖ A far-away observer



R. Xu, Y. Gao, and L. Shao, arXiv:2007.10080

Hot spot emission model

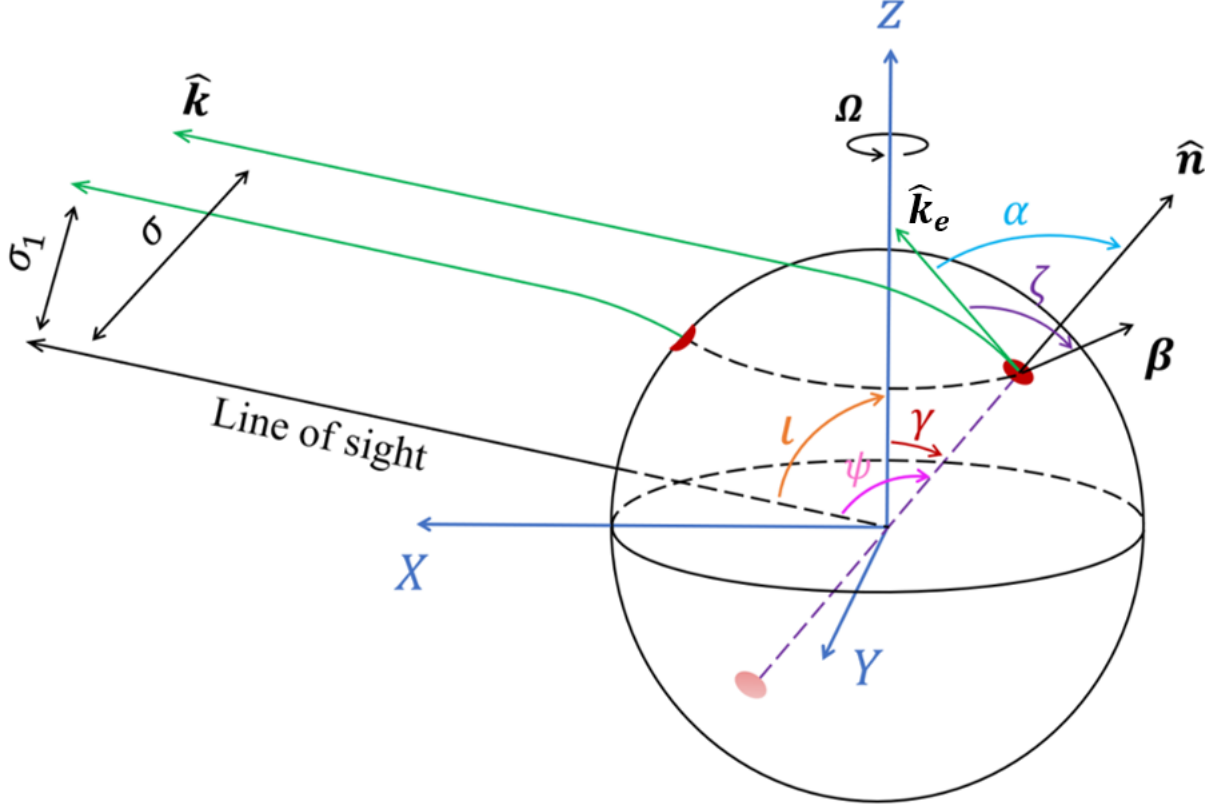
- ❖ A pair of small hot spots
- ❖ A far-away observer

$$dF = I_o(E_o, \hat{\mathbf{k}}) dE_o d\Omega = I_o(E_o, \hat{\mathbf{k}}) dE_o \frac{\sigma d\sigma d\lambda}{D^2}$$

H. O. Silva, N. Yunes, Phys. Rev. D 99, 044034 (2019)

gravitational redshift $\longrightarrow I_o/\nu_o^3 = I_e/\nu_e^3$

bending of light $\longrightarrow d\sigma d\lambda \leftrightarrow dS = R^2 \sin \psi d\psi d\lambda$



R. Xu, Y. Gao, and L. Shao, arXiv:2007.10080

Spherical NS

$$ds^2 = -B(r)dt^2 + A(r)dr^2 + r^2(d\theta^2 + \sin^2 \theta d\phi^2)$$

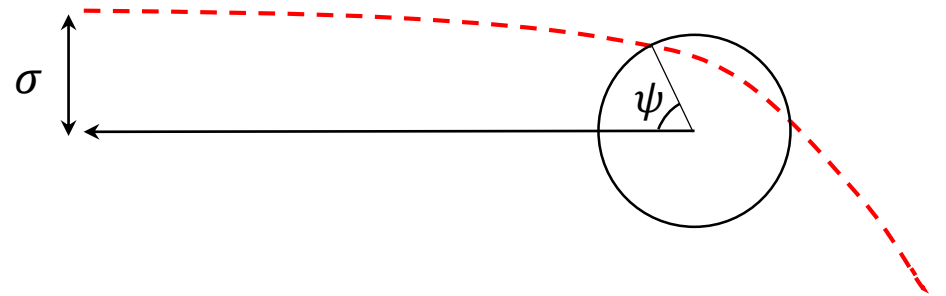
gravitational redshift $\frac{\nu_o}{\nu_e} = \sqrt{B(R)}$

bending of light $\psi = \int_{\infty}^R \frac{d\theta}{dr} dr = \psi(\sigma)$

$$\frac{dt}{dp} = \frac{1}{B}$$

$$-B \left(\frac{dt}{dp} \right)^2 + A \left(\frac{dr}{dp} \right)^2 + r^2 \left(\frac{d\theta}{dp} \right)^2 = 0$$

$$\frac{d\theta}{dp} = \frac{\sigma}{r^2}$$



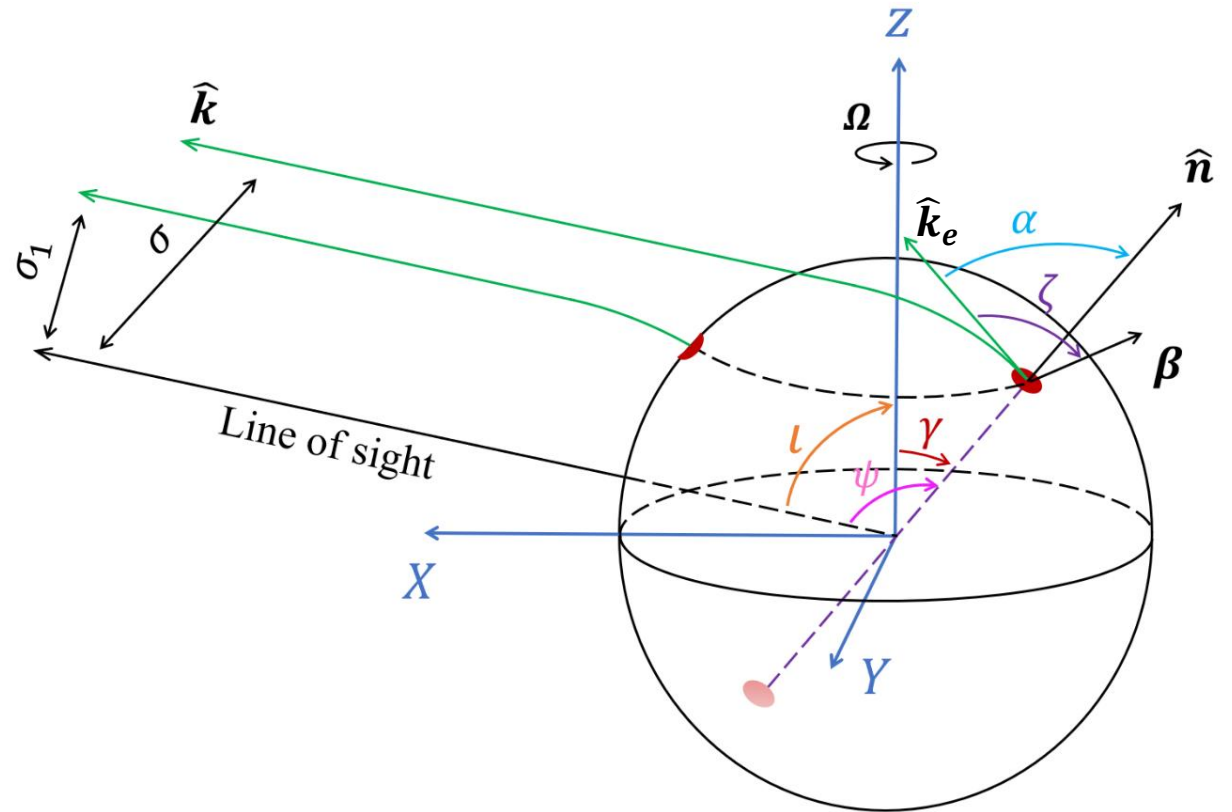
Spherical NS

$$ds^2 = -B(r)dt^2 + A(r)dr^2 + r^2(d\theta^2 + \sin^2 \theta d\phi^2)$$

$$\begin{aligned} dF &= I_o(E_o, \hat{\mathbf{k}}) dE_o \frac{\sigma d\sigma d\lambda}{D^2} \\ &= B^2(R) I_e(E_e, \hat{\mathbf{k}}_e) dE_e \frac{\sigma d\sigma}{\sin \psi d\psi} \frac{dS}{R^2 D^2} \end{aligned}$$

Isotropic radiation: $I_e(E_e, \hat{\mathbf{k}}_e) = I_e(E_e)$

$$\cos \psi = \cos \iota \cos \gamma + \sin \iota \sin \gamma \cos \phi$$

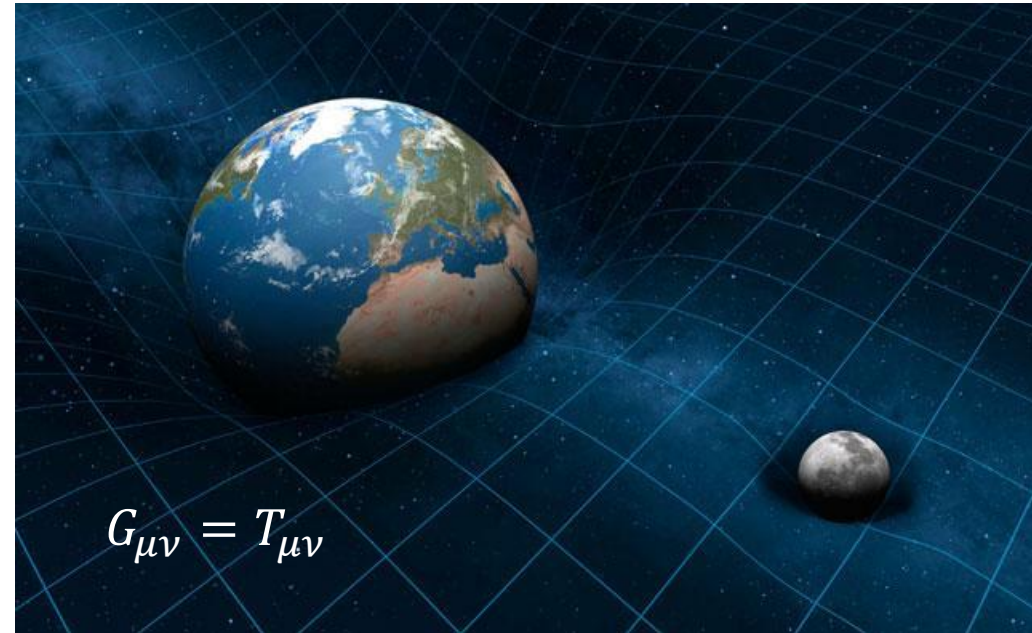


R. Xu, Y. Gao, and L. Shao, arXiv:2007.10080

Spherical NS

$$ds^2 = -B(r)dt^2 + A(r)dr^2 + r^2(d\theta^2 + \sin^2 \theta d\phi^2)$$

General Relativity (GR): $B(r) = 1 - \frac{r_s}{r}, A(r) = \frac{1}{1 - \frac{r_s}{r}}$



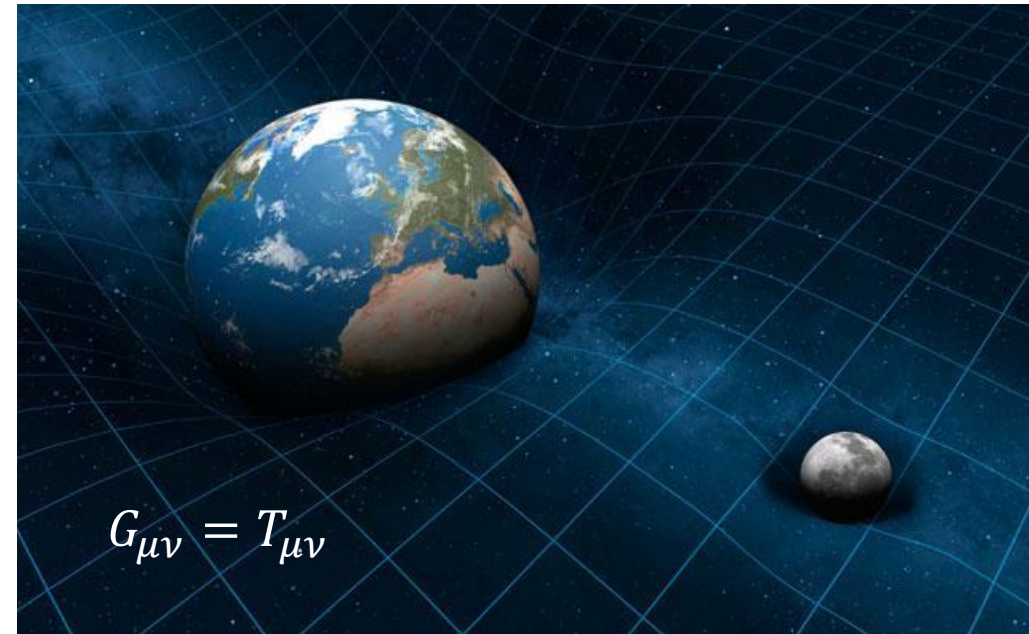
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General Relativity (GR): $B(r) = 1 - \frac{r_s}{r}, A(r) = \frac{1}{1 - \frac{r_s}{r}}$

Scalar-tensor theory: $B(r) = ?, A(r) = ?$

- Cosmology: dark energy, inflation
- Quantum gravity



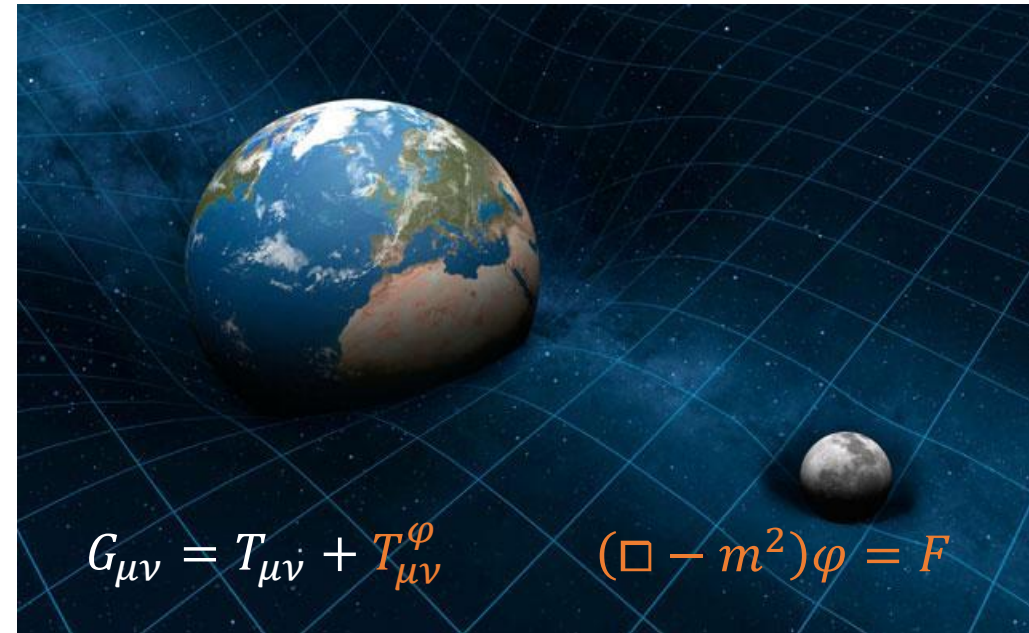
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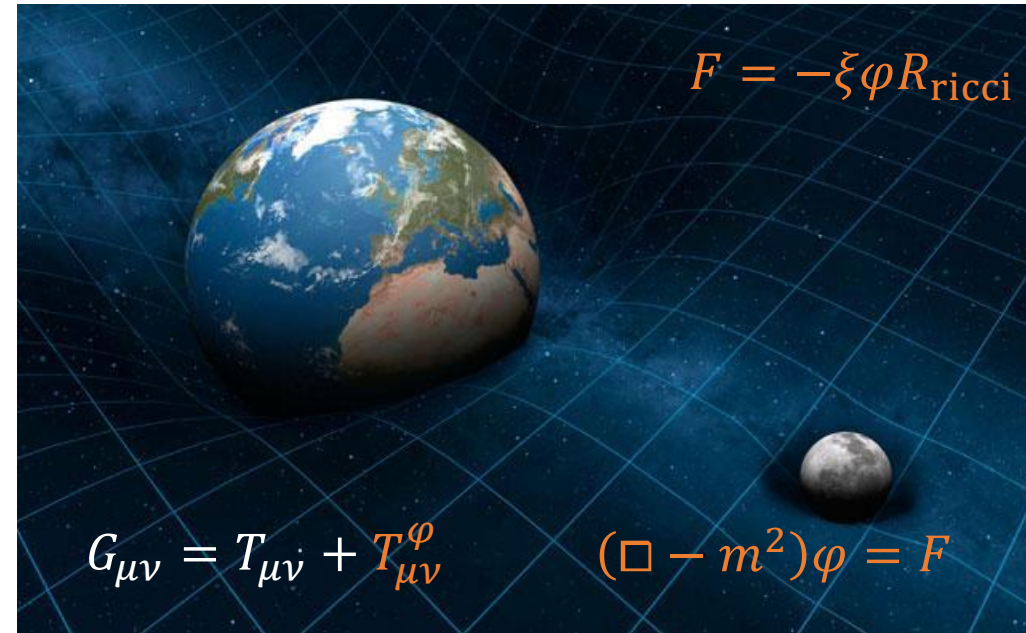
Spherical NS

$$ds^2 = -B(r)dt^2 + A(r)dr^2 + r^2(d\theta^2 + \sin^2\theta d\phi^2)$$

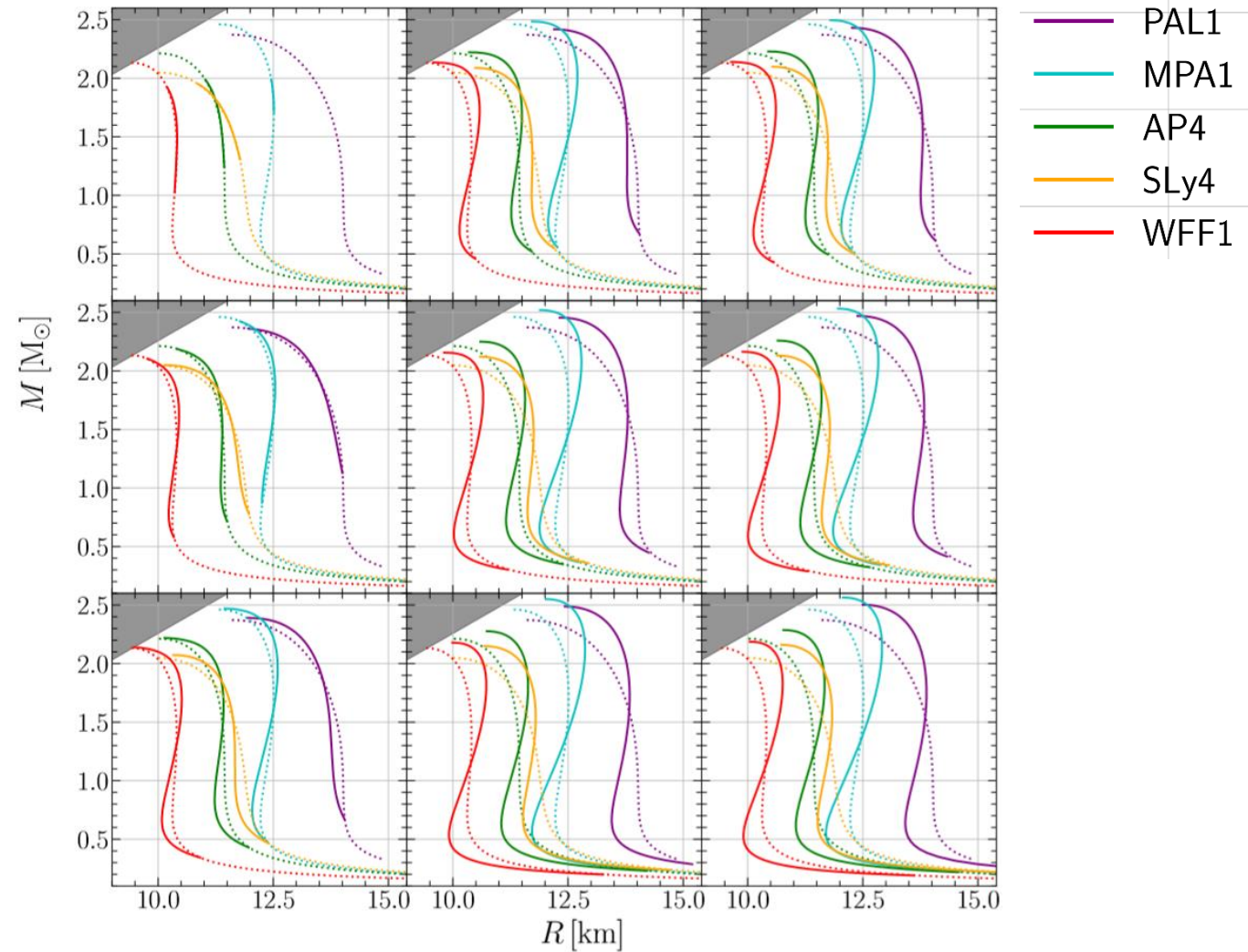
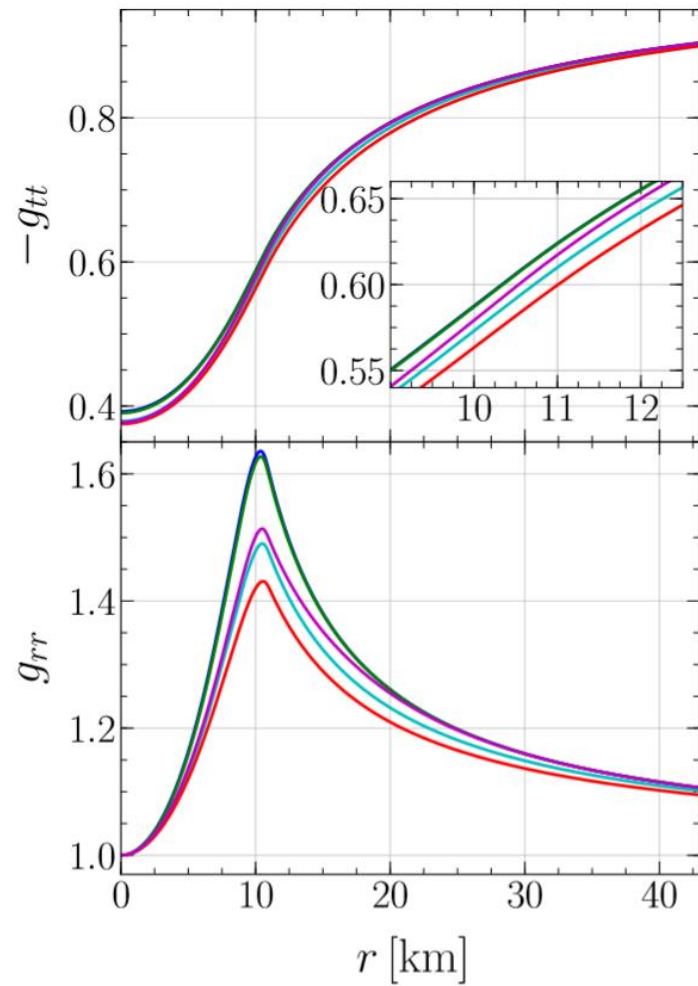
General Relativity (GR): $B(r) = 1 - \frac{r_s}{r}, A(r) = \frac{1}{1 - \frac{r_s}{r}}$

Scalar-tensor theory: $B(r) = ?, A(r) = ?$

- Cosmology: dark energy, inflation
- Quantum gravity



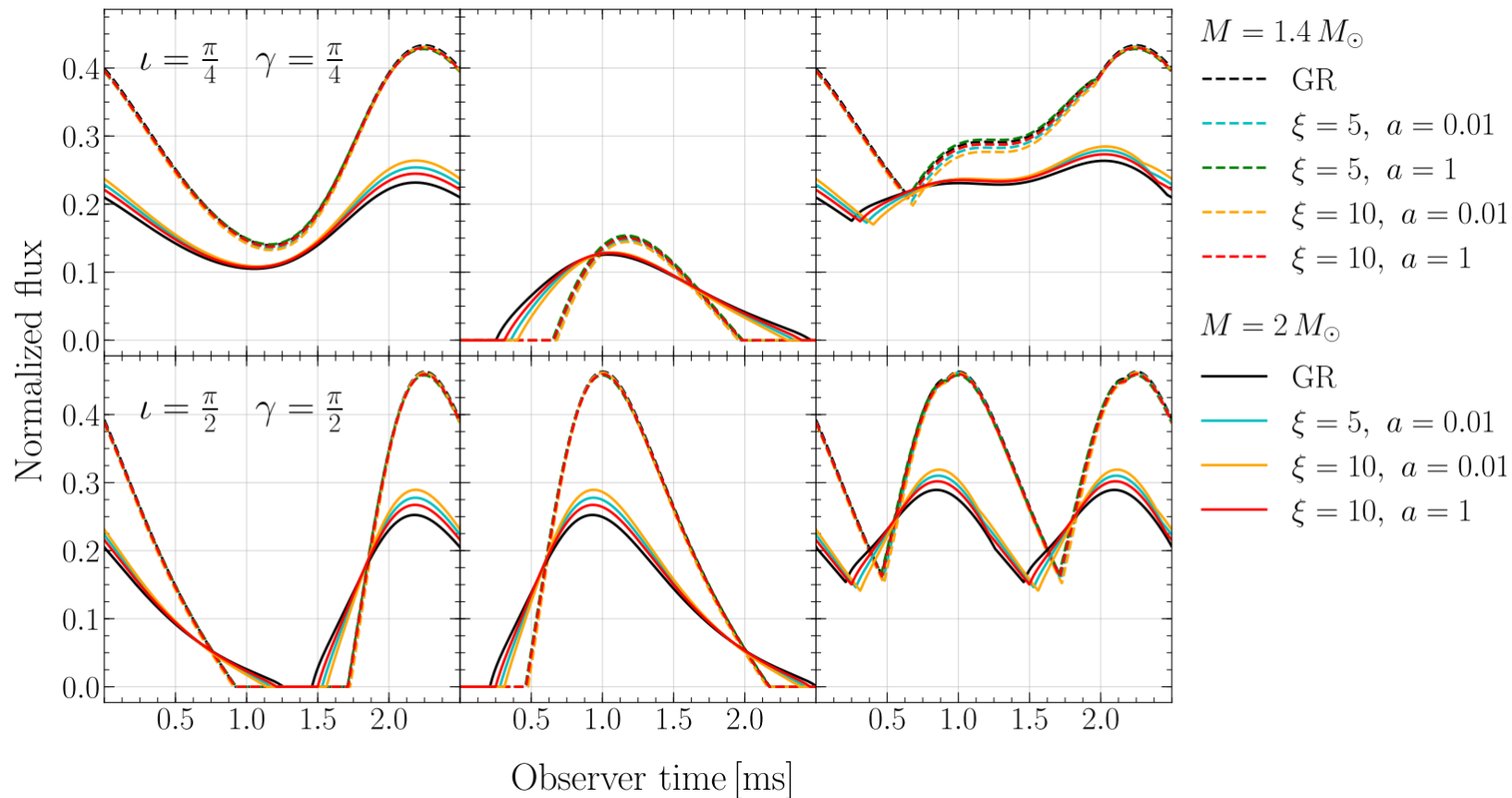
Spherical NSs in scalar-tensor theory



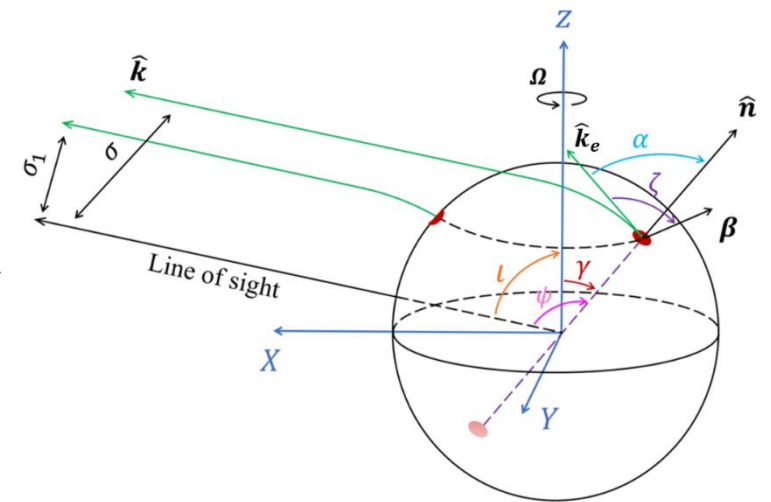
Pulse profiles

$$dF = B^2(R)I_e(E_e, \hat{\mathbf{k}}_e)dE_e \frac{\sigma d\sigma}{\sin \psi d\psi} \frac{dS}{R^2 D^2} \propto \frac{\sigma d\sigma}{\sin \psi d\psi}$$

Isotropic radiation



$$m \approx a \times 1.97 \times 10^{-11} \text{eV}$$



Summary

- X-ray pulse profiles contain information about the spacetime around NSs (gravitational redshift, bending of the light trajectories).
- Scalarized NSs generate spacetime different from the Schwarzschild metric.
- Distinguishing scalarized NSs from those in GR using X-ray pulsar observations is challenging as the differences in the metrics are small for coupling constants of order unity.

Spherical NS

$$ds^2 = -B(r)dt^2 + A(r)dr^2 + r^2(d\theta^2 + \sin^2 \theta d\phi^2)$$

$$\begin{aligned} dF &= I_o(E_o, \hat{\mathbf{k}}) dE_o \frac{\sigma d\sigma d\lambda}{D^2} \\ &= B^2(R) I_e(E_e, \hat{\mathbf{k}}_e) dE_e \frac{\sigma d\sigma}{\sin \psi d\psi} \frac{dS}{R^2 D^2} \end{aligned}$$

