



# Strong magnetic field impacts on the neutrino transport in Core-Collapse Supernovae

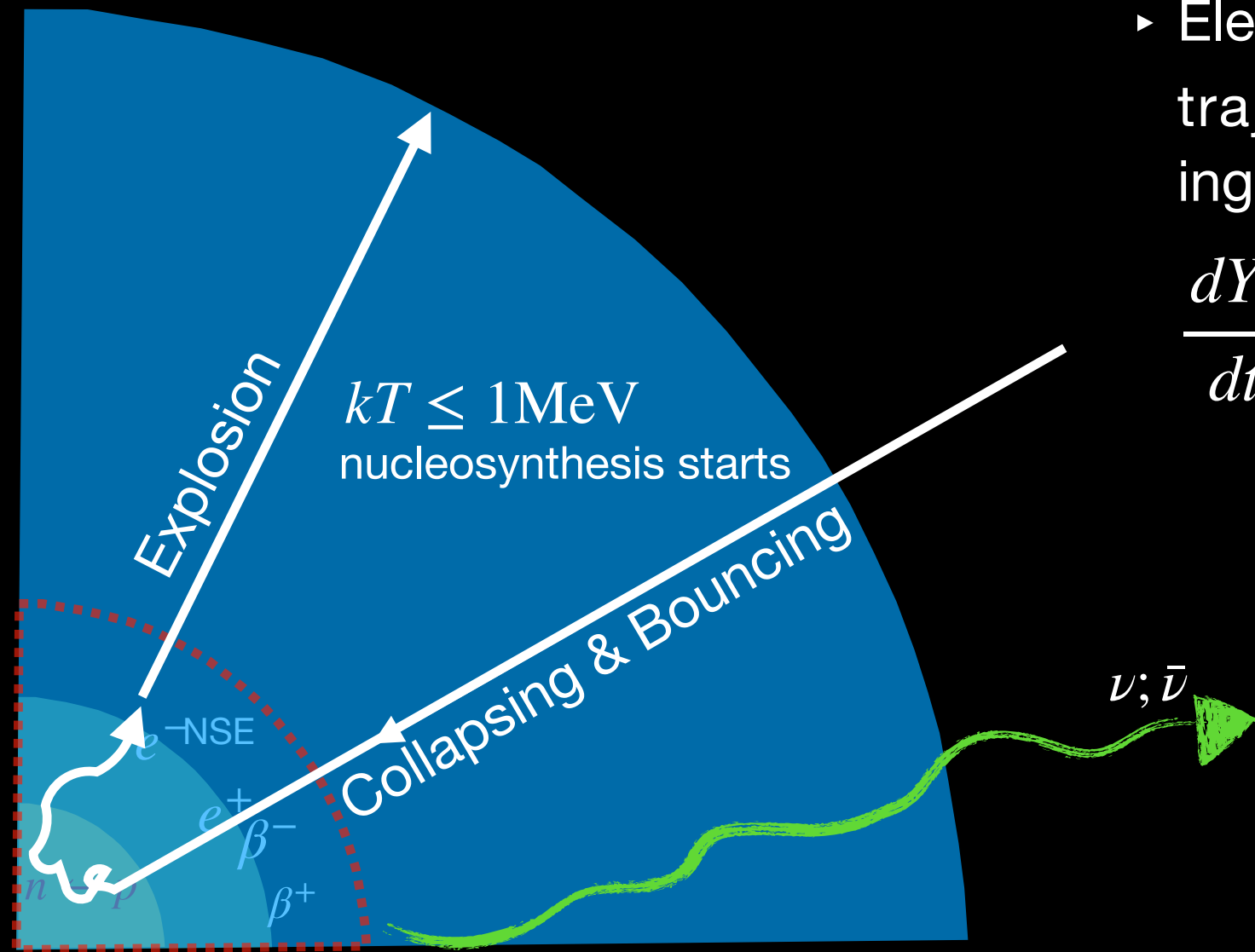
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Peking Univ.

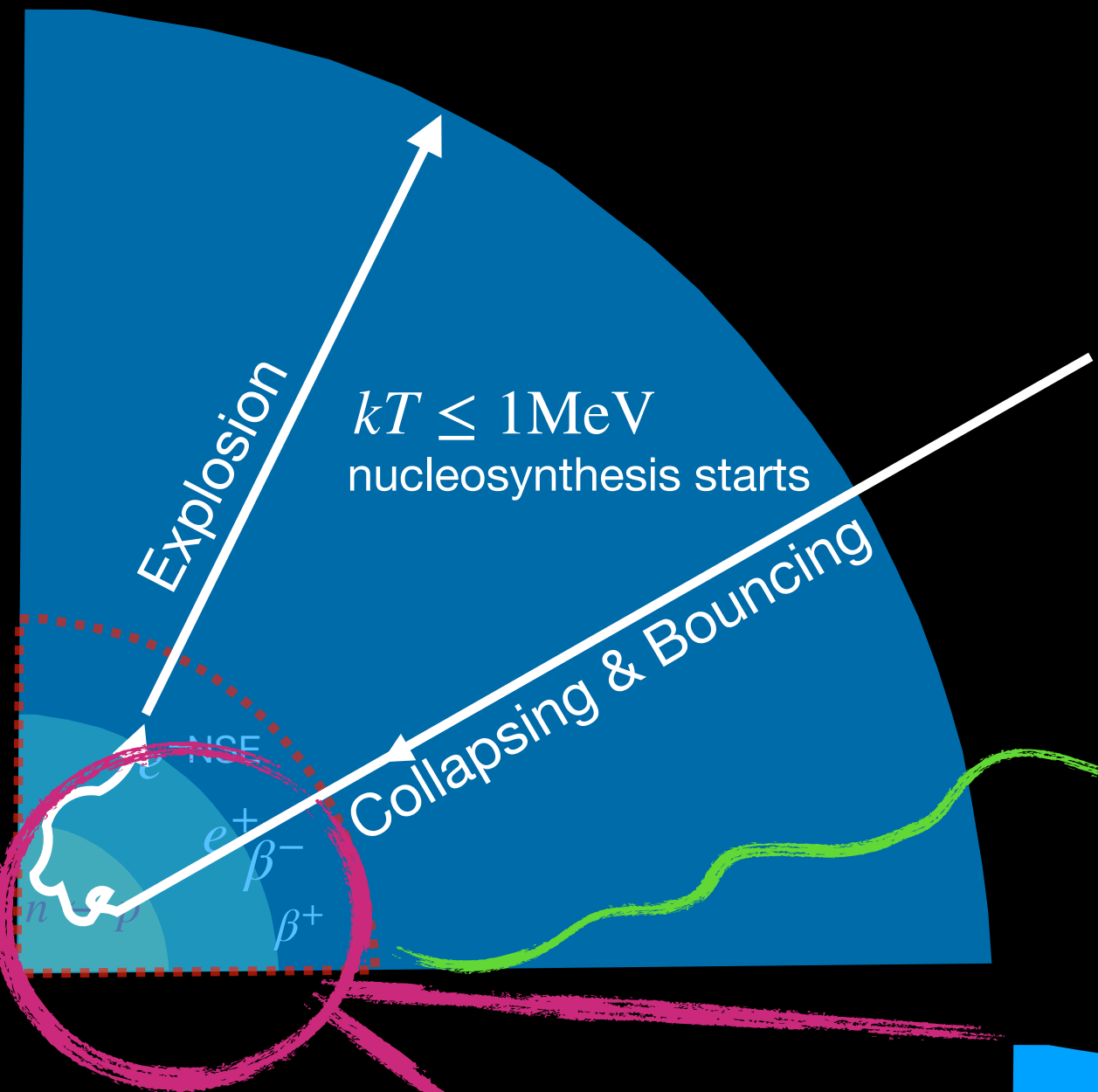
Collaborators: Prof. Shuai Zha (Yunnan Observatories, CAS )  
Prof. Toshitaka Kajino (Beihang Univ. )  
Submitted to ApJ

Dialogue at the Dream Field@Guizhou, 13th May



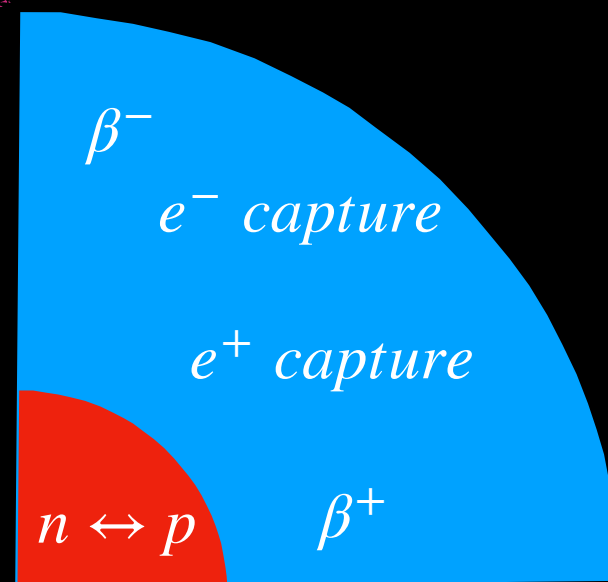
- ▶ Electron fraction  $Y_e$  evolves along with the trajectory.  $Y_e(T, \rho, Y_e)$  describes the ingredient for r-process nucleosynthesis

$$\frac{dY_e}{dt} = -(\lambda_{pe^-} + \lambda_{p\bar{\nu}_e})X_p + (\lambda_{ne^+} + \lambda_{n\nu_e})X_n + \sum_h \left(\frac{X_h}{A_h}\right) (\lambda_{h\nu_e} + \lambda_{he^+} - \lambda_{h\bar{\nu}_e} - \lambda_{he^-})$$



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- ▶  $\beta$  decay and  $e^\pm$  capture: determining the abundance flow and the isotopic ratio
- ▶  $n \leftrightarrow p$  : determines electron fraction  $Y_e$ , which further affect the neutron-richness.

Magnetic field in SNe could be  $10^{14\sim 16}$  G, strongly influence the electron motion as well as weak interactions

## ▶ Electron capture rate with magnetic field

$$\Gamma_{pe^- \rightarrow n\nu_e}^B = \sum_{n=0}^{\infty} (2 - \delta_{n0}) \cdot \int_0^{\infty} \sigma(E_\nu, B) dp_z f_{FD}(\epsilon; \mu, T_\nu) g(E_\nu; \mu_\nu, T_\nu)$$

$$E_e^2 = p_z^2 + m_e^2 + 2eBn$$

( $c = \hbar = 1$ )

Phase space:

$$\sum_{n=0}^{\infty} (2 - \delta_{n0}) \frac{dp_z}{2\pi} \frac{eB}{2\pi} f_{FD}(E_B, T)$$

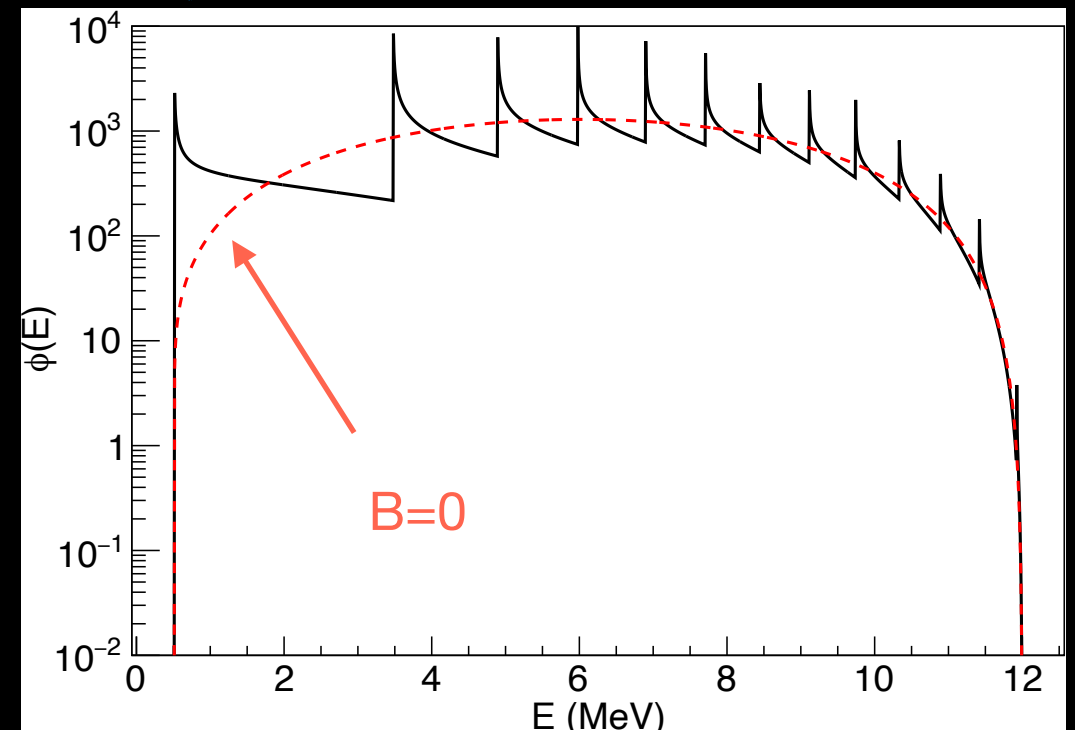
### Cross section of neutrino interaction

$$\sigma_{\nu N}(B) = \sigma_B^1 \left[ 1 + 2\chi \frac{(f \pm g)g}{f^2 + 3g^2} \cos \Theta_\nu \right] + \sigma_B^2 \left[ \frac{f^2 - g^2}{f^2 + 3g^2} \cos \Theta_\nu + 2\chi \frac{(f \mp g)g}{f^2 + 3g^2} \right]$$

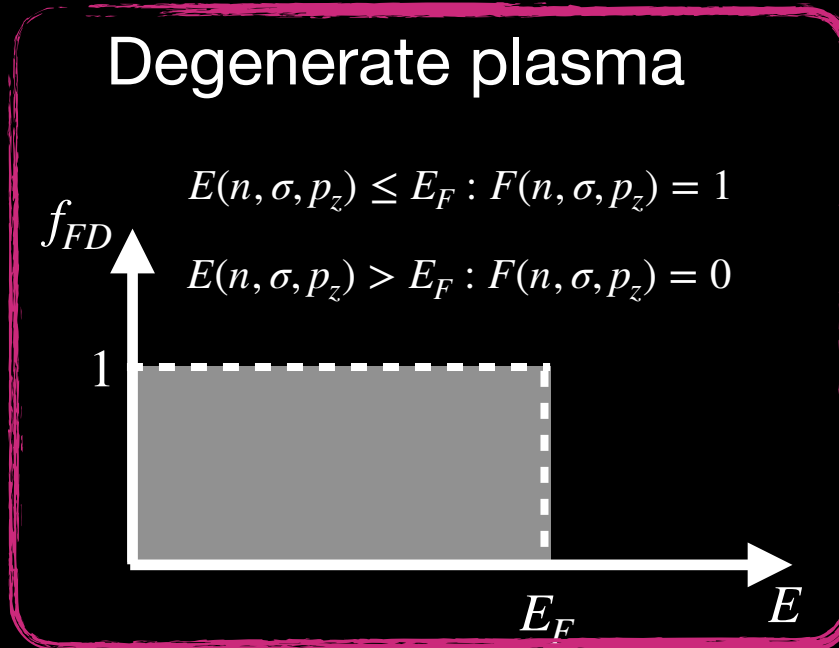
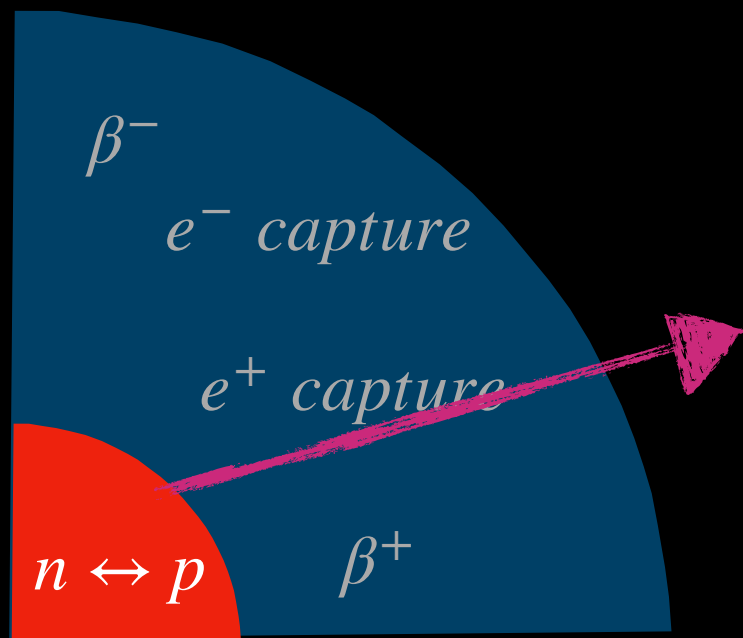
$$\sigma_B^1 = \frac{G_F^2 \cos^2 \theta_C}{2\pi} (f^2 + 3g^2) eB \sum_{n=0}^{n_{max}} \frac{g_n E_e}{\sqrt{E_e^2 - m_e^2 - 2neB}}$$

$$\sigma_B^2 = \frac{G_F^2 \cos^2 \theta_C}{2\pi} (f^2 + 3g^2) eB \frac{E_e}{\sqrt{E_e^2 - m_e^2}}$$

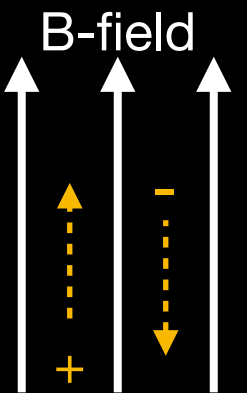
Duan&Qian 2005



Famiano et al, ApJ 898, 163



► Fermi energy within magnetic field

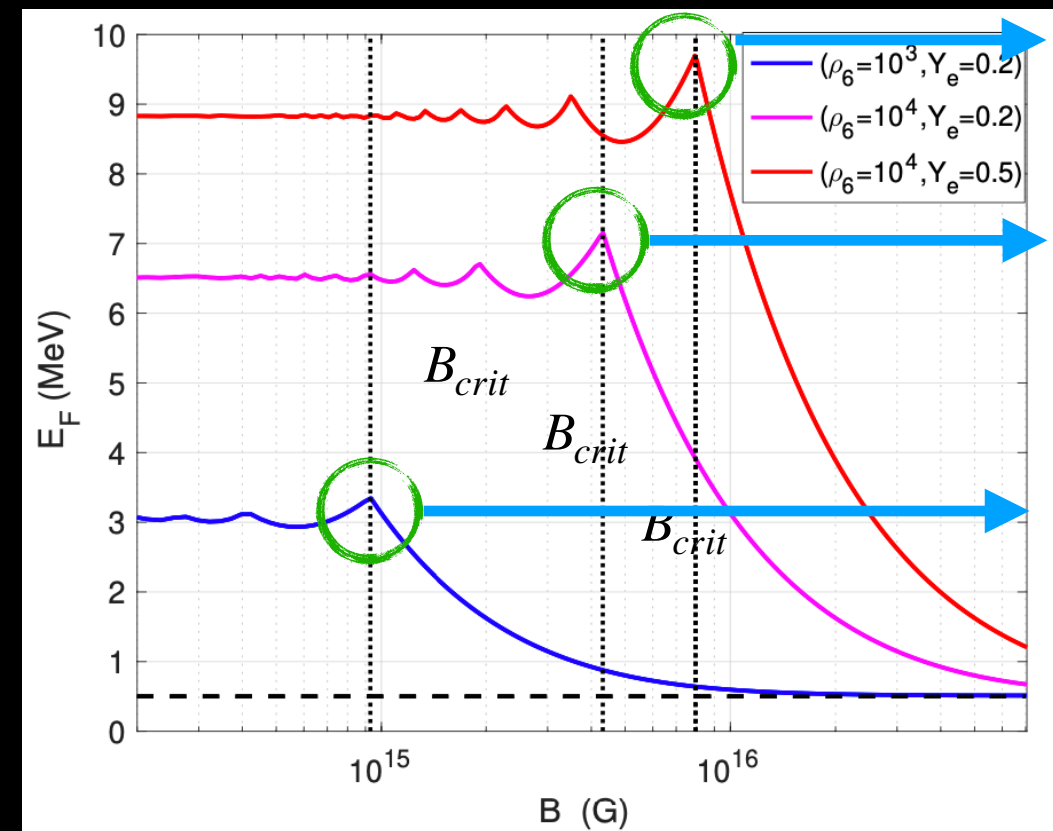


Total number of electrons  
= Summation of electrons on all levels

$$\rho Y_e = \frac{N_+}{V} + \frac{N_-}{V}$$

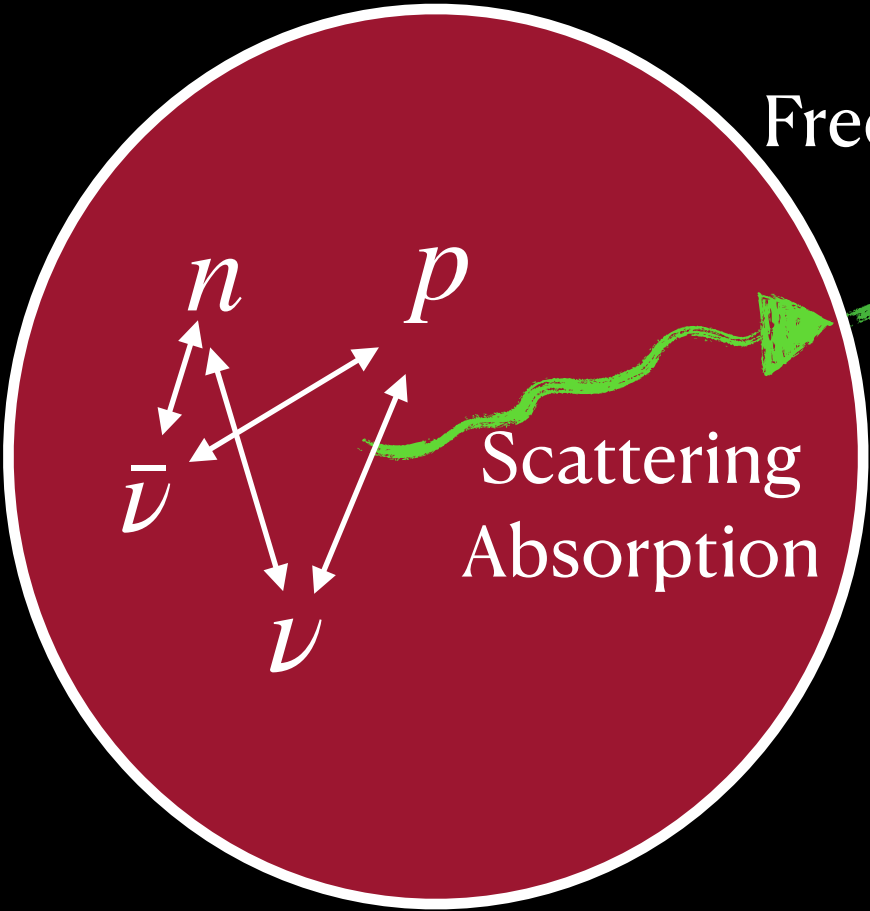
- $B > B_{crit}$ , only lowest Landau level occupied,  $E_F$  decreases monotonically as a function of  $B$
- The new Landau level leads to a peak of  $E_F$ , results in a wiggle shape

$E_F$  as a function of B-field strength



$$(\rho_6 = \rho/10^6)$$

$$(B_{crit} = \frac{\pi}{e} [2\pi(\rho Y_e)^2]^{1/3})$$

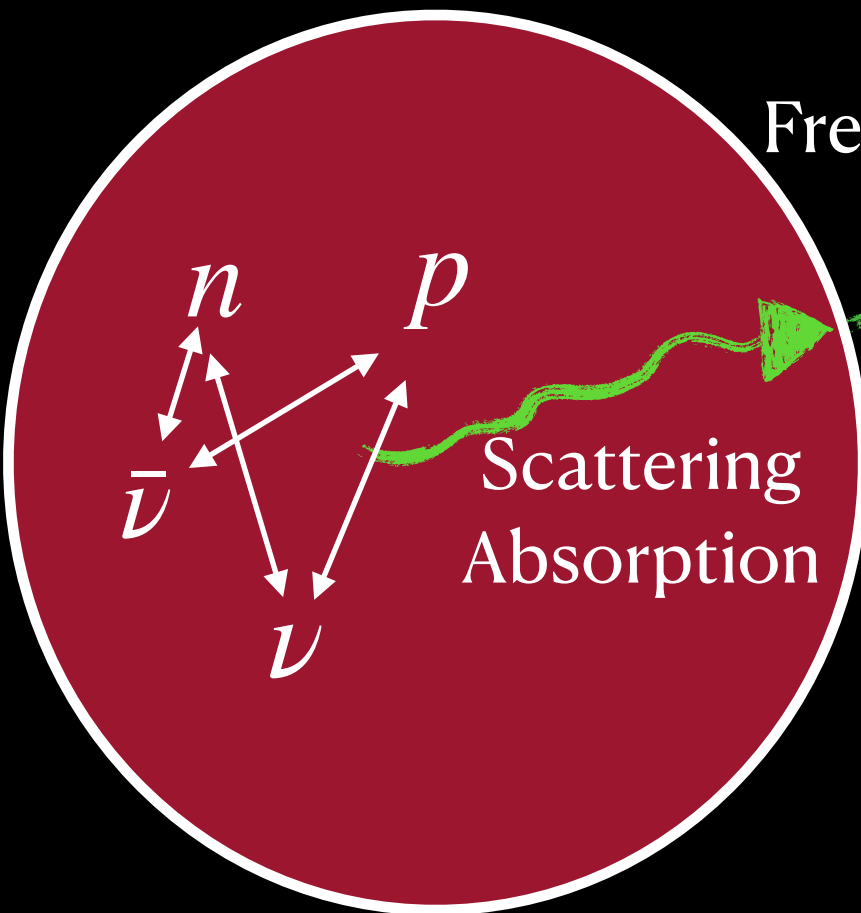


► Definition of  $\nu$  - sphere

Optical Depth along the path  $[s_1, s_2]$ :

$$\tau_{\nu_i}([s_1, s_2]) = \int_{s_1}^{s_2} ds \kappa_{\nu_i}(s)$$

$$\tau = \int_{R_\nu}^{\infty} dr \kappa_t(r) = \frac{2}{3}$$



Free-streaming

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Opacity

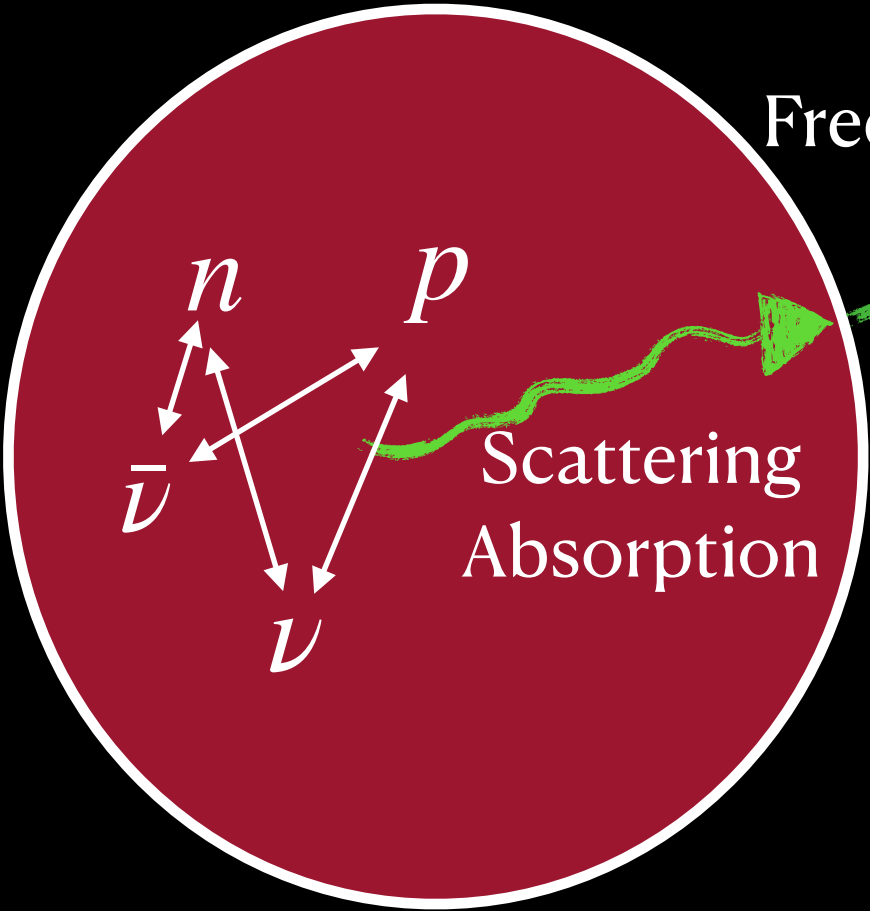
$$\kappa_t(\nu_e) = \kappa_s(\nu_e n) + \kappa_s(\nu_e p) + \kappa_a(\nu_e n)$$

$$\kappa_t(\bar{\nu}_e) = \kappa_s(\bar{\nu}_e n) + \kappa_s(\bar{\nu}_e p) + \kappa_a(\bar{\nu}_e p)$$

$$\kappa_t(\nu_x) = \kappa_s(\nu_x n) + \kappa_s(\nu_x p)$$

s: scattering on n&p,  
a: absorption on n/p

M. Ruffert, H.-Th. Janka, and G. Schafer (1995)  
S. Rosswog and M. Liebendorfer (2003)  
A. Perego, R. M. Cabezón, and R. Käppeli (2016)



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Opacity

$$\begin{aligned} \kappa_t(\nu_e) &= \kappa_s(\nu_e n) + \kappa_s(\nu_e p) + \kappa_a(\nu_e n) \\ \kappa_t(\bar{\nu}_e) &= \kappa_s(\bar{\nu}_e n) + \kappa_s(\bar{\nu}_e p) + \kappa_a(\bar{\nu}_e p) \\ \kappa_t(\nu_x) &= \kappa_s(\nu_x n) + \kappa_s(\nu_x p) \end{aligned}$$

s: scattering on n&p,  
a: absorption on n/p

change with B-field

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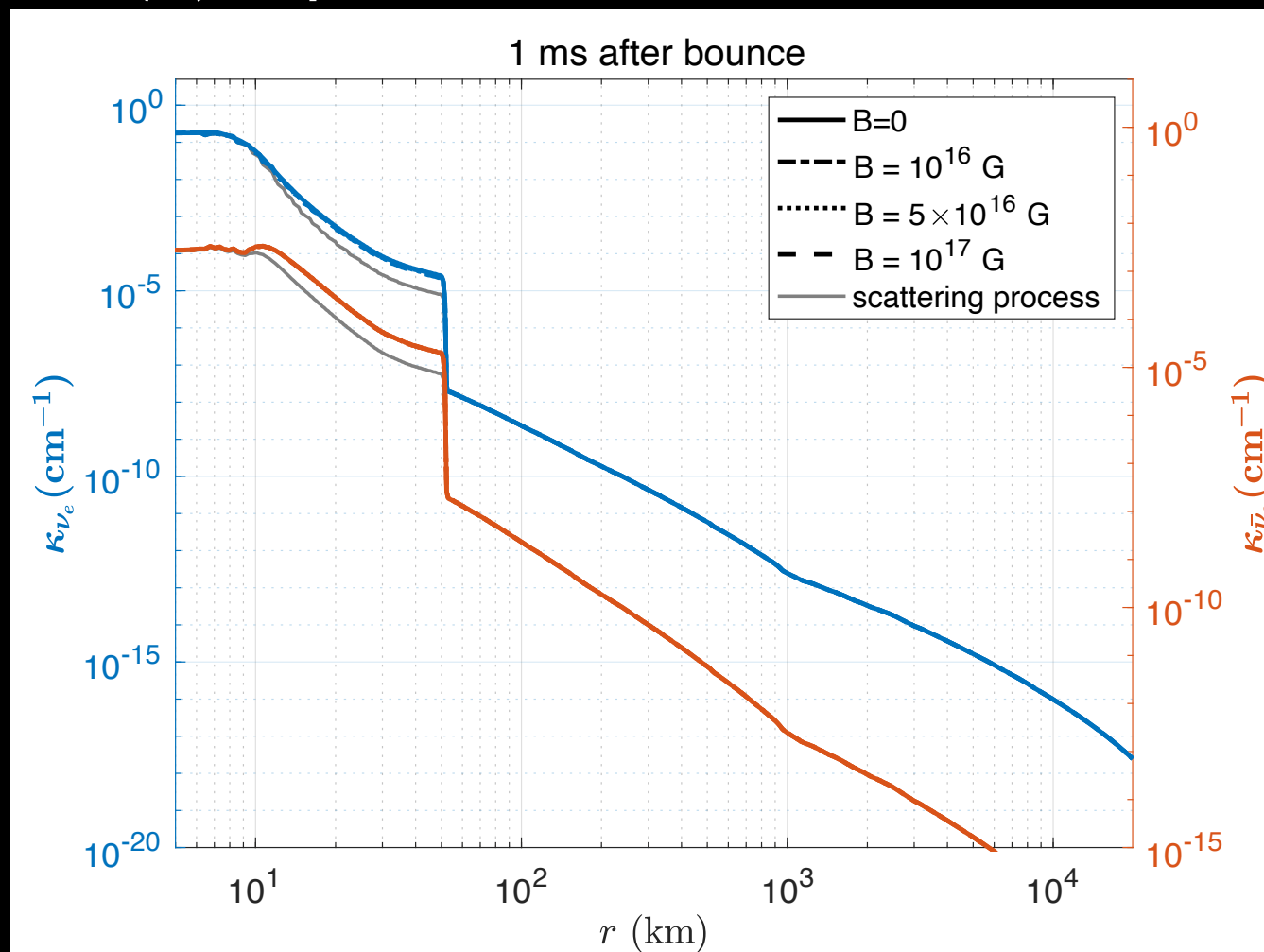
$$\kappa_a^B(\nu_e n) = A \rho Y_{np} \left( \frac{1}{m_e c^2} \right)^2 \frac{\int_0^\infty \sigma_{\nu N}(E_e, B) E_\nu^4 f_{FD}(E_\nu, \mu_\nu; T_\nu)}{\int_0^\infty E_\nu^2 f_{FD}(E_\nu, \mu_\nu; T_\nu)} \left[ 1 - \frac{1}{\exp(F_5(\eta_{\nu_e})/F_4(\eta_{\nu_e}) - \eta_e)} \right]$$



GR1D: 1D Core-Collapse SNe code O'Connor & Ott375 2010; O'Connor 2015

EoS: Lattimer & Swesty LS180 (1991) Progenitor: 9.6  $M_{\odot}$  massive star (Heger)

## ► $\nu(\bar{\nu})$ Opacities

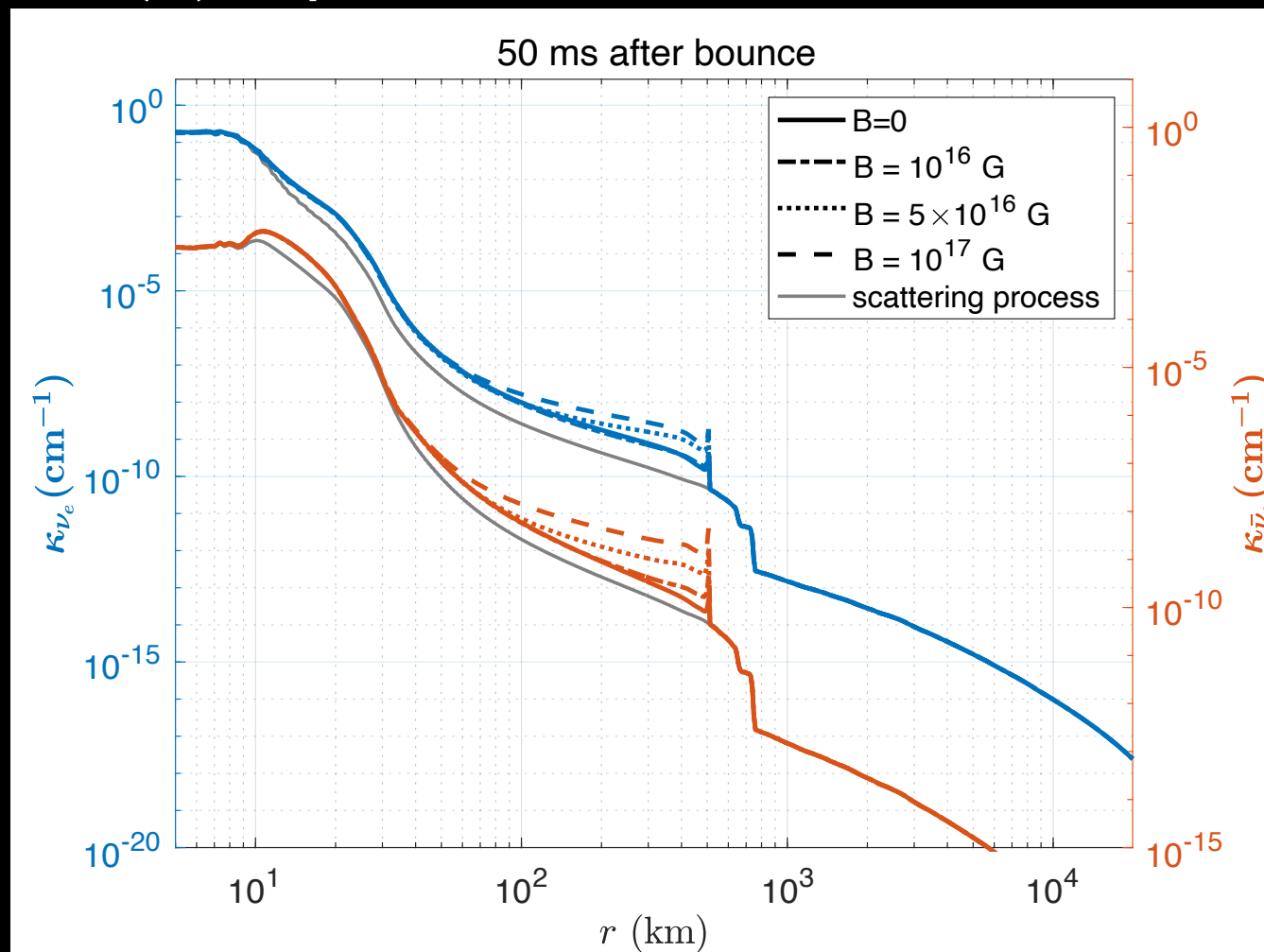


Y. Luo et al submitted

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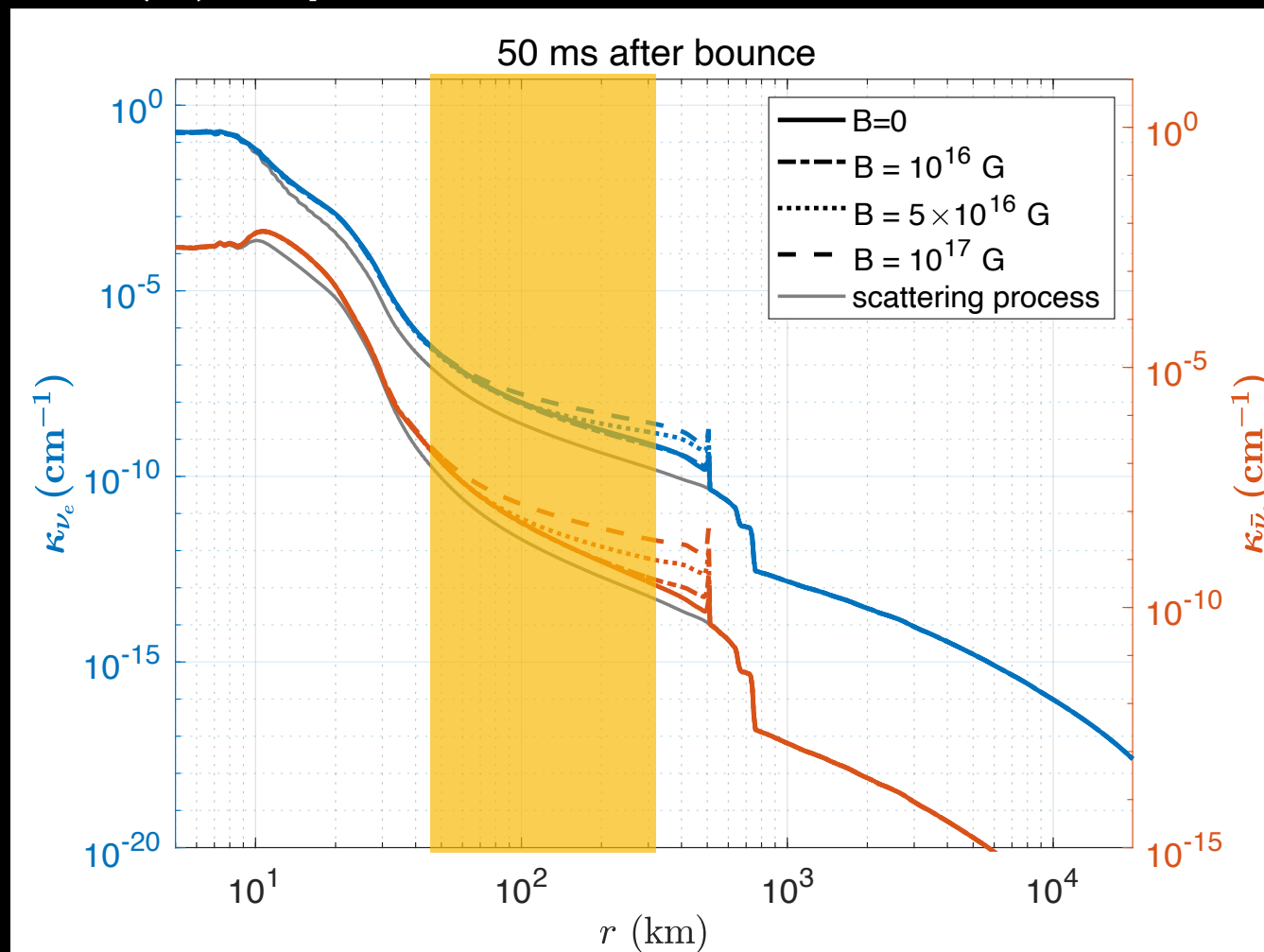


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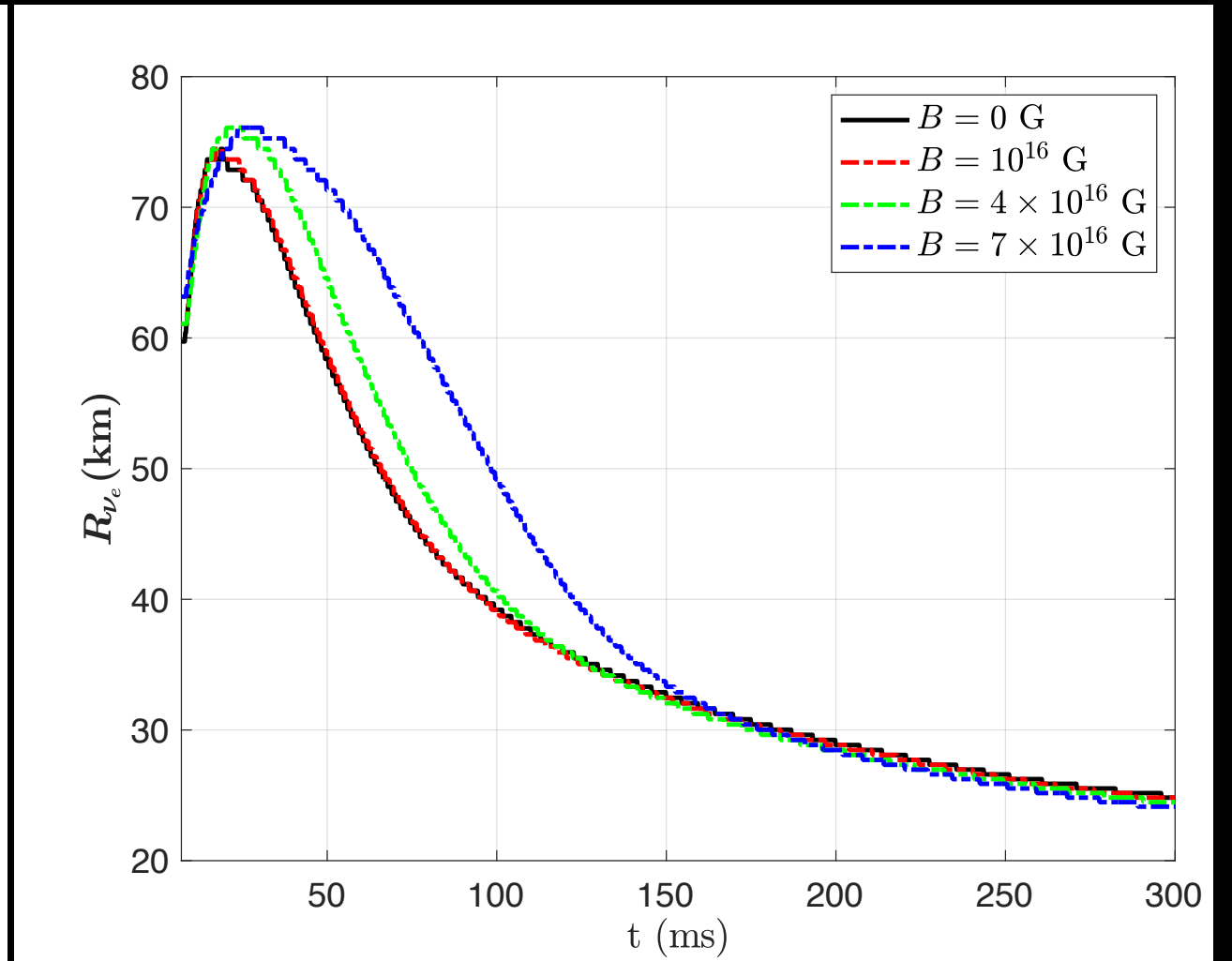
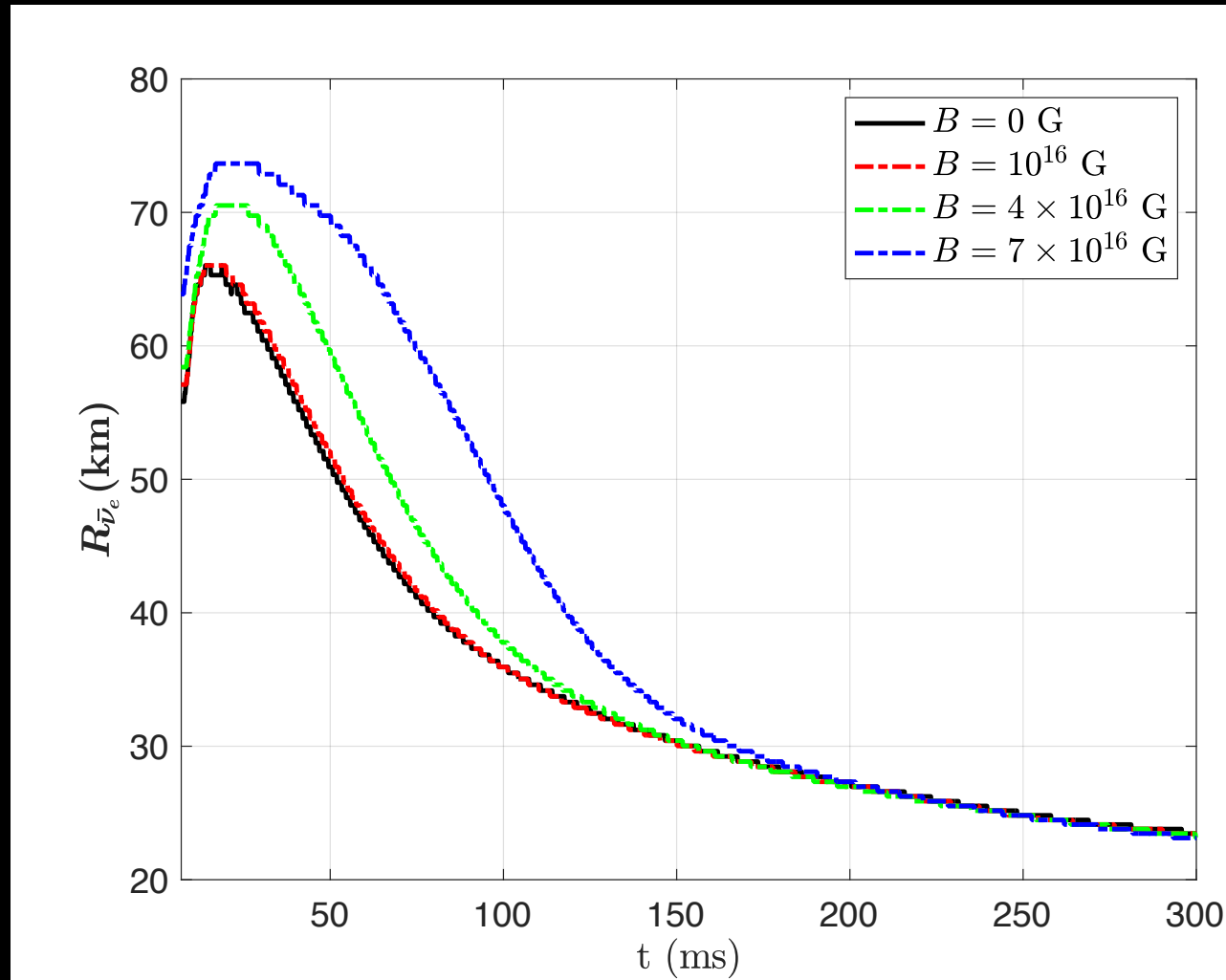
Y. Luo et al submitted

- No significant change @High density&temperature region (B is not strong to make  $e^{\pm}$  confine on LLL)
- Quantized phase space of  $e^{\pm}$ 
  - Enhancement of the number density
  - Enhancement of the interaction rate

## ► $\nu(\bar{\nu})$ Spheres evolution

Y. Luo et al submitted

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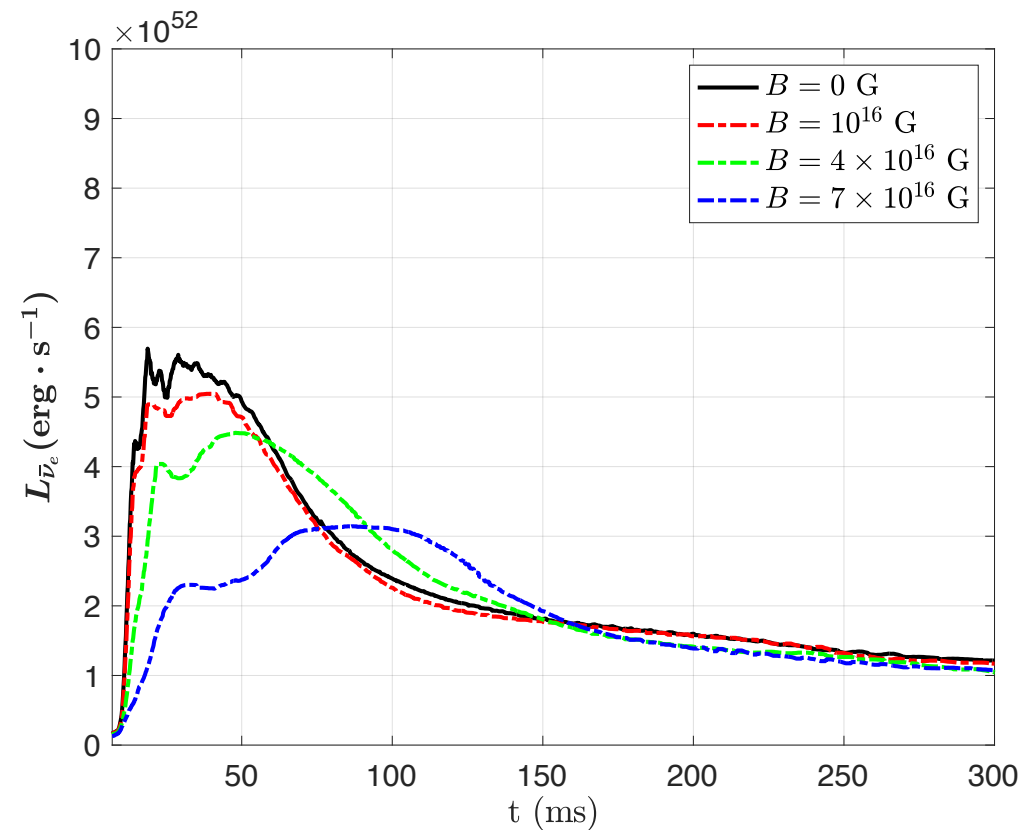
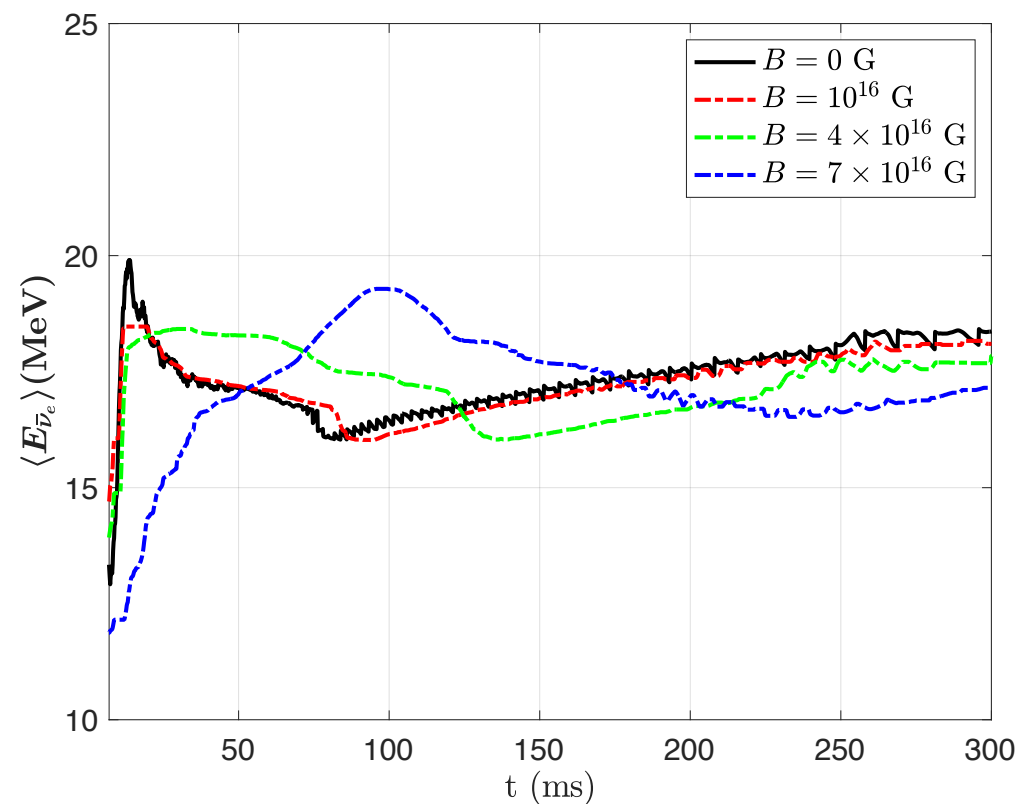
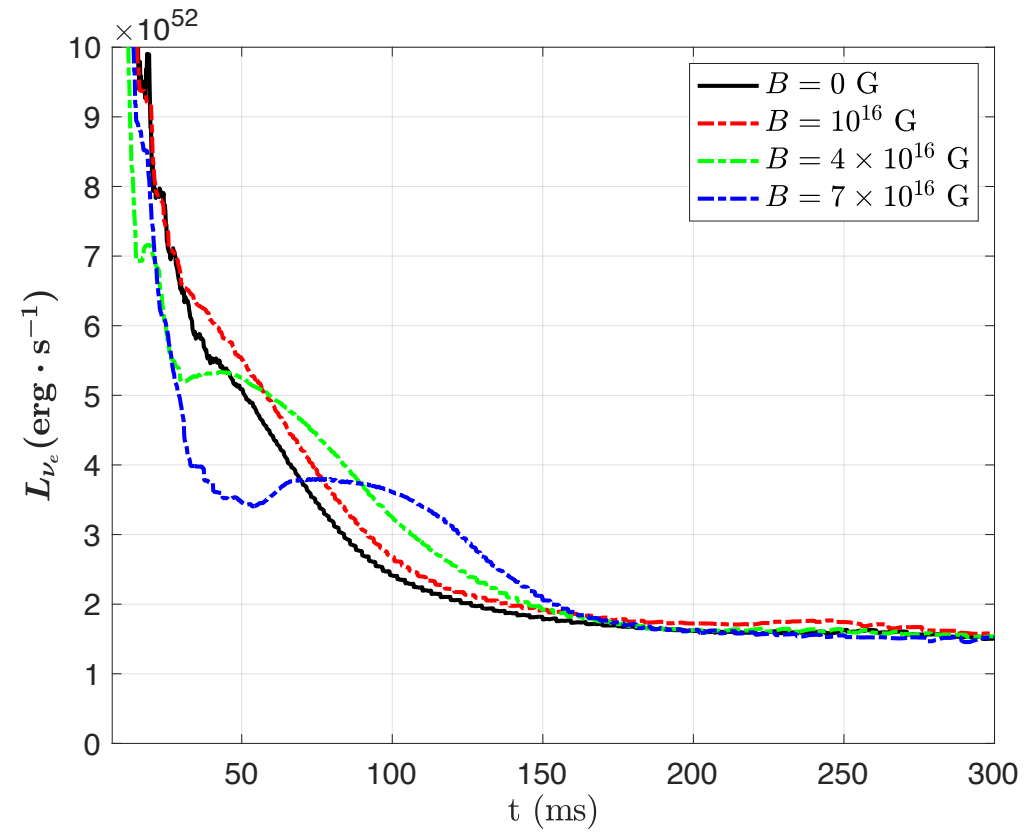
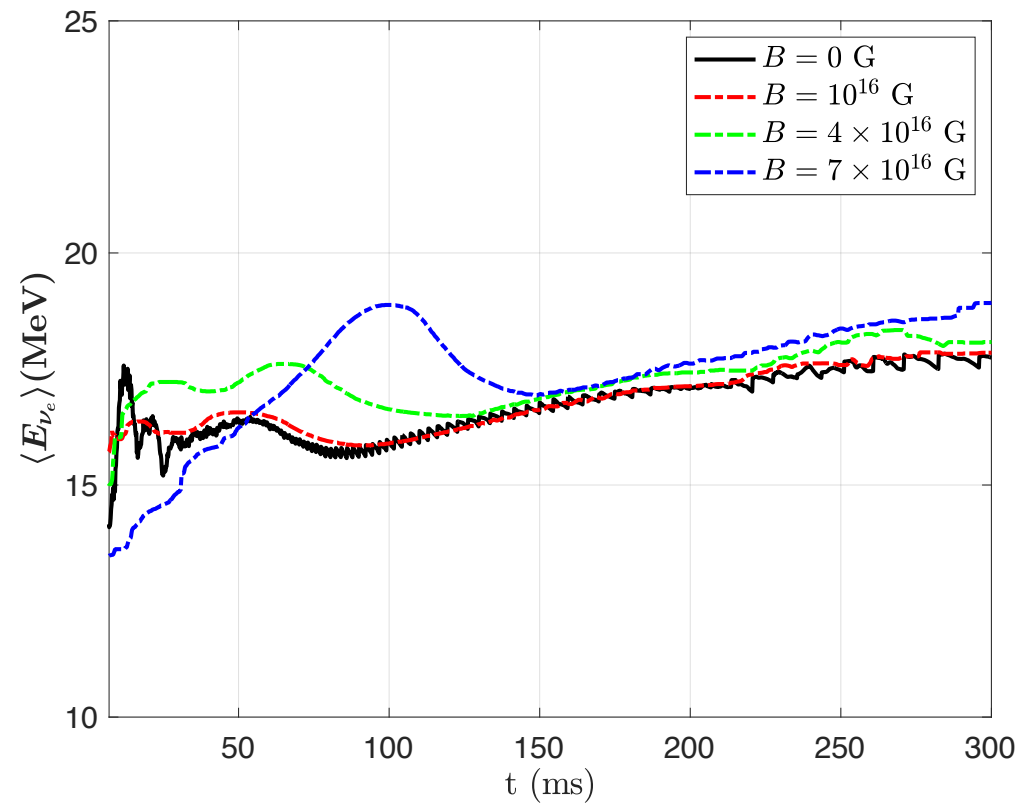


- The leakage scheme with B-field is modified in GR1D
- B-field is set as a const, but with  $r_{\text{cut}} = 100$  km

- Enhanced  $\nu(\bar{\nu})$  spheres after bounce until 150 ms
- Enhanced opacities directly enlarge  $\nu(\bar{\nu})$  spheres
- $\bar{\nu}$ -sphere is more sensitive

## ▶ $\nu(\bar{\nu})$ Mean energy

## ▶ $\nu(\bar{\nu})$ Luminosity



**Dream Field**

# Why Magnetic field

Mo is a valuable element to study all nucleosynthetic processes in the solar-system.

Isotopic fraction in %

$^{92}\text{Mo}$	$^{94}\text{Mo}$	$^{95}\text{Mo}$	$^{96}\text{Mo}$	$^{97}\text{Mo}$	$^{98}\text{Mo}$	$^{100}\text{Mo}$
14.53	9.15	15.84	16.67	9.6	24.39	9.82
100	99.1	0	0	0	0	0
0	0.9	69.6	100 (over)	63.7	82.2	4.5
0	0	30.4	0	36.3	17.8	95.5

Bisterzo et al., MNRAS  
418, 284 (2011); ApJ.  
787, 10 (2014).

**Red = p-process**

**Green = s-process**

**Blue = r-process**

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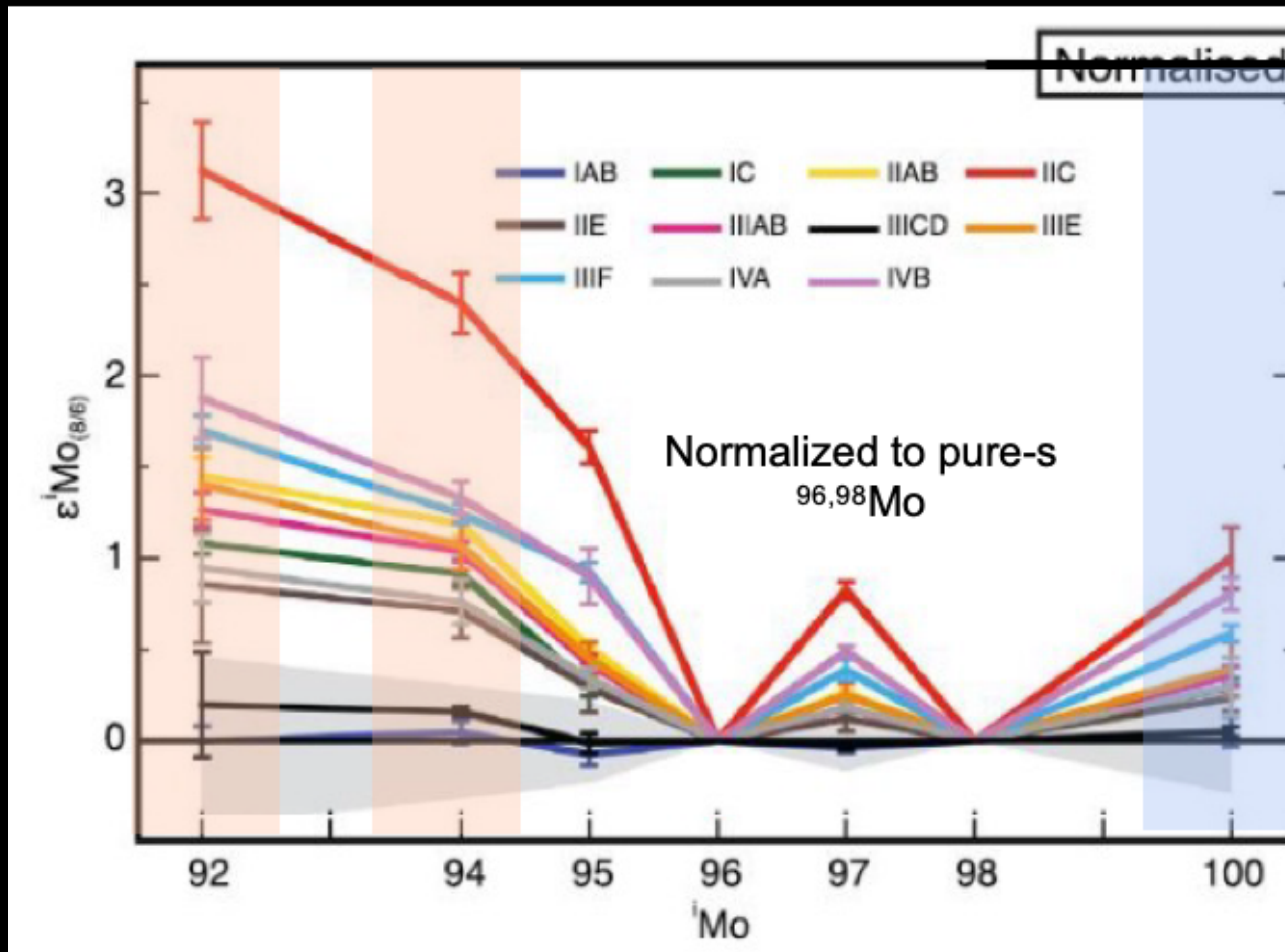
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Poole, EPSL, 473, 215 (2017).

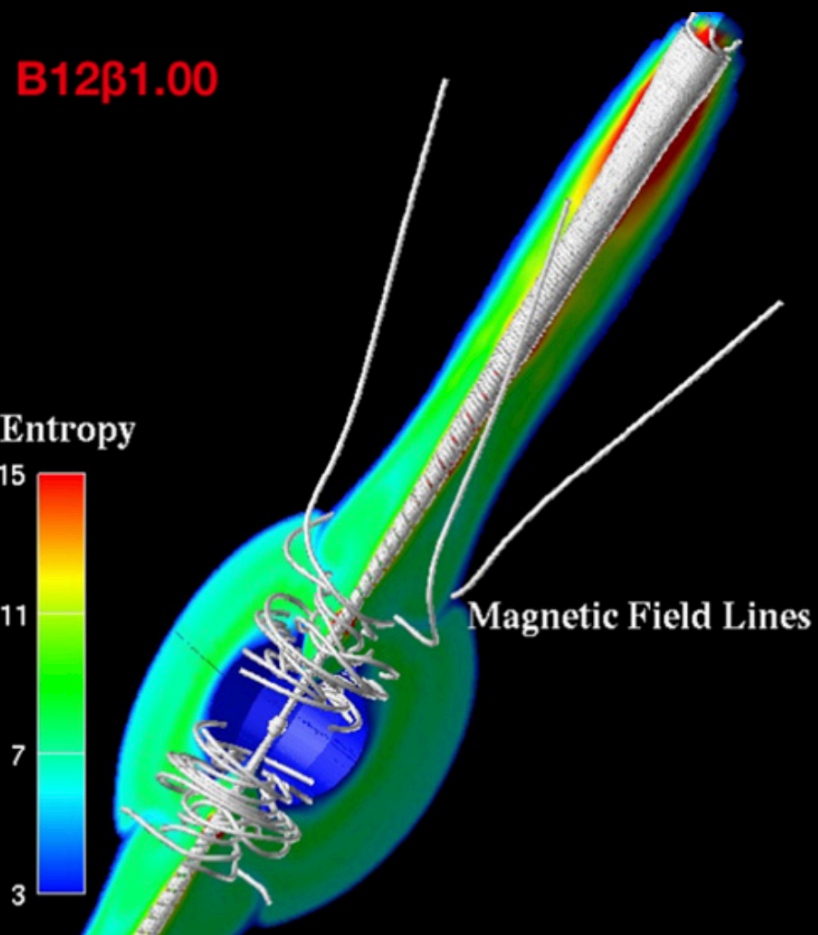
Correlated Anomaly in Meteorites  
between  $^{92,94}\text{Mo}$  and  $^{100}\text{Mo}$



Origin is in  
the same astrophysical site.



## MHD-Jet SNe



Nishimura et al 2015

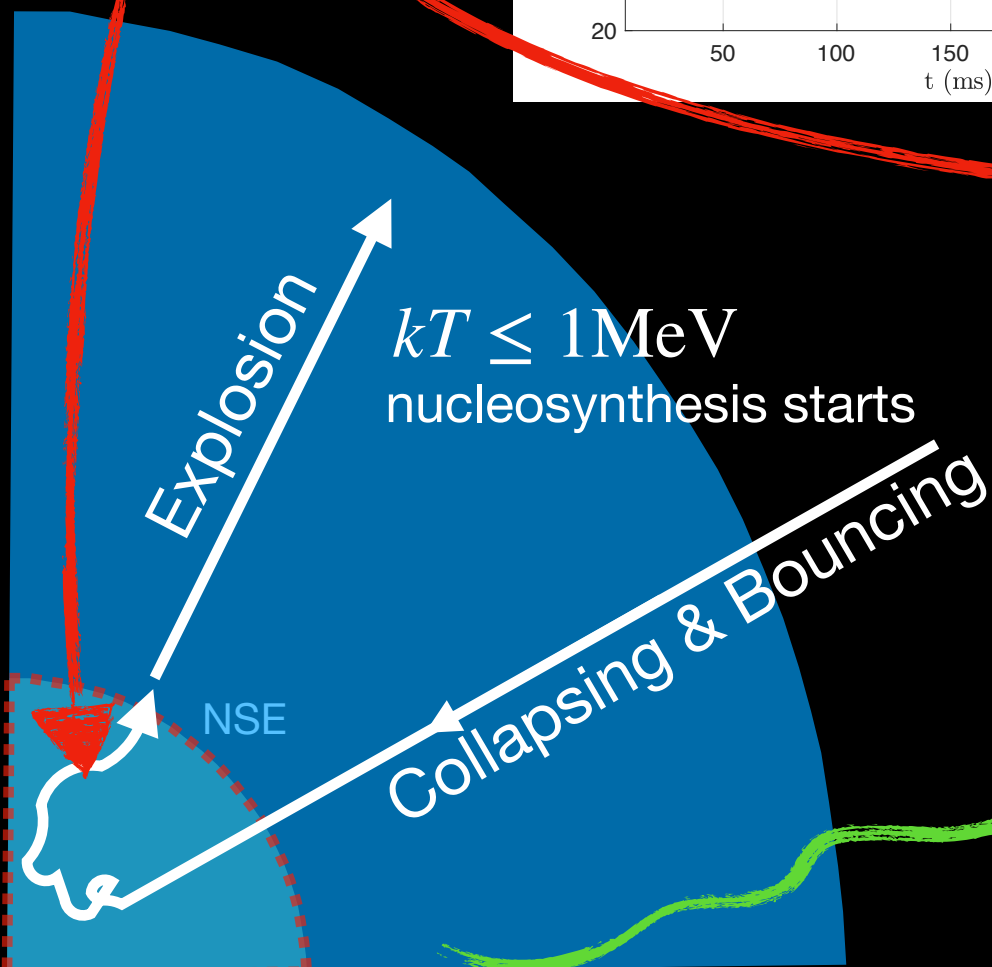
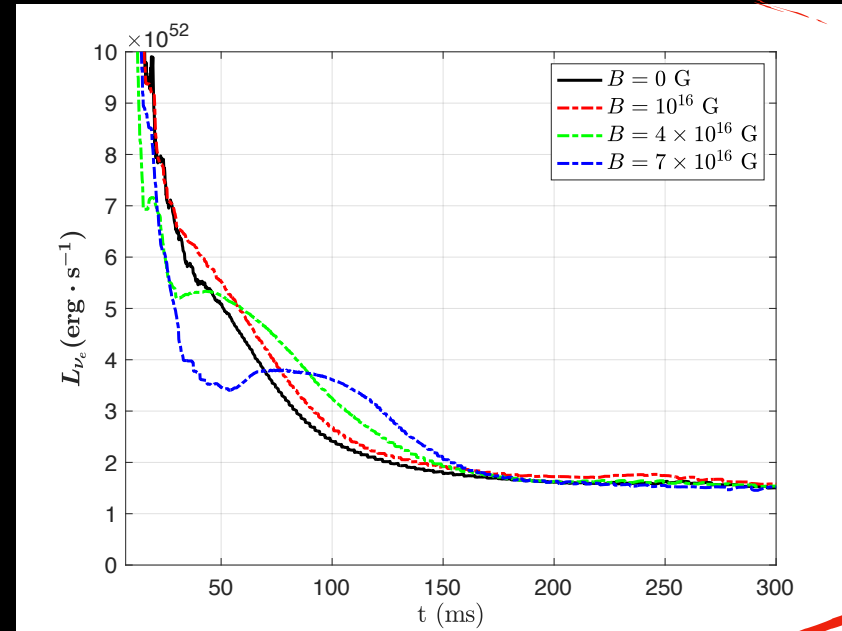
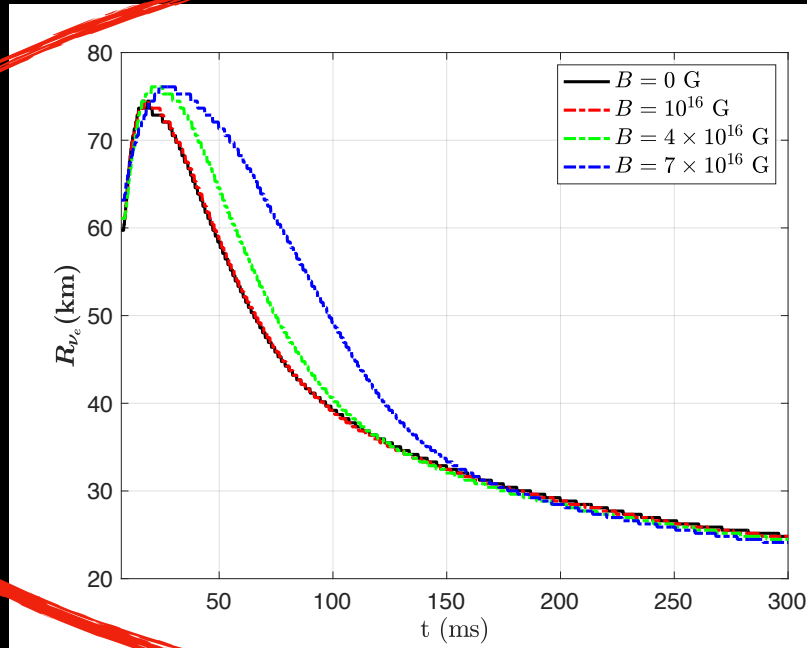
A site with both  $\nu p$ -process and  $r$ -process

High entropy  $\longrightarrow$  Neutron-rich

SNe  $\longrightarrow$  Neutrinos

# Why Magnetic field

- ▶ Inside the  $\nu$ -sphere, anti-neutrinos are more sensitive to B-field,  $\nu(\bar{\nu})$ -spheres are enlarged by B-field while luminosities are suppressed due to less energy release rate.



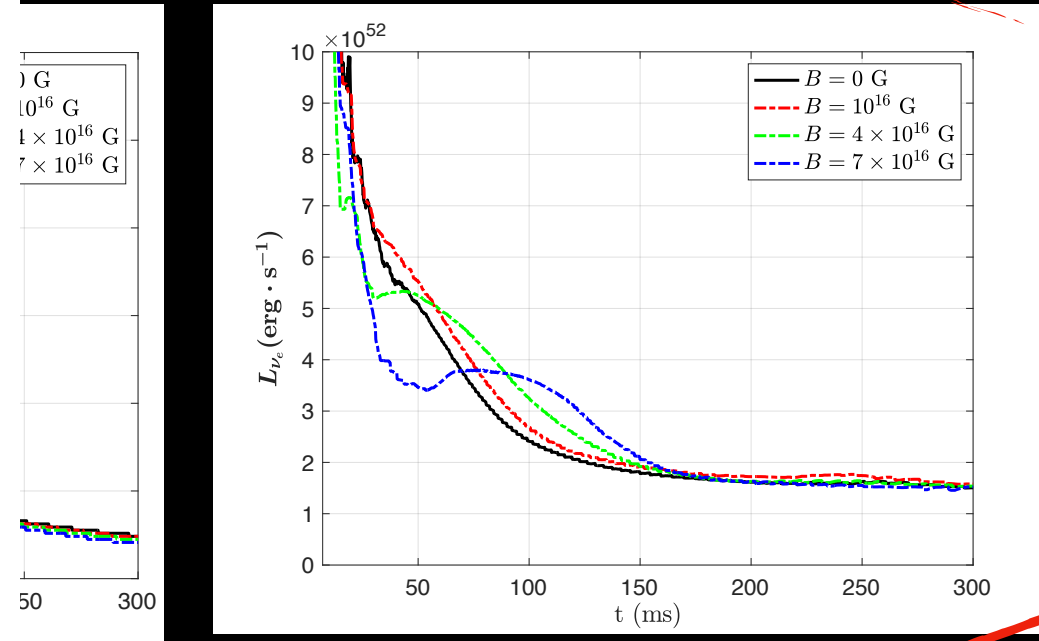
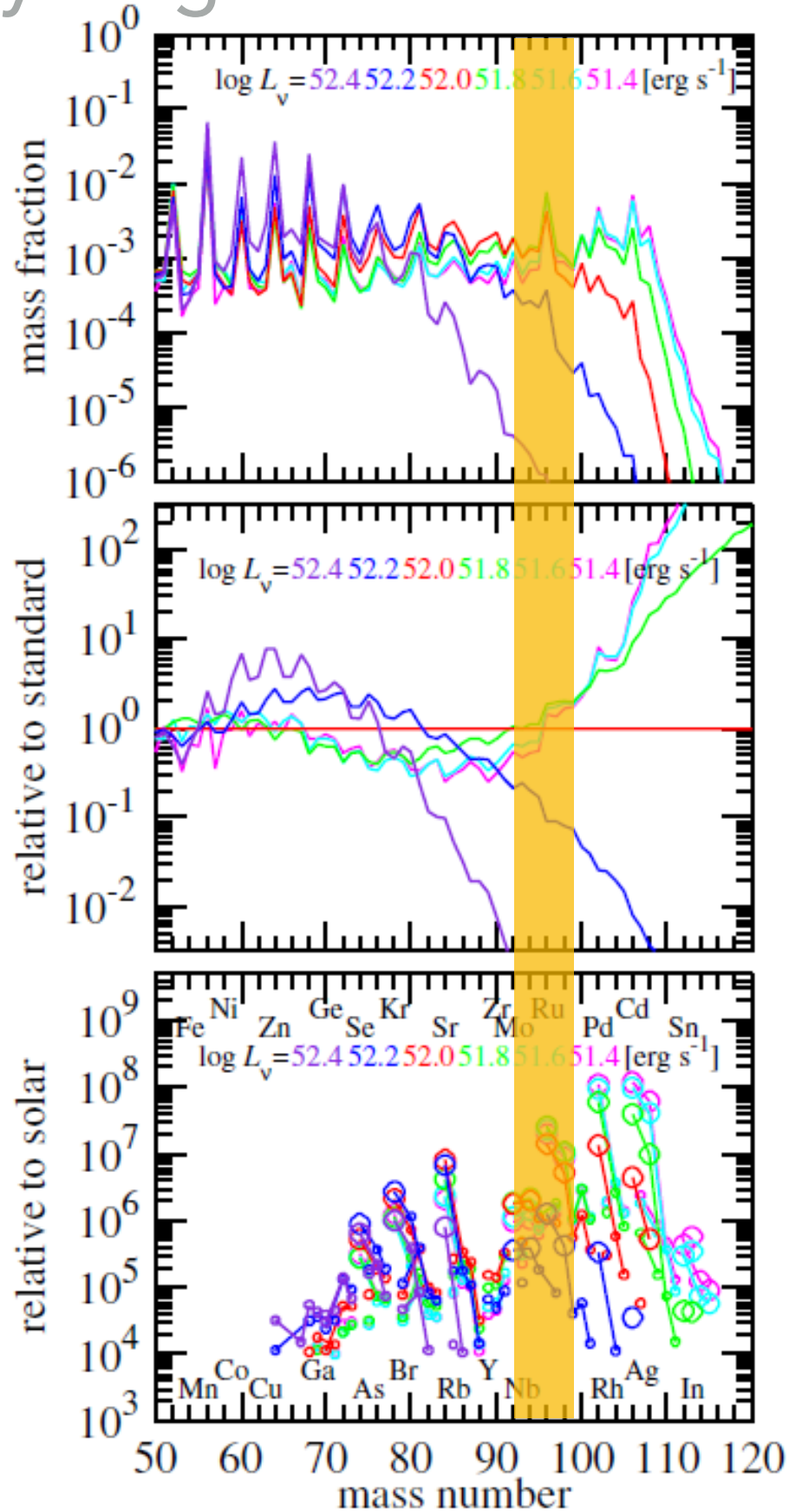
$\nu; \bar{\nu}$

$$\frac{d\phi_{\nu_e}}{dE_{\nu_e}} = \frac{L_{\nu_e}}{8\pi^2 R_\nu^2} \frac{E_{\nu_e}^2}{\exp(E_{\nu_e}/T_{\nu_e}) + 1}$$

# Why Magnetic field

- ▶ In
- en
- ra

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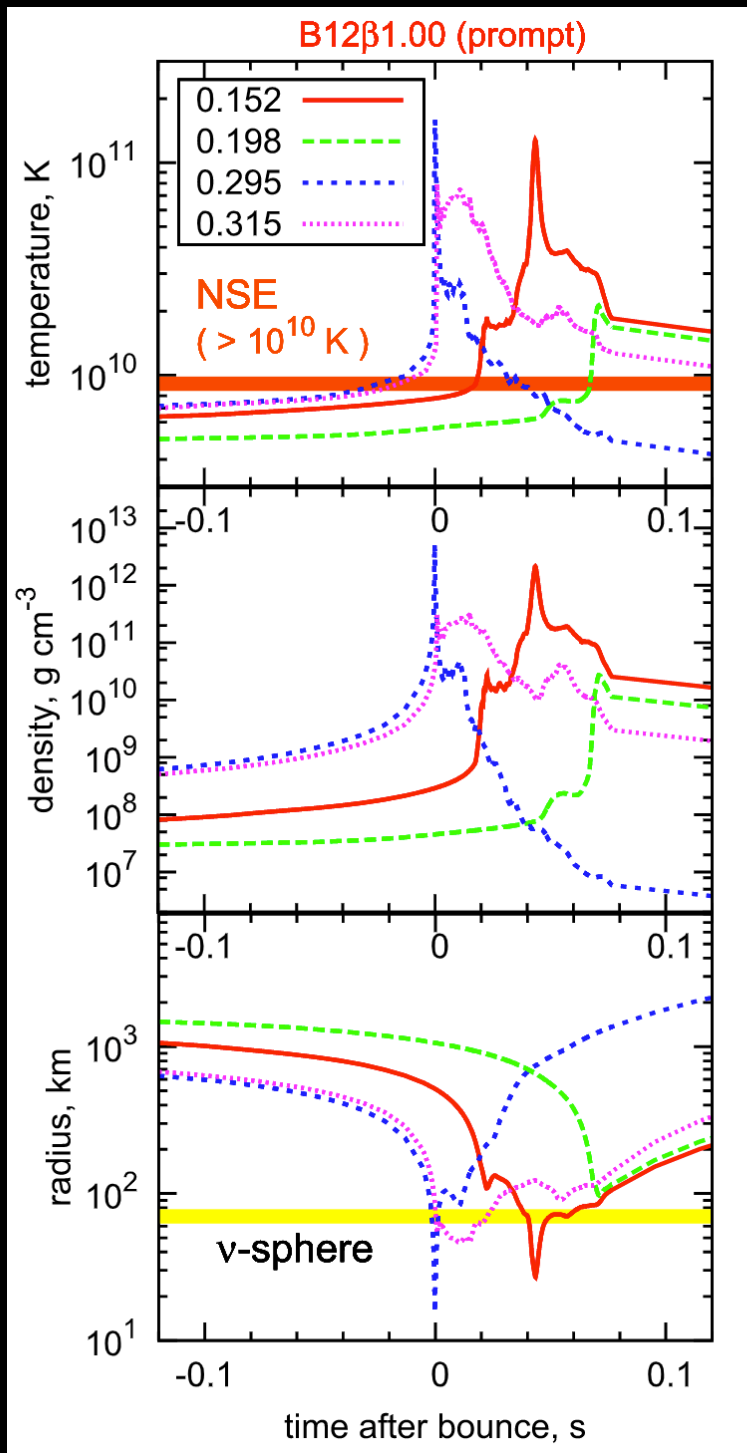


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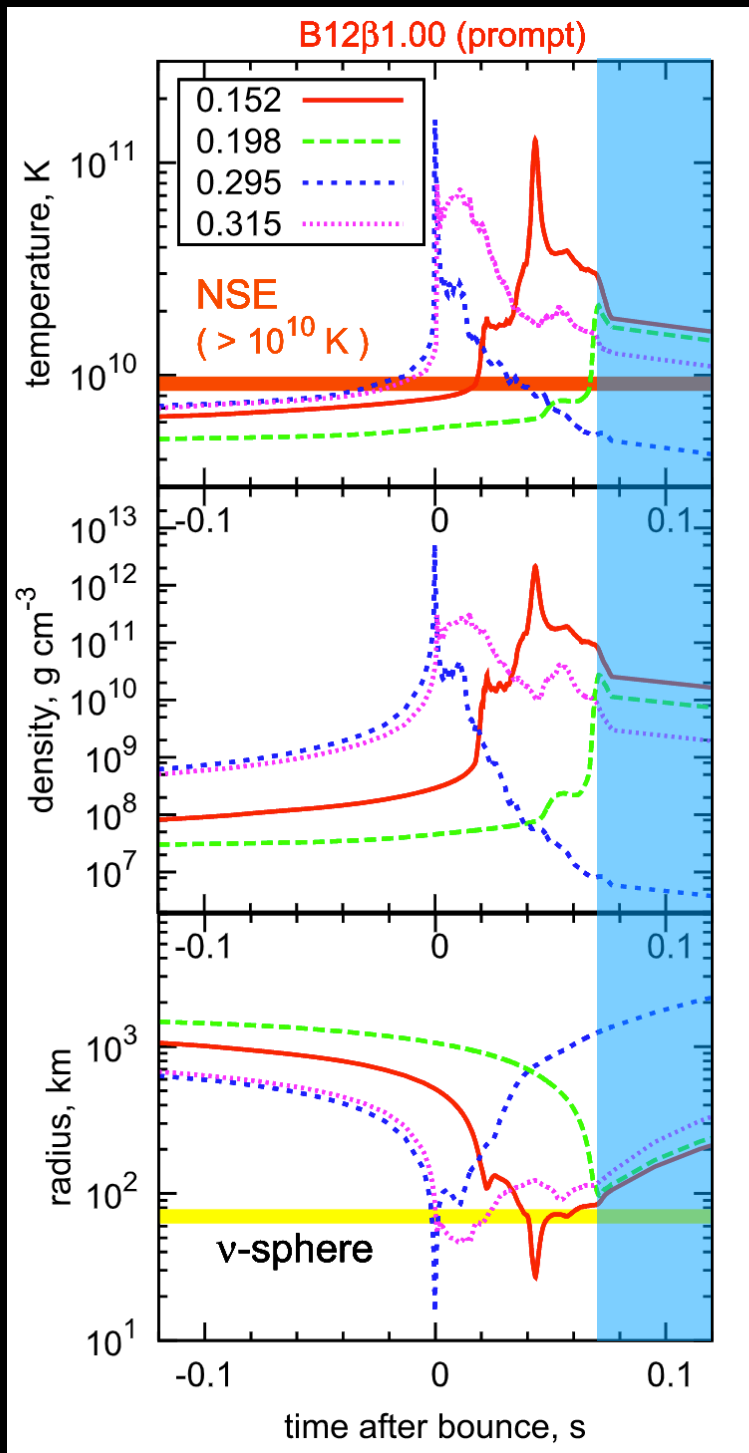
►  $Y_e$  evolution outside the  $\nu$  – sphere

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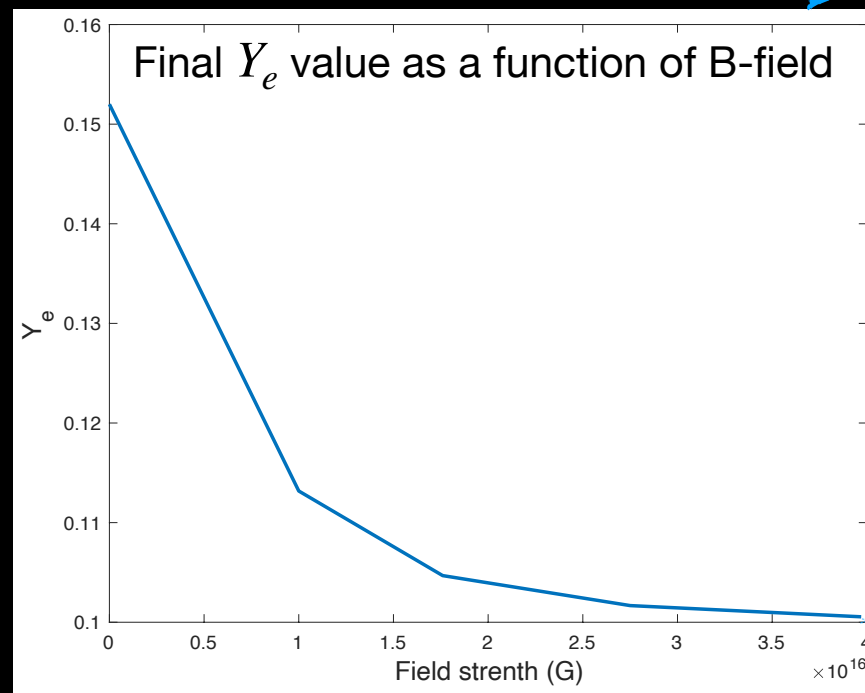
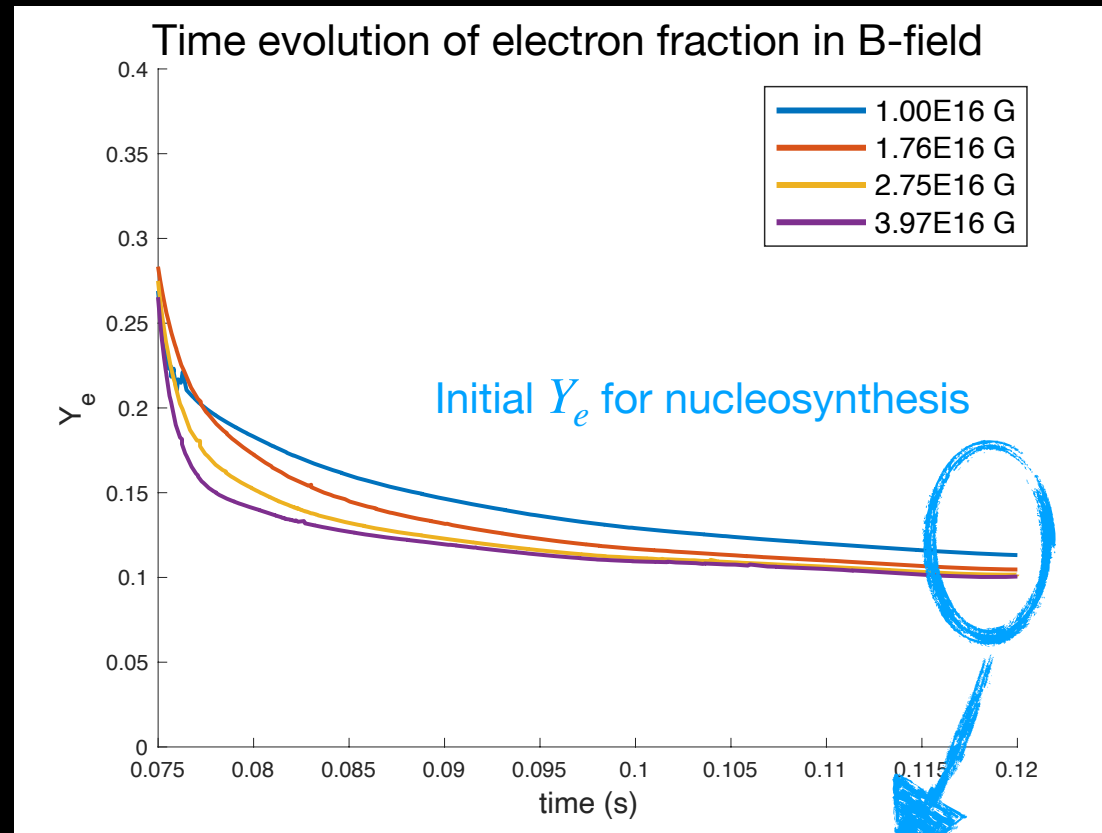


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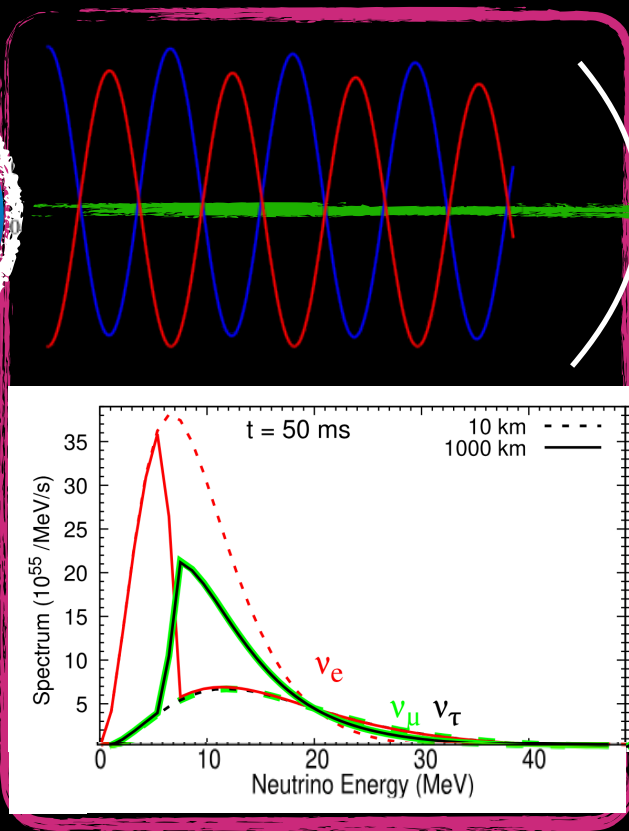
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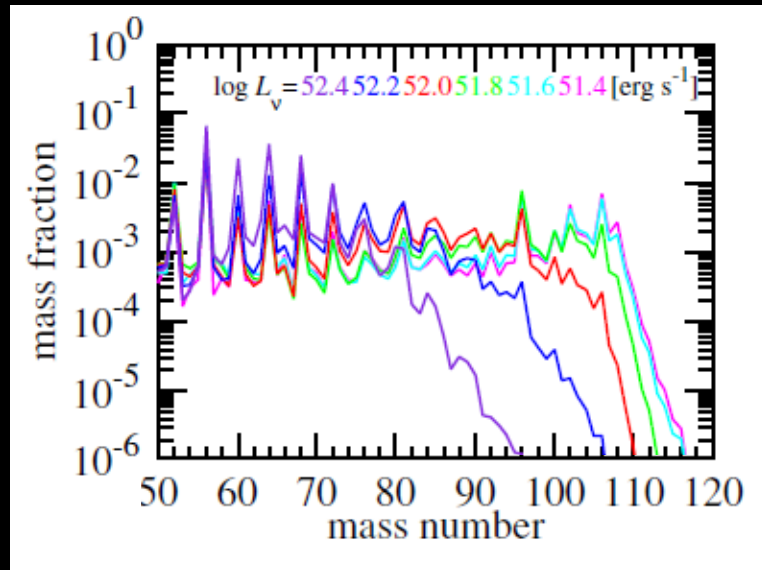
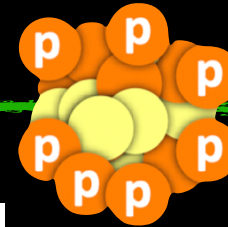
Nishimura et al 2015



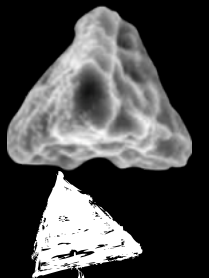
More neutron rich is expected



$\nu p$ -process  $\nu$ -process  
 $r$ -process



Meteorites



Prediction & Constraints

