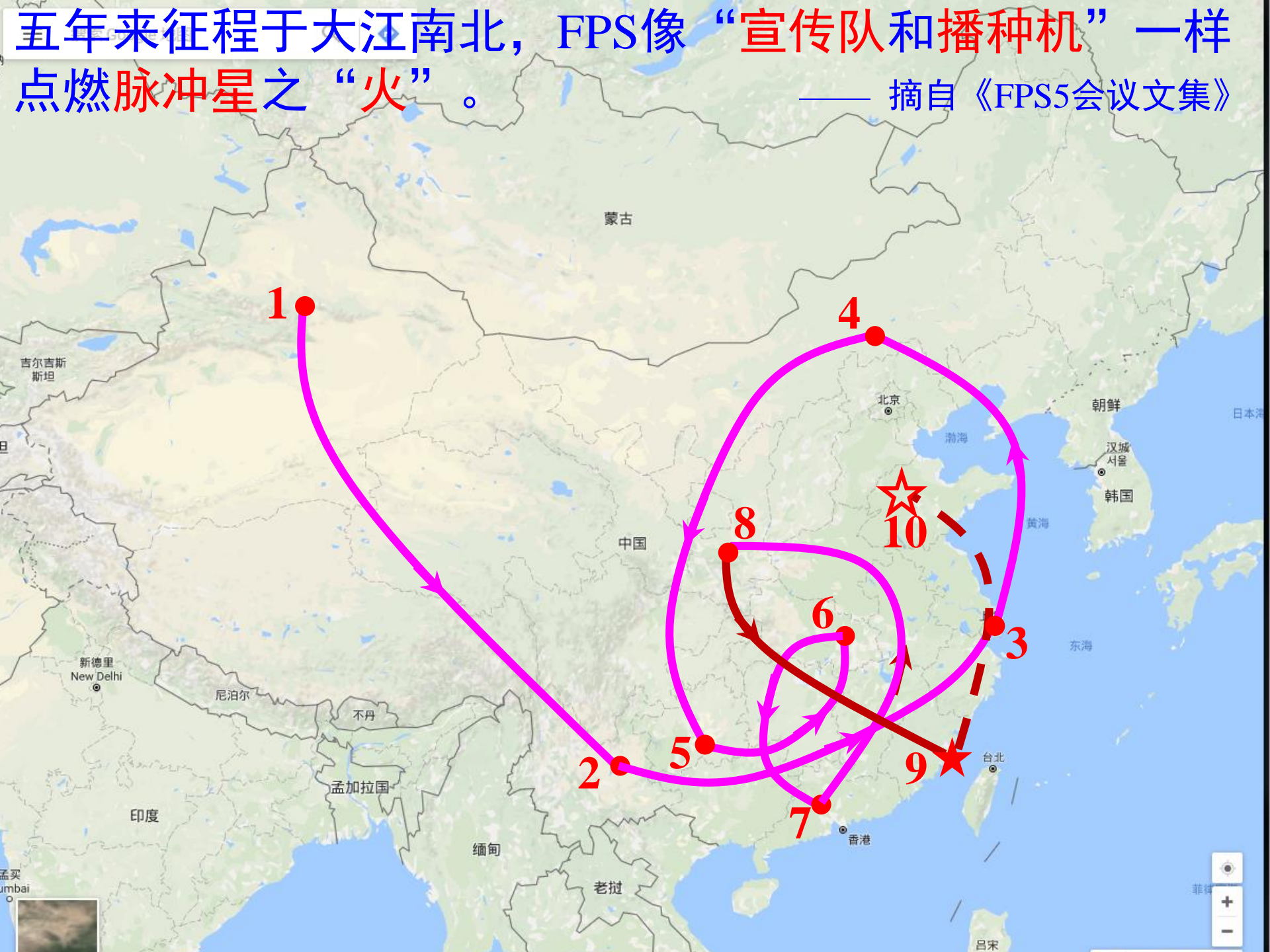


五年来征程于大江南北，FPS像“宣传队和播种机”一样
点燃脉冲星之“火”。
——摘自《FPS5会议文集》



CMB 公园路线图

Renxin Xu (徐仁新)^{1,2}

¹School of Physics, ²KIAA; PKU

(北京大学物理学院)

“FPS9”

Aug. 28-30, 2020; Xiamen Aqua Resort

A key to all (PSR/FRB/GRB/SNE...): what's CBM?

- **What** if normal baryonic matter is *compressed*?

Normal matter is squeezed so great that 2-flavoured nuclei come in close contact during a supernova!





Wellcome to Neutron Star

Strangeon Star



Kill Electrons

NEUTRON STAR

STRANGE STAR

Strange Quark Star



2-flavoured

3-flavoured

PARK of Gravity-CBM

the park of gravity-compressed baryonic matter





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To kill e^- 's: from Chandra to Landau

- A real historical note: Landau in 1931/1932!

ON THE THEORY OF STARS.

By L. Landau.

(Received 7 January 1932).

From the theoretical point of view the physical nature of Stellar equilibrium is considered.

The astrophysical methods usually applied in attacking the problems of stellar structure are characterised by making physical assumptions chosen only for the sake of mathematical convenience. By this is characterised, for instance, Mr. Milne's proof of the impossibility of a star consisting throughout of classical ideal gas; this proof rests on the assertion that, for arbitrary L and M , the fundamental equations of a star consisting of classical ideal gas admit, in general, no regular solution. Mr. Milne seems to have overlooked the fact, that this assertion results only from the assumption of opacity being constant throughout the star, which assumption is made only for mathematical purposes and has nothing to do with reality. Only in the case of this assumption the radius R disappears from the relation between L , M and R necessary for regularity of the solution. Any reasonable assumptions about the opacity would lead to a relation between L , M and R . It is quite exempt from the physical criteria of Eddington's mass-luminosity-relation.

It seems reasonable to try to attack the problems of stellar structure by methods of theoretical physics. To investigate the physical nature of stars for a given purpose we must at first investigate the physical nature of a given mass without generating a star, for which equilibrium being the condition (for given temperature). The physical nature of gravitation is negative and inversely proportional to some

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L. Landau

we have no need to suppose that the radiation of stars is due to some mysterious process of mutual annihilation of protons and electrons, which was never observed and has no special reason to occur in stars. Indeed we have always protons and electrons in atomic nuclei very close together, and they do not annihilate themselves; and it would be very strange if the high temperature did help, only because it does something in chemistry (chain reactions!). Following a beautiful idea of Prof. Niels Bohr's we are able to believe that the stellar radiation is due simply to a violation of the law of energy, which law, as Bohr has first pointed out, is no longer valid in the relativistic quantum theory, when the laws of ordinary quantum mechanics break down (as it is experimentally proved by continuous-rays-spectra and also made probable by theoretical considerations).¹ We expect that this must occur when the density of matter becomes so great that atomic nuclei come in close contact, forming one gigantic nucleus.

On these general lines we can try to develop a theory of stellar structure. The central region of the star must consist of a core of highly condensed matter, surrounded by matter in ordinary state. If the transition between these two states were a continuous one, a mass $M < M_0$ would never

state (i. e. without ... Because, as far as we can conclude that the transition is separated by some distance and its vapour is caused by some kind of ... the existence of ... the two states.

In the above con-

siderations is yet to be constructed, and only such a theory can show how far they are true.

February 1931, Zurich.

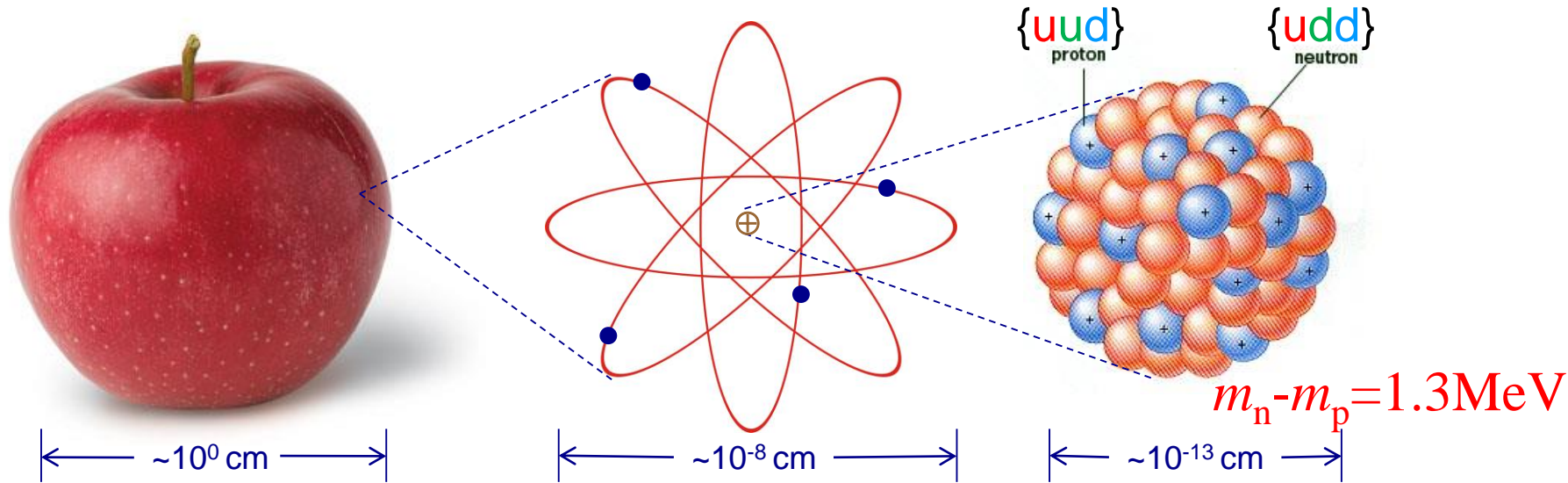
¹ L. Landau and R. Peierls, ZS. f. Phys. 69, 56, 1931.

$$A_c \approx \lambda_c^3 / \text{fm}^3 \sim 10^9$$
$$\text{Compton } \lambda_c \sim 2.4 \times 10^3 \text{ fm}$$

Landau L. 1932, *Sov. Phys.*, 1, 285

To kill e 's: from Chandra to Landau

- Let's do an **exercise**...to squeeze an apple!



Total baryon number $A_{\text{apple}} \sim 100\text{g}/u \sim 10^{26}$.

Electrons no-relativistic before squeezing, but what after?

A giant “**nucleus**”: $\sim 50\mu\text{m}$, $\sim \rho_{\text{nucl}}$, $E_e \sim 200\text{MeV}$ if e keeps

Gravity-squeezed core: $A \gg A_{\text{apple}}$ and *not gravity-free!*

Two **uncertainties**: Quarks de-*confined*? *Strangeness*?



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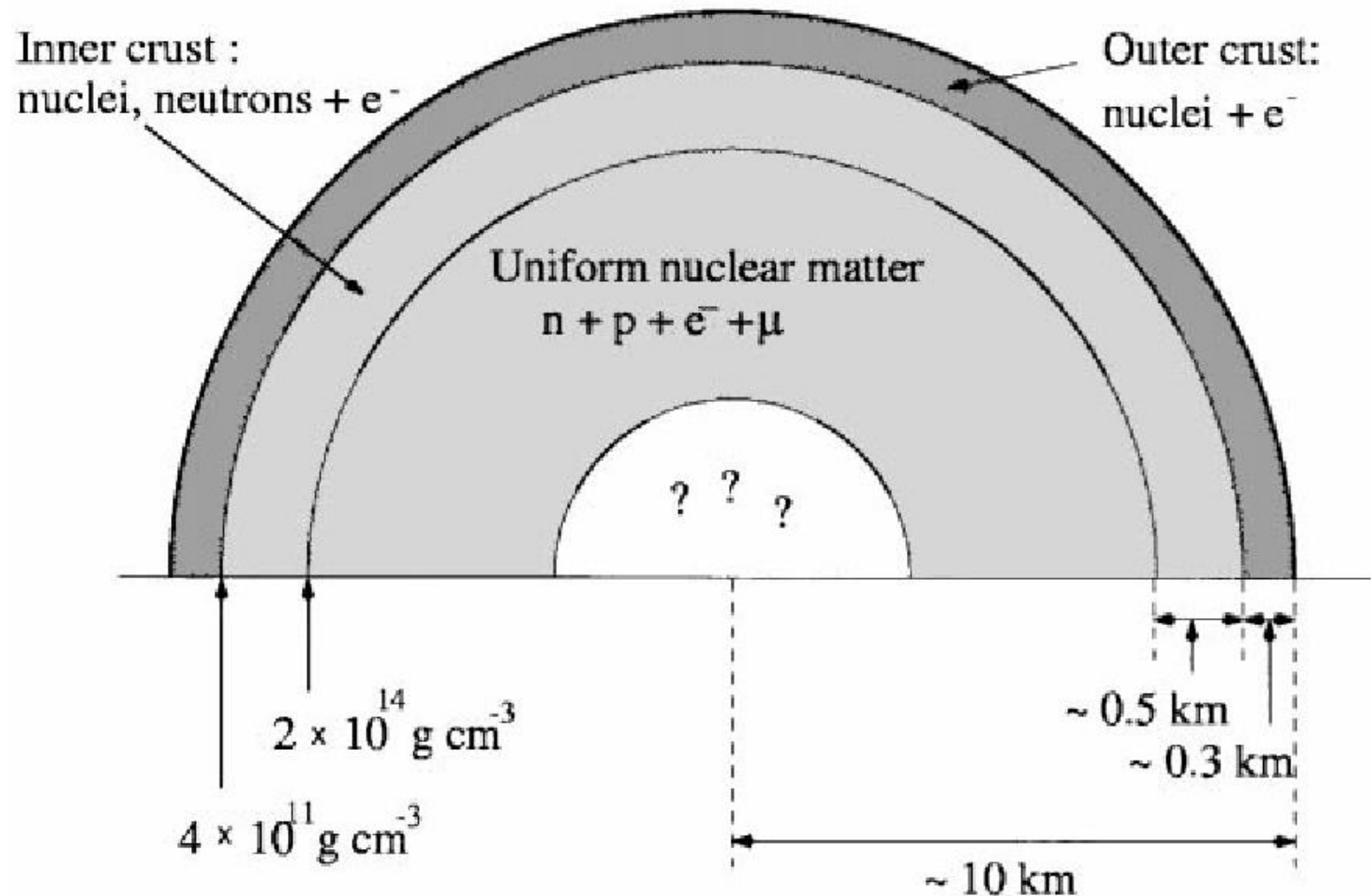
PARK of Gravity-CBM

the park of gravity-compressed baryonic matter



Neutron Stars

- **Neutronization**: $p + e \rightarrow n + \nu_e$ ($u + e \rightarrow d + \nu_e$) if **2-favoured**





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Why strange?

Hobby

- God may love a matter state with flavor-maximization...

For strong matter around the *nuclear density*, the separation between quarks, $\Delta\ell$, could be ~ 0.5 fm, determined by α_s !

From Heisenberg's uncertainty relation, $\Delta\ell \cdot \Delta p \approx \hbar$, one may have an energy scale for strong matter, E_{scale} , ~~(c,b,t)~~

$$E_{\text{scale}} \approx \hbar c / \Delta\ell \approx 0.2 \text{ GeV} \cdot \text{fm} / 0.5 \text{ fm} = 0.4 \text{ GeV.}$$

Note that... we may expect 3-flavored strong matter because

$$E_{\text{scale}} \gg \Delta m_{\text{uds}} \equiv (m_s - m_{\text{ud}})c^2!$$

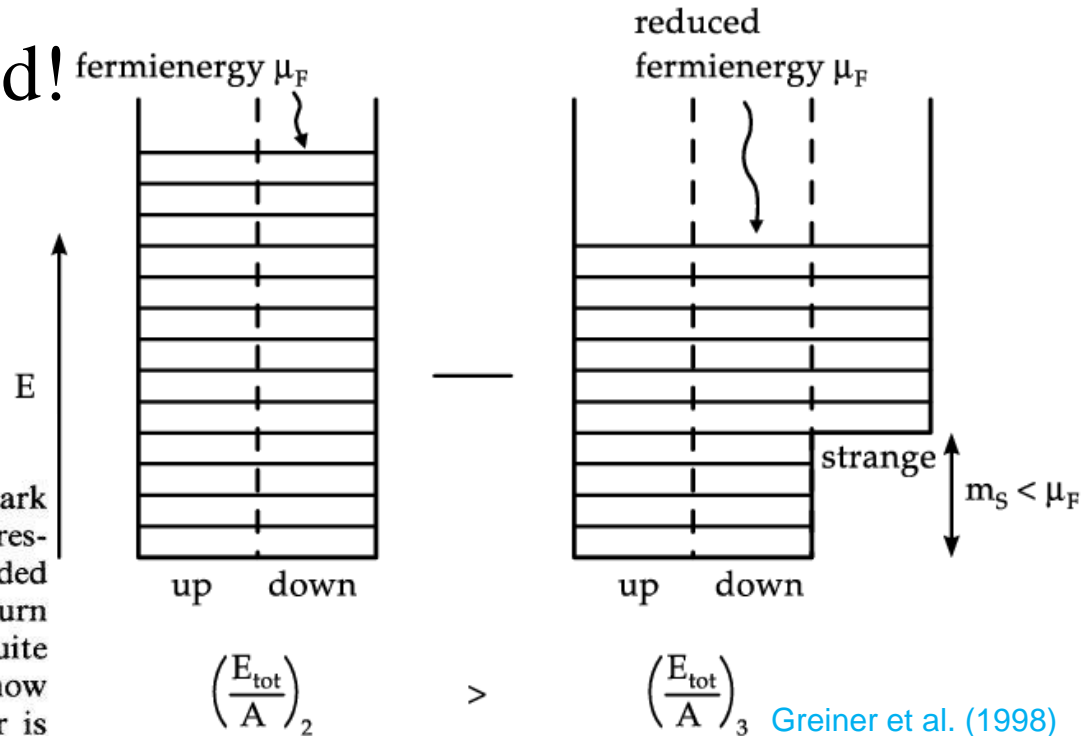
- Strong matter should be **3**-flavored (*u,d,s*), but why are normal stable atomic nuclei **2**-flavored (*u,d*)? ← Nuclear Symmetry Energy

Strange Quark Stars

- **Free** quarks 3-flavoured! As noted in Witten (1984), SQM matter in bulk may constitute the true ground state.
- A. Stable macroscopic quark matter

The most extreme possibility is that macroscopic quark matter is bound and stable at zero temperature and pressure. At first sight, one might think that this is excluded by the fact that ordinary nuclei do not spontaneously turn into the hypothetical dense quark state. This is not quite so, however. Observations of nuclear physics only show that in the absence of strange quarks, nuclear matter is more stable than quark matter. Addition of strangeness does not help stabilize nuclear matter, because strange baryons are heavier than nonstrange baryons. For quark matter, the story is different. (This point has been noted before in Refs. 19 and 26.) The likely Fermi momentum in quark matter is 300–350 MeV, more than the strange-quark mass, so it is energetically favored for some of the nonstrange quarks to become strange quarks, lowering the Fermi momentum and the energy.

This effect can easily be estimated in the simplest form of the bag model.²¹ A single quark flavor of Fermi



$$\left\{ \begin{array}{l} m_u = 2 \sim 8 \text{ MeV} \\ m_d = 5 \sim 15 \text{ MeV} \\ m_s \sim 100 \text{ MeV} \end{array} \right.$$

$$\rho = 2\rho_N$$

↓

$$\mu_F \sim 500 \text{ MeV}$$



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Strangeon Stars

- Can quarks be *free* at a few $\rho_{\text{nucl}}???$

The state of cold quark matter:
a model-independent view

Renxin Xu (徐仁新)

School of Physics, Peking University

Compact stars in the QCD phase diagram
(CSQCD II), PKU

May 24th, 2009.

Stiff EoS \Rightarrow low ρ_c

Quark cluster = Strangeon

- What *if* strong *color*-interaction exists?
Strong interaction \Rightarrow *quark cluster*?

Diquark: color SU(3), *Coulomb*-like
 $3 \times 3 = 6(\text{repulsion}) + 3^*(\text{attraction})$

Let's estimate the length scale of and
interaction strength in a quark cluster:

$$l_q \sim \frac{1}{\alpha_s} \frac{\hbar c}{mc^2} \sim 1\text{fm} / \alpha_s, \quad E_q \sim 300\alpha_s^2 \text{MeV},$$

if quarks are dressed, with mass $\sim 300\text{MeV}$.

What about α_s ?

strangeon ['streɪdʒɪɒn] = **strange** + nucle**on** with strangeness $S = -B$

Strangeon Stars

- A linked bag-model of strangeon matter...

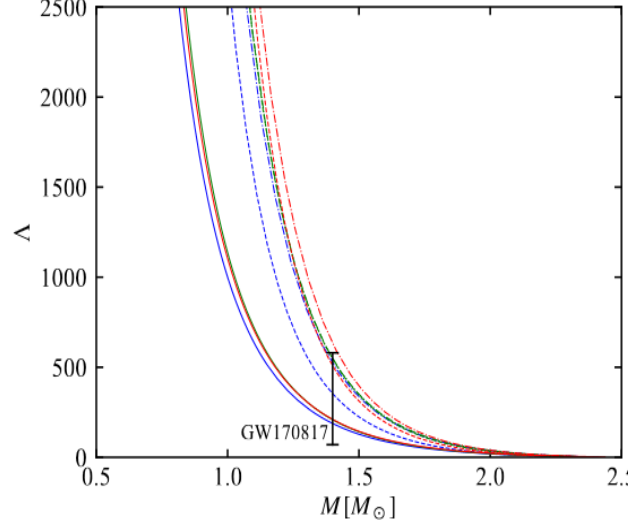
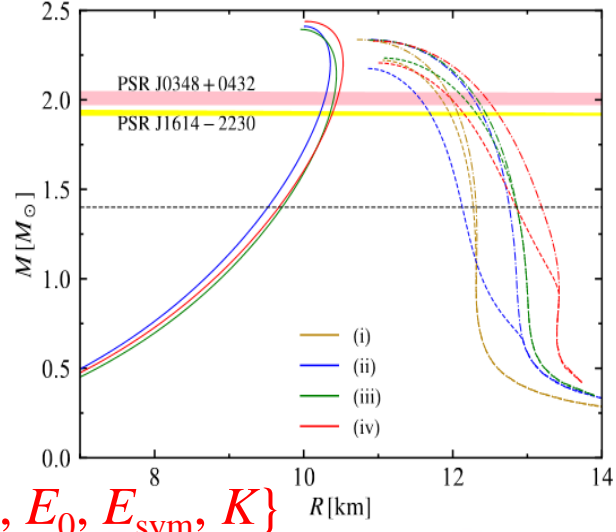
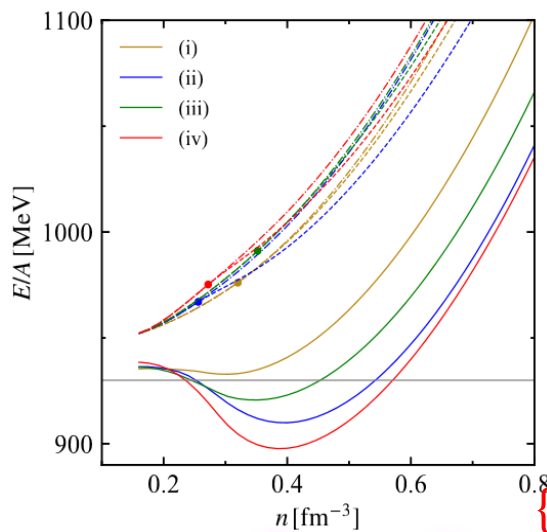
In the Fermi-gas approximation, the energy per lattice cell is obtained with

Miao et al. (arXiv: 2008.06932)

$$E = \sum_j (\Omega_j + N_j \mu_j) + BV - \frac{z_0}{r_{\text{bag}}} \frac{\omega}{4\pi}, \quad (1)$$

$$\Omega_i = \Omega_{i,V}V + \Omega_{i,S}S + \Omega_{i,C}C, \quad (2)$$

$$B = B_0 + B_1\xi + B_2\xi^2 + B_3\xi^3, \quad (8)$$

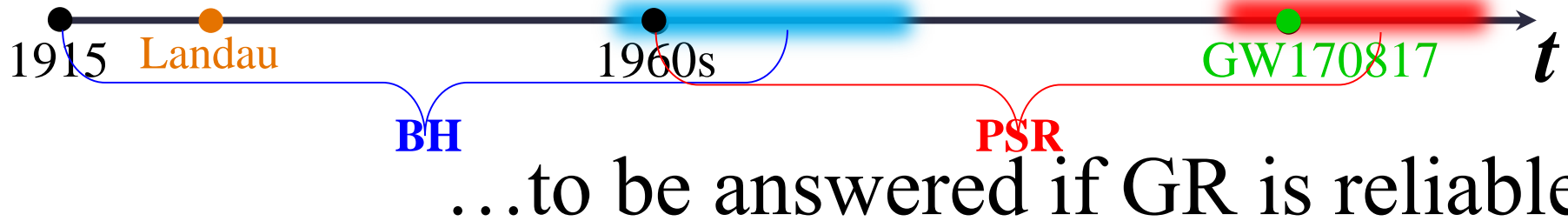


$\{n_0, E_0, E_{\text{sym}}, K\}$

	C_1	$\hat{m}_s [\text{MeV}]$	$B_2 [\text{MeV}/\text{fm}^3]$	$B_3 [\text{MeV}/\text{fm}^3]$	$z_0(n_0)$	$L [\text{MeV}]$
(i)	2.7	220	136.7	50	2.944	45.1
(ii)	2.7	220	112.7	100	2.926	52.7
(iii)	2.7	280	125.0	100	2.908	56.6
(iv)	3.2	280	162.3	100	2.843	62.8

Summary

- **BH** astrophysics was active, but it is a golden era of **NS/PSR** with multi-messenger astron.:



...to be answered if GR is reliable!

- The basic units inside pulsar-like stars could be **3-flavour *symmetric strangeons*** rather than **2-flavour *asymmetric nucleons*** if the Nature really loves symmetry when building the world.
- To test strangeon model further... **THANKS!**