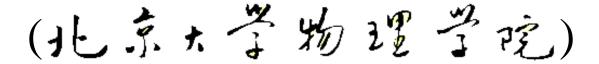


(上帝喜歡味對稱嗎?)



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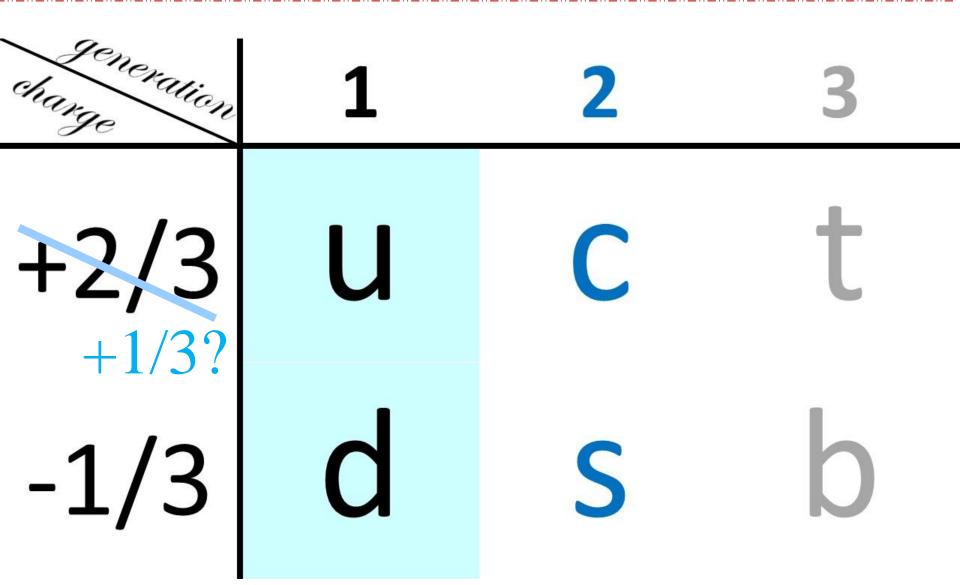
"QCS2023"

Yangzhou University, Sept. 23~26, 2023

From *isospin symmetry* and charge neutrality to the *atomic world*



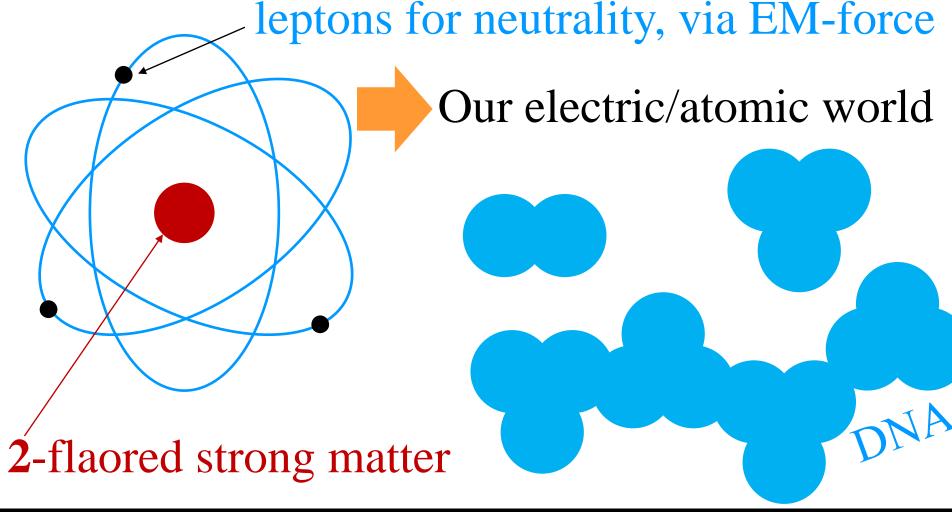
Isospin symmetry/charge neutrality



Isospin symmetry/charge neutrality

•We would therefore have atoms and beautiful world

leptons for neutrality, via FM-force



Isospin symmetry/charge neutrality

Certainly, including the beautiful site of QCS2023



Can we have both *flavor symmetry and charge* neutrality when facing "gigantic nucleus"?



•What is a gigantic nucleus? anticipated by Landau'32

ON THE THEORY OF STARS.

By L. Landau.

(Received 7 January 1932).

From the theoretical point of view the physical nature of Stellar equilibrium is considered.

The astrophysical methods usually applied in attacking lar structure are characterised by making s chosen only for the sake of mathephys matic By this is characterised, for instance, Mr. of the impossibility of a star consisting throu ical ideal gas; this proof rests on the rbitrary L and M, the fundamental equaasser tions sisting of classical ideal gas admit, in solution. Mr. Milne seems to have overgener looked at this assertion results only from the ty being constant throughout the star, assum s made only for mathematical purposes which and do with reality. Only in the case of this radius R disappears from the relation betwe necessary for regularity of the solution. sumptions about the opacity would lead to a relation between L, M and R, which relation would be eyi <mark>Oavido vioht lean dae</mark>l criticisms put for w Eddington's mass-luminosity-relation.

It seems ets problem of stellar structure by methods of theoretical physics, i. e. to investigate the physical nature of stellar purpose we must at first investigate the statistical equilibrium of a given mass without generation of energy, the condition for which equilibrium being the minimum of free energy F (for given temperature). The part of free energy due to gravitation is negative and inversely proportional to some

Landau L. 1932, Sov. Phys., 1, 285

S Landa

we have no need to suppose that the radiation of stars is due to some mysterious process of mutual annihilation of protons and electrons, which was never observed and has no special reason to occur in stars. Indeed we have always protons and electrons in atomic nuclei very close together, and they do not annihilate themselves; and it would be very strange if the high temperature did help, only because it does something in chemistry (chain reactions!). Following a beautiful idea of Prof. Niels Bohr's we are able to believe that the stellar radiation is due simply to a violation of the law of energy, which law, as Bohr has first pointed out, is no longer valid in the relativistic quantum theory, when the laws of ordinary quantum mechanics break down (as it is experimentally proved by continuous - rays - spectra and also made probable by theoretical considerations). We expect that this must occur when the density of matter becomes so great that atomic nuclei come in close contact, forming one gigantic nucleus.

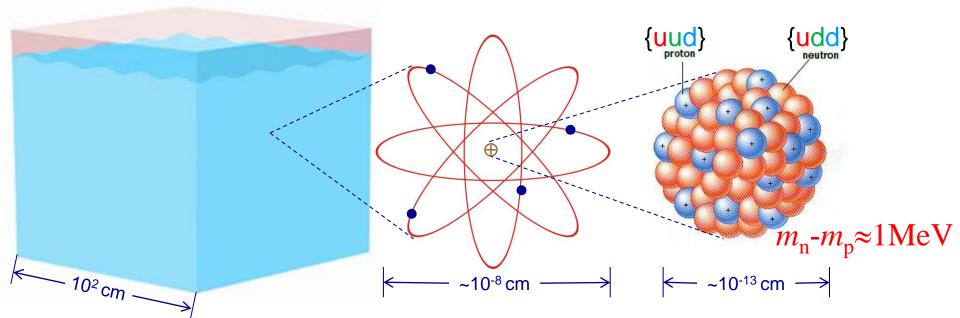
these general lines we can try to develop a theory of the structure. The central region of the star must consist of a core of highly condensed matter, surrounded by matter in ordinary state. If the transition between these two state were a continuous one, a mass $M < M_0$ would never form a star, because the normal equilibrium state the without particular region, would be out to all. Because, as far as we know, it is not one fact, we must conclude that the condensed and non-condensed states are separated by some as table state in h single manner as supplied its a four condensed which could be easily explained by some and of nuclear attraction. This would lead to the existence of a nearly discontinuous monadary between the two states.

The theory of stellar structure founded by the stove of siderations is yet to be constructed, and only such a theory can show how far they are true.

February 1931, Zurich.

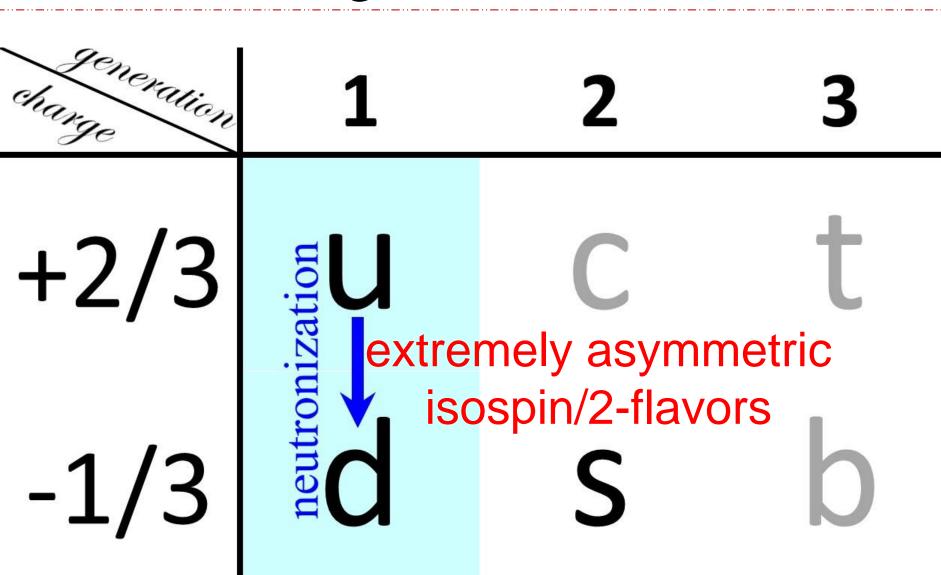
¹ L. Laudan und R. Peierls, ZS. L. Phys. 69, 56, 1931.

•Let's do an exercise...to squeeze one ton of water!



Before: baryon $A_{\text{water}} \sim 10^6 \text{g/u} \sim 10^{30}$, Electrons $E_e < 1 \text{MeV}$ After the squeezing:

- a, A giant "nucleus": $A_{\text{water}}^{1/3}$ fm~10μm at ~ ρ_{nucl}
- **b**, Pauli principle $\Rightarrow E_{\rm e} \sim 300 {\rm MeV} >> m_{\rm n} m_{\rm p} > m_{\rm e} c^2 !!!$



•Nucleons as the units: atomic nucleus \rightarrow NS

•First QCS: 2014,

KIAA @ PKU

•QCSII: 2017,

Yukawa Institute,

Kyoto

•QCSIII: 2019, Asia Pacific Center for Theoretical Physics,

Busan Kyoto

•QCSIV: 2023,

Yangzhou Univ.,

Sept. 23~26,

qcs2023@126.com+

Quarks and Compact Stars KIAA at Peking University, Beijing; Oct. 20-22, 2014

Though the standard model of particle physics is proved to be successfully perfect as the discovery of Higgs boson, the nature of strong interaction at low energy scale, which is essential for us to understand atomic nuclei and compact stars as well as the early Universe, is still far from the end. Certainly the quark degree of freedom could not be negligible, and compact stars provide a unique test-

non-perturbative behaviors of the s KIAA at Peking University hosts a bilateral meeting strong interto strengthen the researches and foster collaborations action. between China and Japan.

Organizers:

Toshiki Maruyama, Toshitaka Tatsumi and Renxin Xu

http://kiaa.pku.edu.cn/qcs2014

中子数 N

ground for studying the

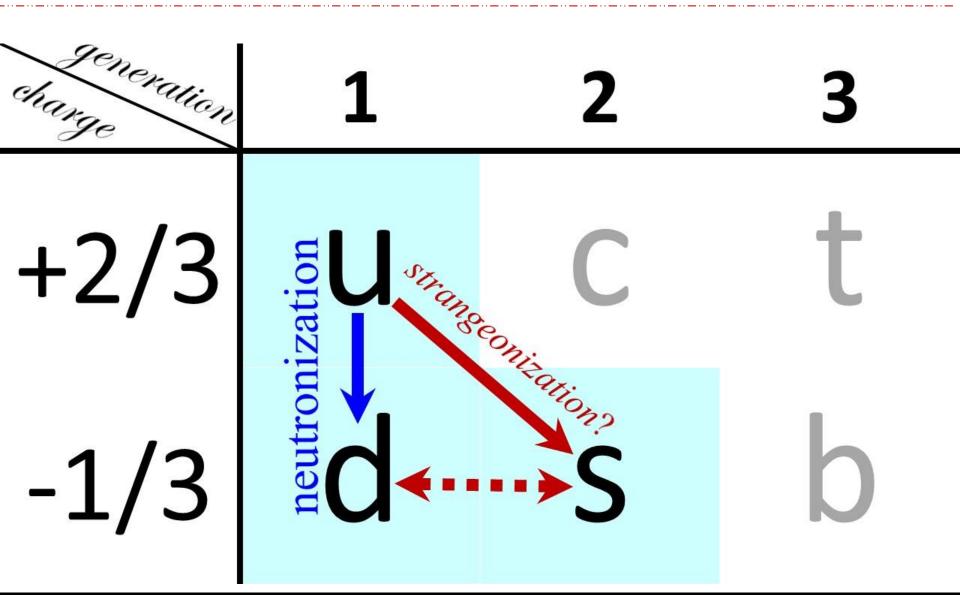
 10^{2}

 10^{57}

What if Nature loves *flavor symmetry*?

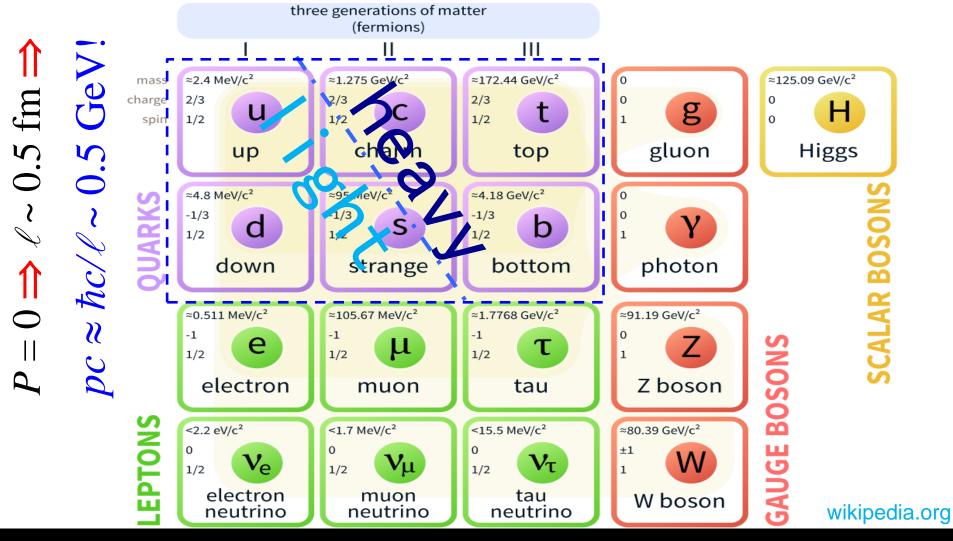


What if Nature loves fla-symmetry?



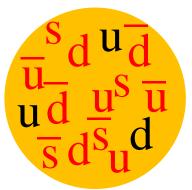
What if Nature loves fla-symmetry?

•Can we have both *fish*/fla.-sym. & *bear's paw*/neutrality?

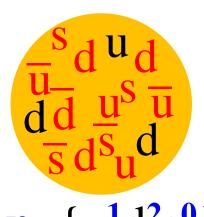


What if Nature loves fla.-symmetry?

•What's a proton? What's a neutron?



$$p = \{u^2d^1s^0\}$$



A perturbative calculation of quantum chromodynamics (QCD) may predict a nucleon sea with *light-flavor symmetry*, but the observed flavor asymmetry in the light-quark sea would be the result of the non-perturbative nature.

Strangeness and Hadron Structure

John Ellis^a

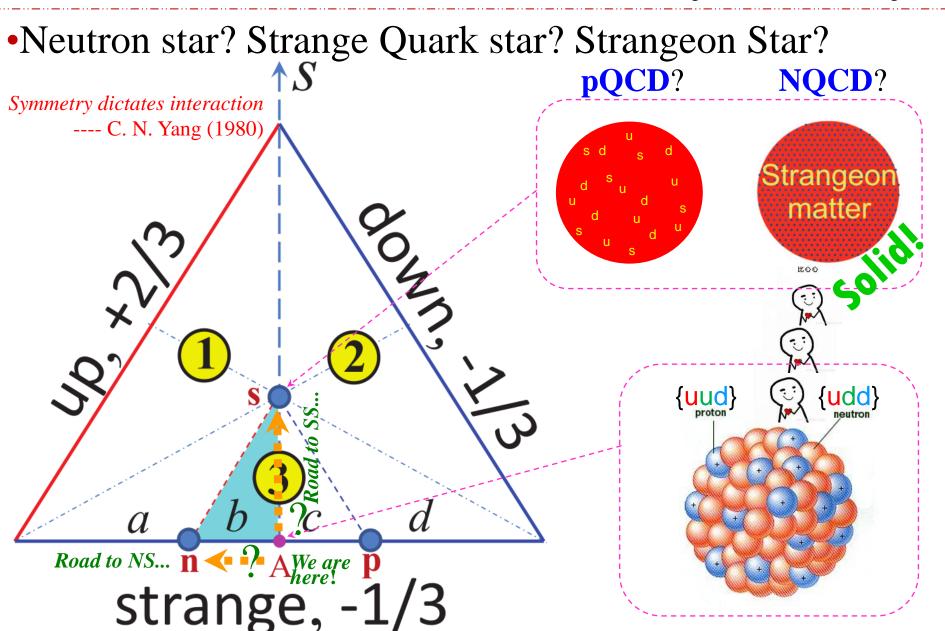
^aTheoretical Physics Division, CERN CH - 1211 Geneva 23

The nucleon wave function may contain a significant component of $\bar{s}s$ pairs, according to several measurements including the π -nucleon σ term, charm production and polarization effects in deep-inelastic scattering. In addition, there are excesses of ϕ production in LEAR and other experiments, above predictions based the naive Okubo-Zweig-Iizuka rule, that may be explained if the nucleon wave function contains a polarized $\bar{s}s$ component. This model also reproduces qualitatively data on Λ polarization in deep-inelastic neutrino scattering. The strange component of the proton is potentially important for other physics, such as the search for astrophysical dark matter.

CERN-TH/2000-112

hep-ph/0005322

What if Nature loves fla.-symmetry?



What if Nature loves fla.-symmetry?

- •Q: Strangeon matter matters? I don't
- •A: "Old" physics but particular consequence for us!



Trinity of Strangeon

Strangeon stars: compact objects

baryon number: $A \approx 10^{57}$!

Strangeon cosmic ray: UHECR

 $A = 10^{10} \iff E_{\text{rest}} \approx 10^{19} \text{ eV}$

Strangeon dark matter: cosmology

 $n_{\text{sdm}} \simeq 10^{-16} / A_{30} \text{ km}^{-3}$, Earth $\sim 80 / A_{30} \text{ yr}^{-1}$ Moon $\sim 6 / A_{30} \text{ yr}^{-1}$

Strangeon forms first $\Rightarrow A_{\text{sdm}} \ge A_c \sim 10^9$?

Summary

Summary

- •It has been more than *90 years* since Landau superficially anticipated "neutron"/neutral stars, but the real **basic unit** inside a gigantic nucleus is still a matter of debate, in fact.
- •In the triangle, the units could be **3**-flavour *symmetric strangeons* rather than **2**-flavour *asymmetric nucleons* if the Nature really loves symmetry when building the world.
- •To test the strangeon models in the future...

THANKS