

Formation of millisecond pulsars with He WDs, ultra-compact X-ray binaries, and gravitational wave sources

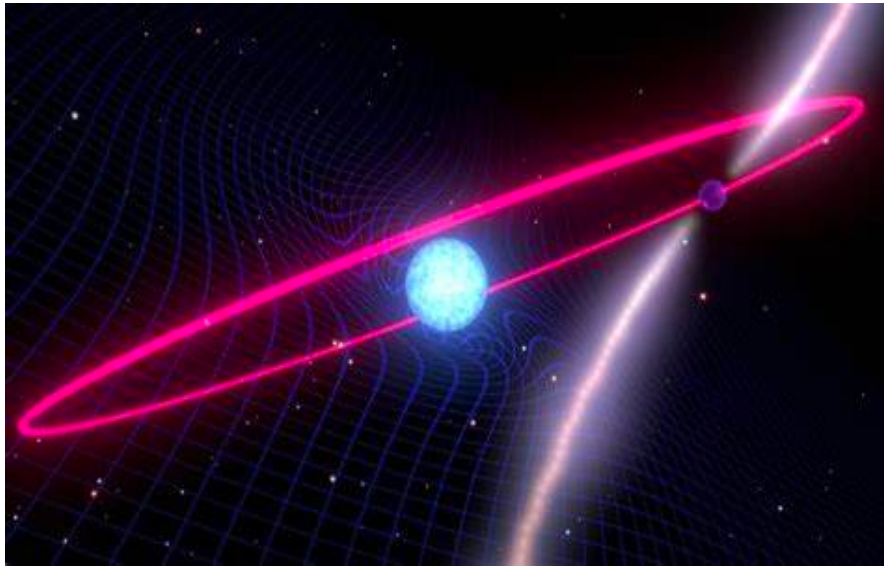
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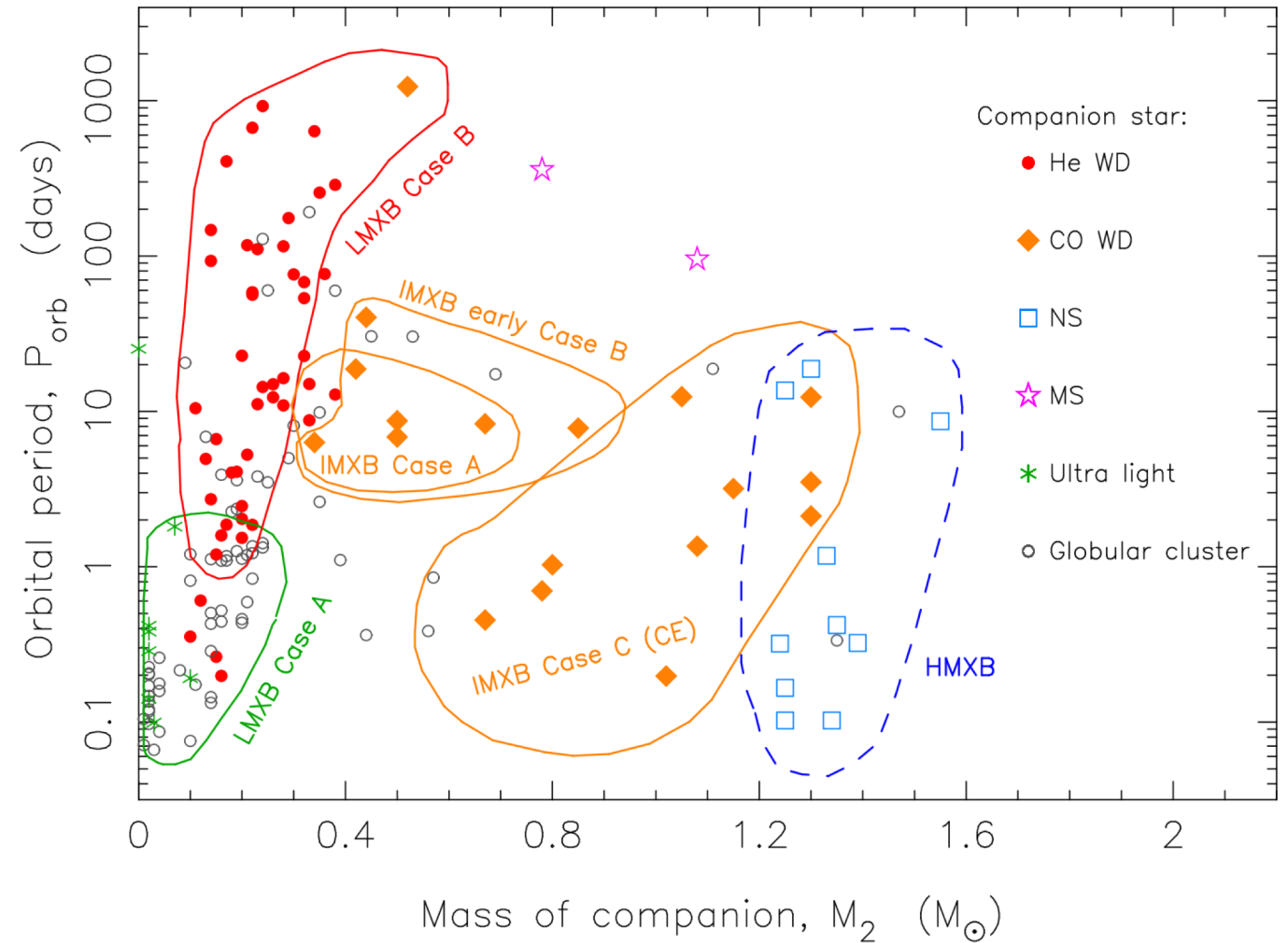
Yunnan observatories

Thomas Tauris, Xuefei Chen, Zhanwen Han

Binary millisecond pulsars

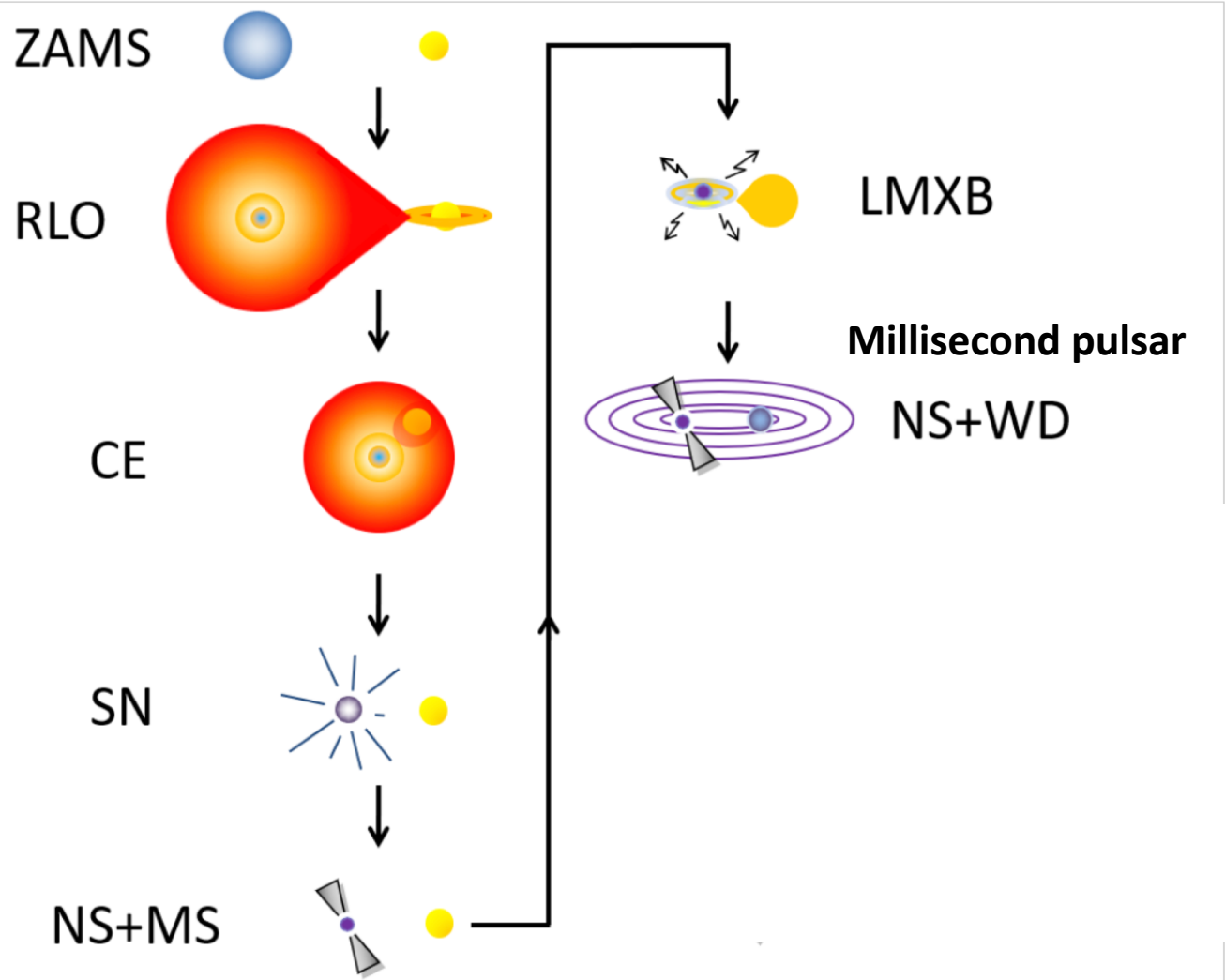


Different types of Companions



Credit: Thomas Tauris

Formation of MSP+He WD binaries



Binary millisecond pulsars with extremely low mass WDs

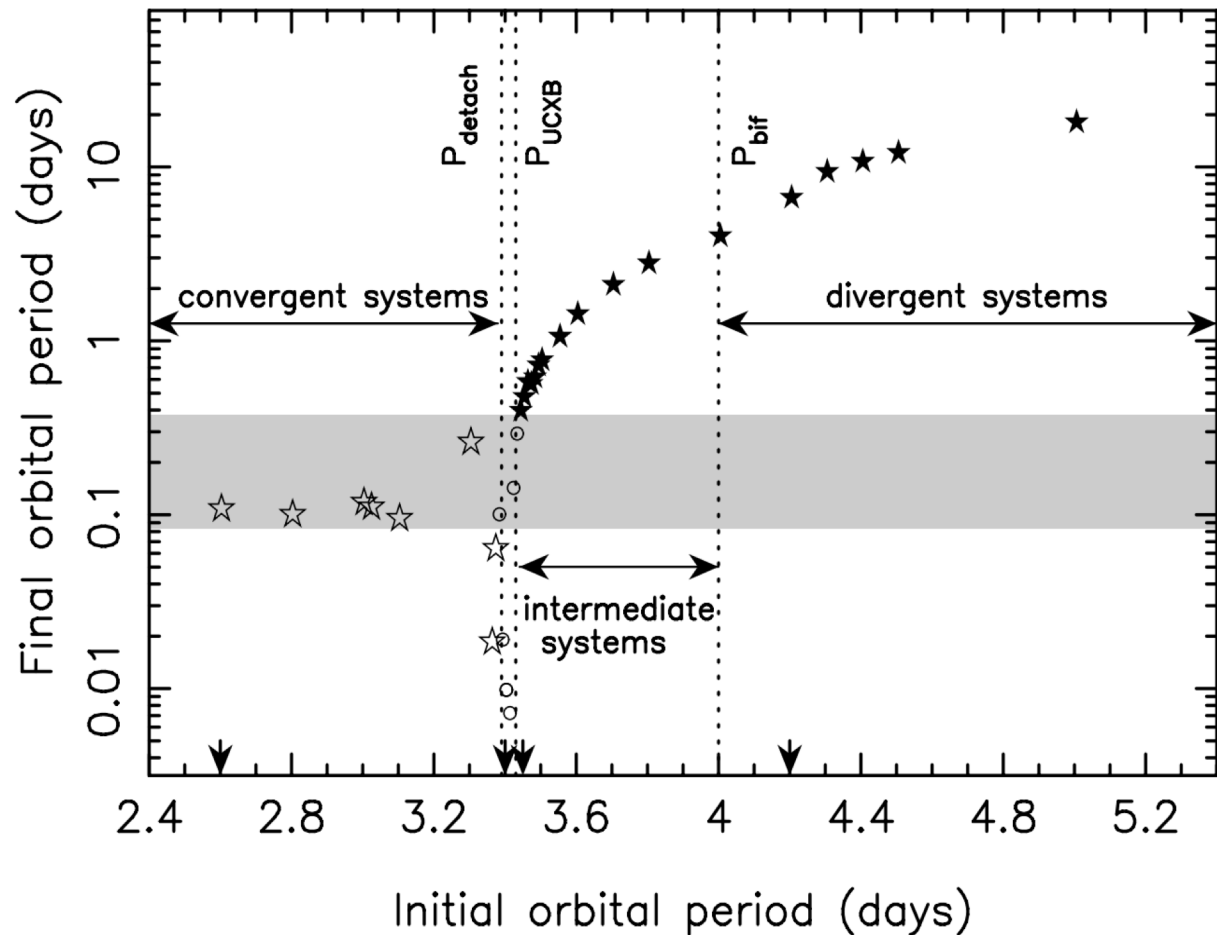
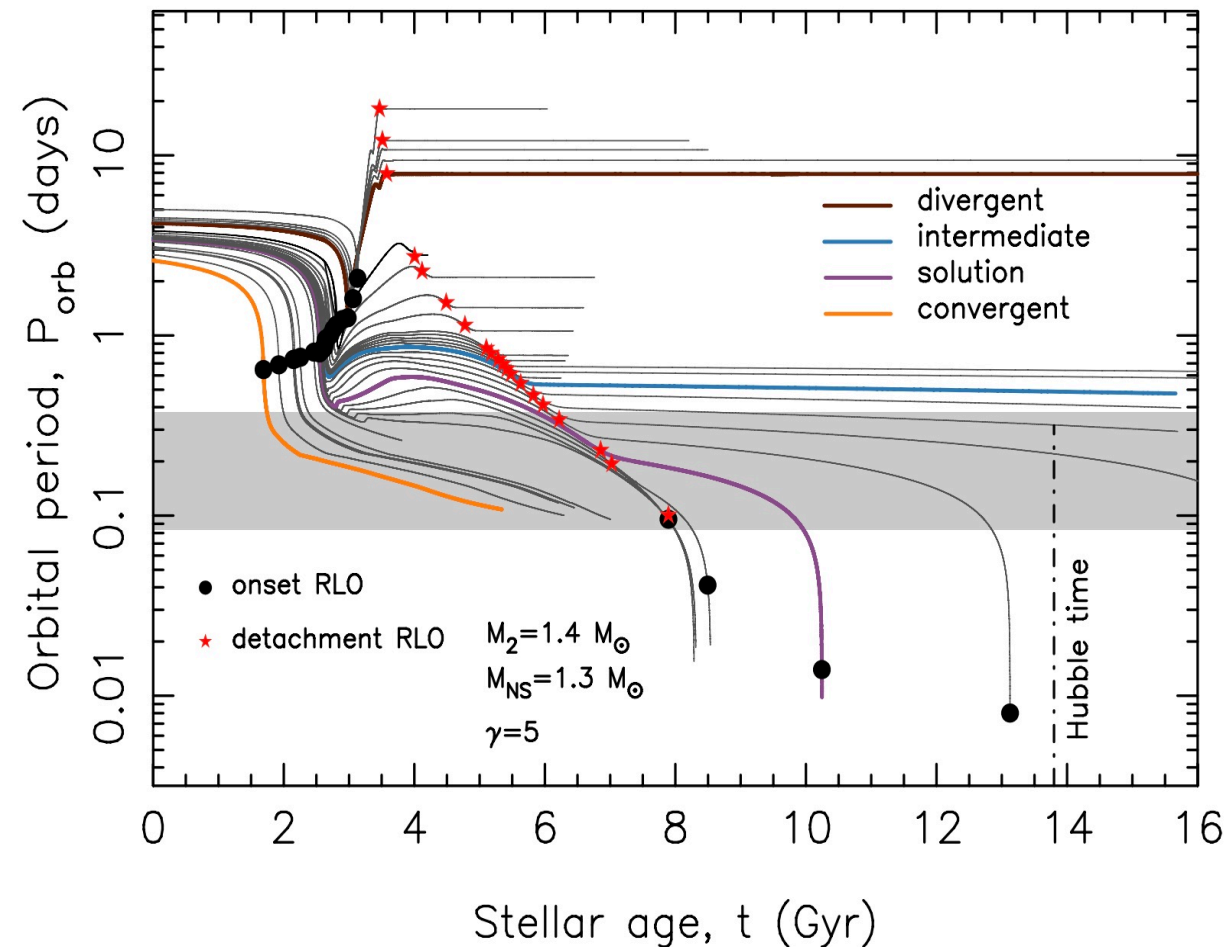
Pulsar name	P_{orb} (h)	M_{WD} (M_{\odot})	M_{NS} (M_{\odot})	Eccentricity	P_{spin} (ms)	\dot{P} (s s^{-1})	WD age (Gyr)	Optical data reference
PSR J0348+0432	2.46	0.17	2.01	2.6×10^{-6}	39.1	2.41×10^{-19}	2.1 ± 0.5	Antoniadis et al. (2013)
PSR J0751+1807	6.31	0.14	1.34	7.1×10^{-7}	3.48	7.79×10^{-21}	–	Lundgren et al. (1995)
PSR J1738+0333	8.52	0.18	1.47	4.0×10^{-6}	5.85	2.41×10^{-19}	–	Antoniadis et al. (2012)
PSR J1816+4510	8.66	0.21 ^a	2.0 ^a	8×10^{-6}	3.19	4.31×10^{-20}	–	Kaplan et al. (2013)
PSR J0024–7204U	10.3	0.15 ^b	1.5 ^b	$<10^{-4}$	4.34	–	0.6	Edmonds et al. (2001)
PSR J1748–2446M	10.6	0.17 ^b	1.5 ^b	$<10^{-4}$	3.57	–	–	
PSR J1748–2446V	12.1	0.15 ^b	1.5 ^b	$<10^{-4}$	2.07	–	–	
PSR J0024–7204Y	12.5	0.17 ^b	1.5 ^b	$<10^{-4}$	2.20	–	–	
PSR J1641+3627D	14.2	0.22 ^b	1.5 ^b	$<10^{-4}$	3.12	–	–	
PSR J1012+5307	14.5	0.16	1.64	$<8.4 \times 10^{-7}$	5.26	1.71×10^{-20}	–	Callanan et al. (1998)

Porb: 2~14 hr; WD mass \leq 0.22 Msun

Istrate et al. 2014

Fine-tuning Problem

Istrate et al. 2014



Convergent:

$$P_{\text{orb}}^{\text{final}} < P_{\text{orb}}^{\text{initial}}$$

Divergent:

$$P_{\text{orb}}^{\text{final}} > P_{\text{orb}}^{\text{initial}}$$

Intermediate:

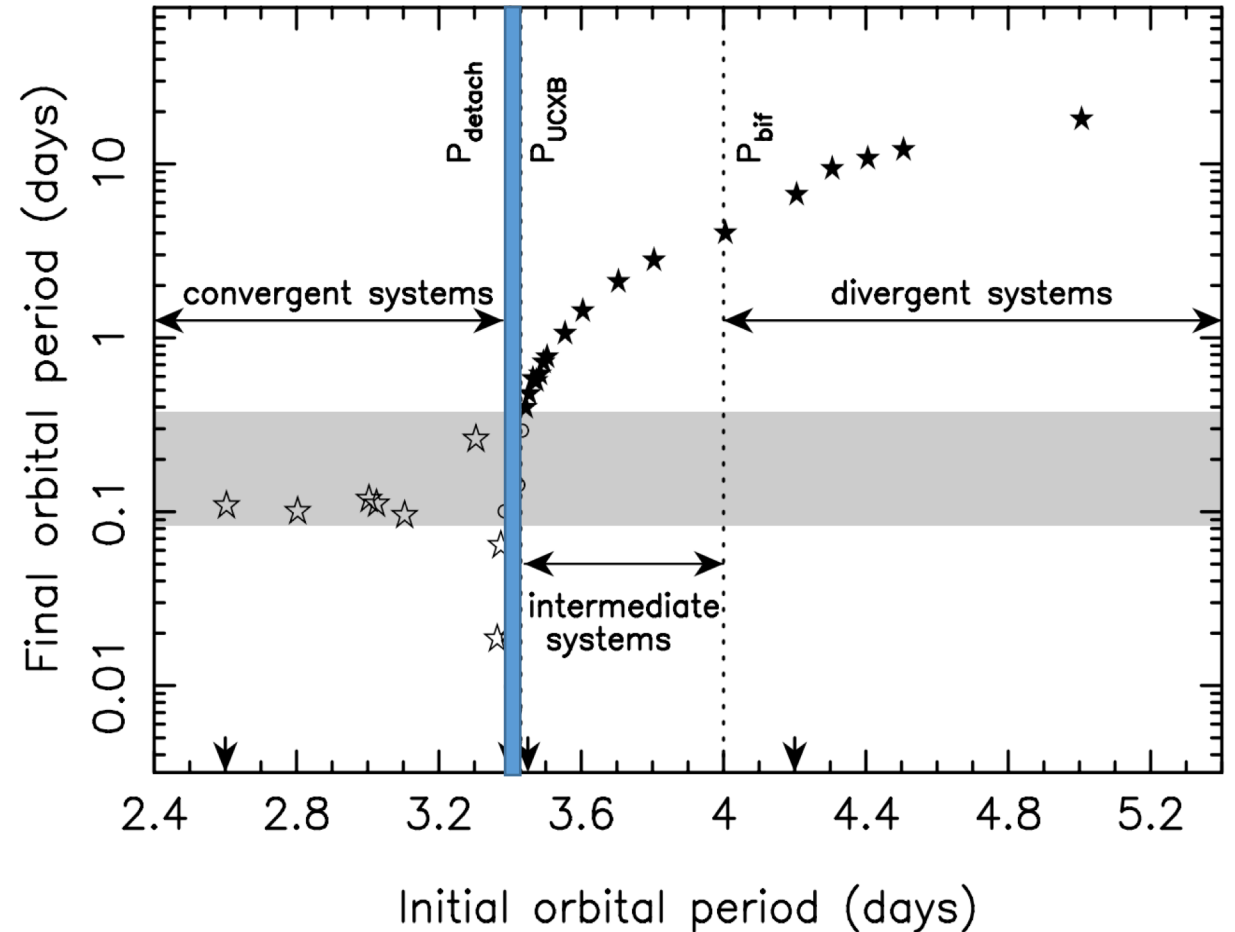
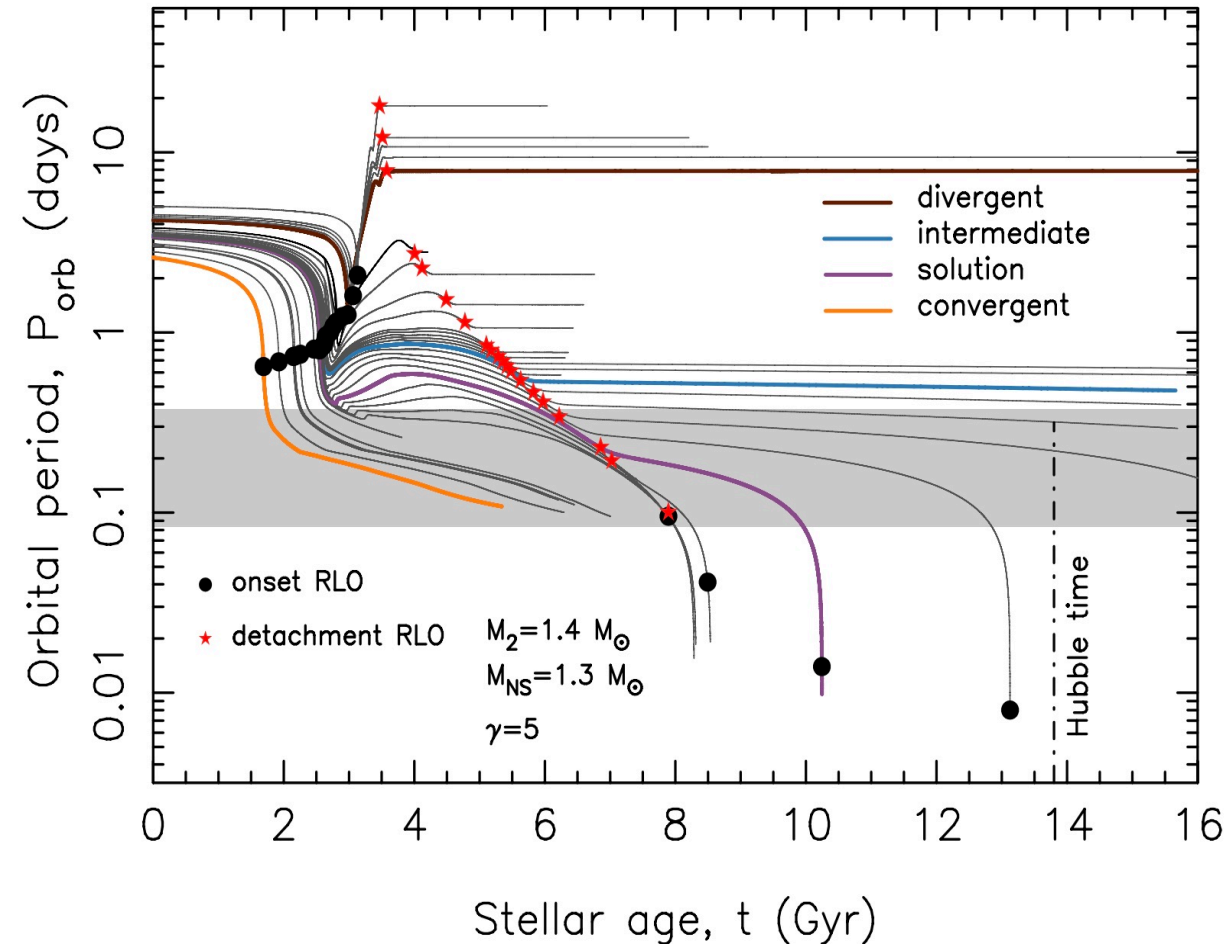
$$9 \text{ h} < P_{\text{orb}}^{\text{final}} < P_{\text{orb}}^{\text{initial}}$$

Grey Area:

Orbital period range of observed MSP+ ELM WDs

Fine-tuning Problem

Istrate et al. 2014



Convergent:

$$P_{\text{orb}}^{\text{final}} < P_{\text{orb}}^{\text{initial}}$$

Divergent:

$$P_{\text{orb}}^{\text{final}} > P_{\text{orb}}^{\text{initial}}$$

Intermediate:

$$9 \text{ h} < P_{\text{orb}}^{\text{final}} < P_{\text{orb}}^{\text{initial}}$$

Initial orbital period: 3.39--3.43 days

The parameter space of LMXBs for producing detached MSP - ELM WD is too small.

Binary evolution model

- NS mass: 1.3 Msun; Donor mass: 1.0 Msun ~ 4.0 Msun;
Orbital Period: 0.40 ~ 1000 day
- Solar Metallicity
- Stellar Wind: Reimers (1975)
- Donor star: initially synchronized.
- Rotation:

Chemical mixing: Secular shear instability, Eddington-sweet circulation, dynamical shear instability and Goldreich-Schubert-Fricke instability

Transport of angular momentum: Spruit-Tayler dynamo

The logo for the MESA stellar evolution code, consisting of the letters 'MESA' in a blue, stylized, sans-serif font.

Stellar evolution code

Angular momentum loss

- Gravitational Wave Radiation

$$\frac{dJ_{\text{gr}}}{dt} = -\frac{32}{5} \frac{G^{7/2}}{c^5} \frac{M_{\text{NS}}^2 M_2^2 (M_{\text{NS}} + M_2)^{1/2}}{a^{7/2}}$$

- Mass loss

accretion efficiency: 30%

Eddington accretion rate:

$$\dot{M}_{\text{Edd}} = \frac{4\pi G M_{\text{NS}}}{c \cdot 0.20(1 + X) \cdot \eta}$$

Material not accreted by NS takes away the specific angular momentum of NS.

- Spin-orbit coupling (Paxton et al. 2015)
- Magnetic braking

New magnetic braking prescription

Van et al. 2019

$$\frac{dJ_{\text{mb}}}{dt} = \frac{dJ_{\text{mb,Sk}}}{dt} \left(\frac{\omega}{\omega_{\odot}} \right)^{\beta} \left(\frac{\tau_{\text{conv}}}{\tau_{\odot,\text{conv}}} \right)^{\xi} \left(\frac{\dot{M}_{\text{W}}}{\dot{M}_{\odot,\text{W}}} \right)^{\alpha}$$

$$\frac{dJ_{\text{mb,Sk}}}{dt} = -3.8 \times 10^{-30} M_2 R_{\odot}^4 \left(\frac{R_2}{R_{\odot}} \right)^{\gamma_{\text{mb}}} \omega^3 \text{ dyne cm}$$

- Reproduce the persistent LMXBs
- Improvements in reproducing transient LMXBs

New magnetic braking prescription

Van et al. 2019

$$\frac{dJ_{\text{mb}}}{dt} = \frac{dJ_{\text{mb,Sk}}}{dt} \left(\frac{\omega}{\omega_{\odot}}\right)^{\beta} \left(\frac{\tau_{\text{conv}}}{\tau_{\odot,\text{conv}}}\right)^{\xi} \left(\frac{\dot{M}_{\text{W}}}{\dot{M}_{\odot,\text{W}}}\right)^{\alpha}$$

$$\frac{dJ_{\text{mb,Sk}}}{dt} = -3.8 \times 10^{-30} M_2 R_{\odot}^4 \left(\frac{R_2}{R_{\odot}}\right)^{\gamma_{\text{mb}}} \omega^3 \text{ dyne cm}$$

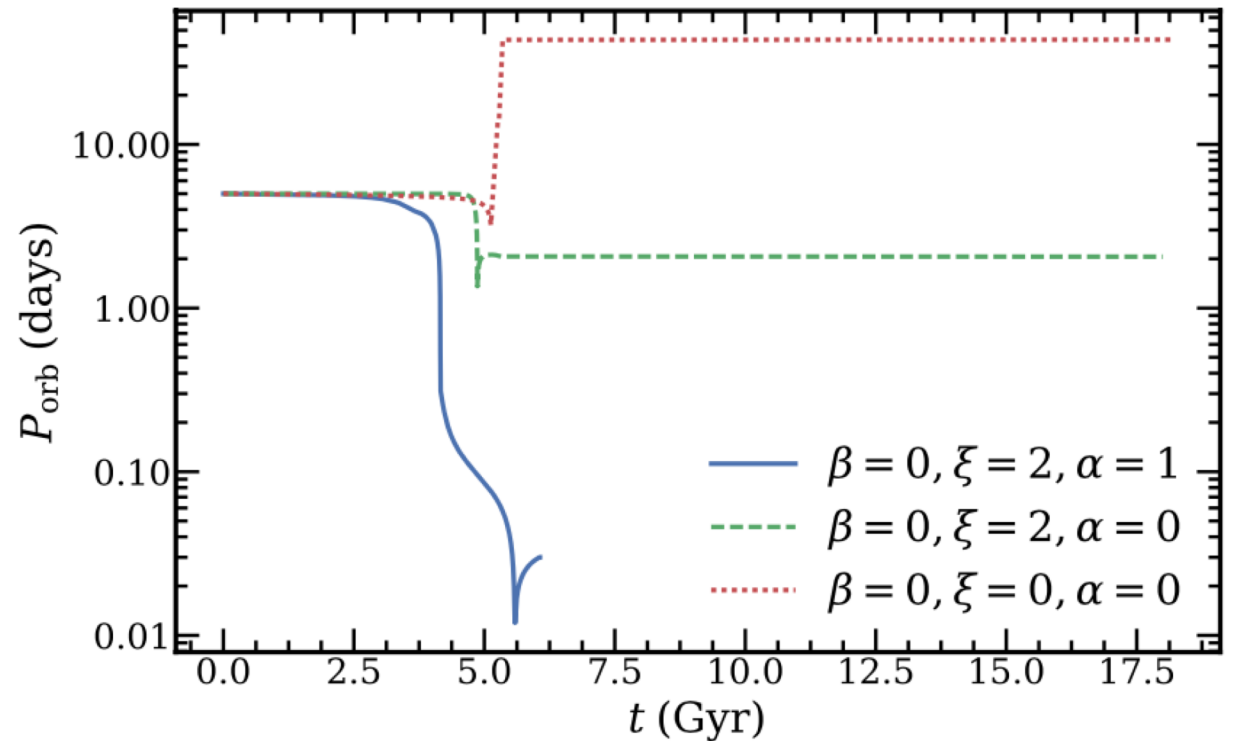
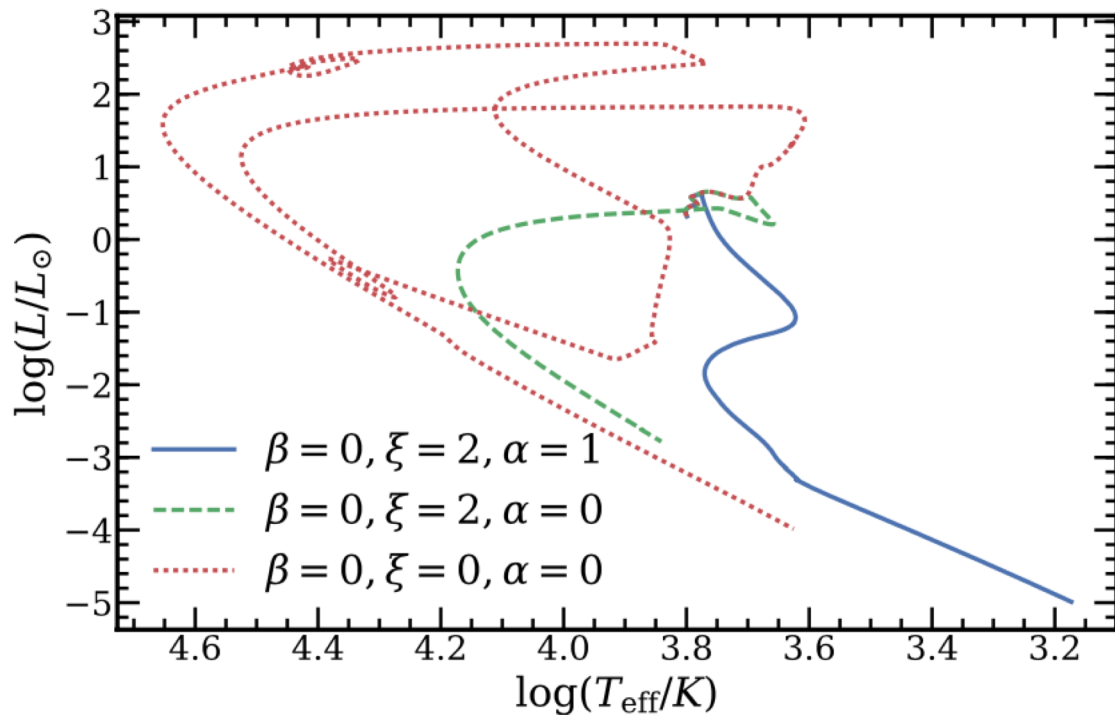
Case	β	ξ	α
1 – Default Skumanich	0	0	0
2 – Convection boosted	0	2	0
3 – Intermediate	0	2	1
4 – Wind boosted	2	4	1

Models with $\beta = 0$, $\xi = 2$ and $\alpha = 1$ can better reproduce the observed results of X-ray binaries (Van et al. 2019).

Influence of MB on binary evolution

NS mass: 1.30 Msun, Donor mass: 1.25 Msun, Orbital period: 5.0 d

Chen et al. MNRAS, 2021

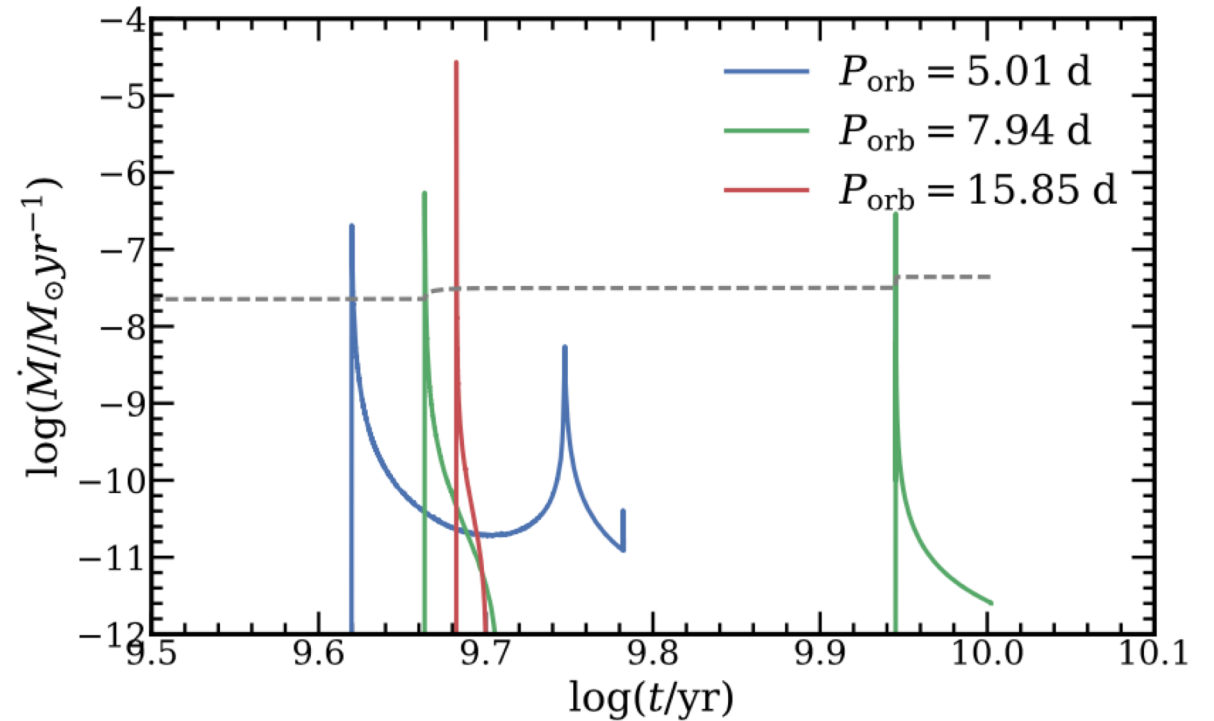
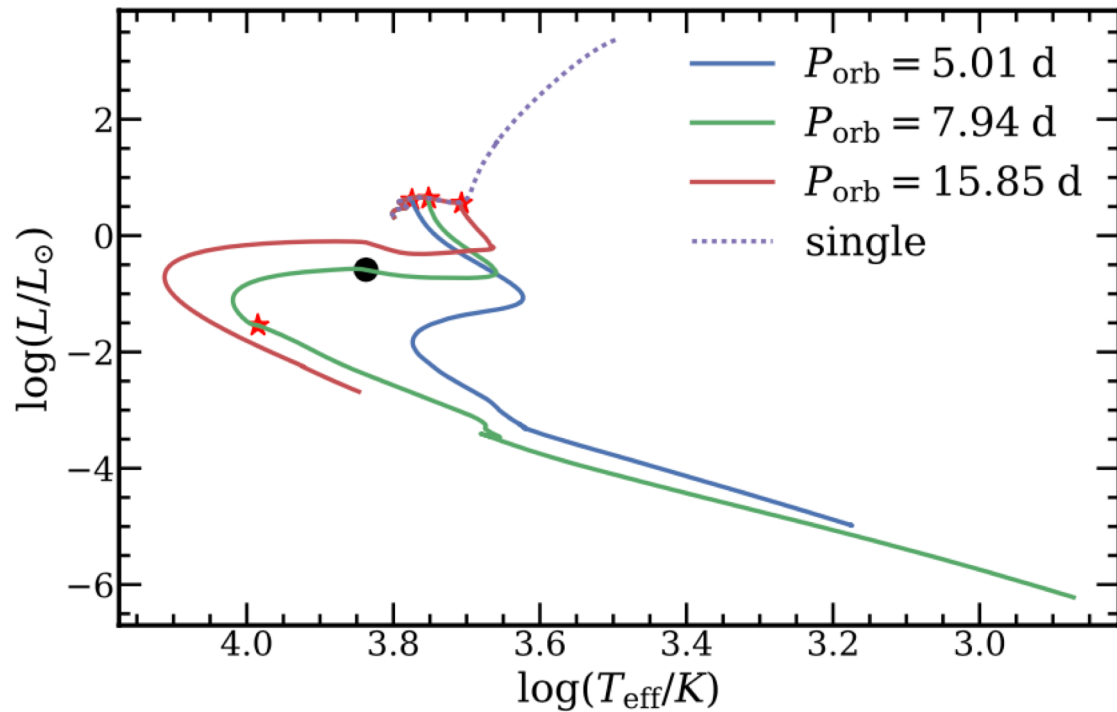


MB has an important influence on the evolution of LMXBs.

Examples of Binary evolution

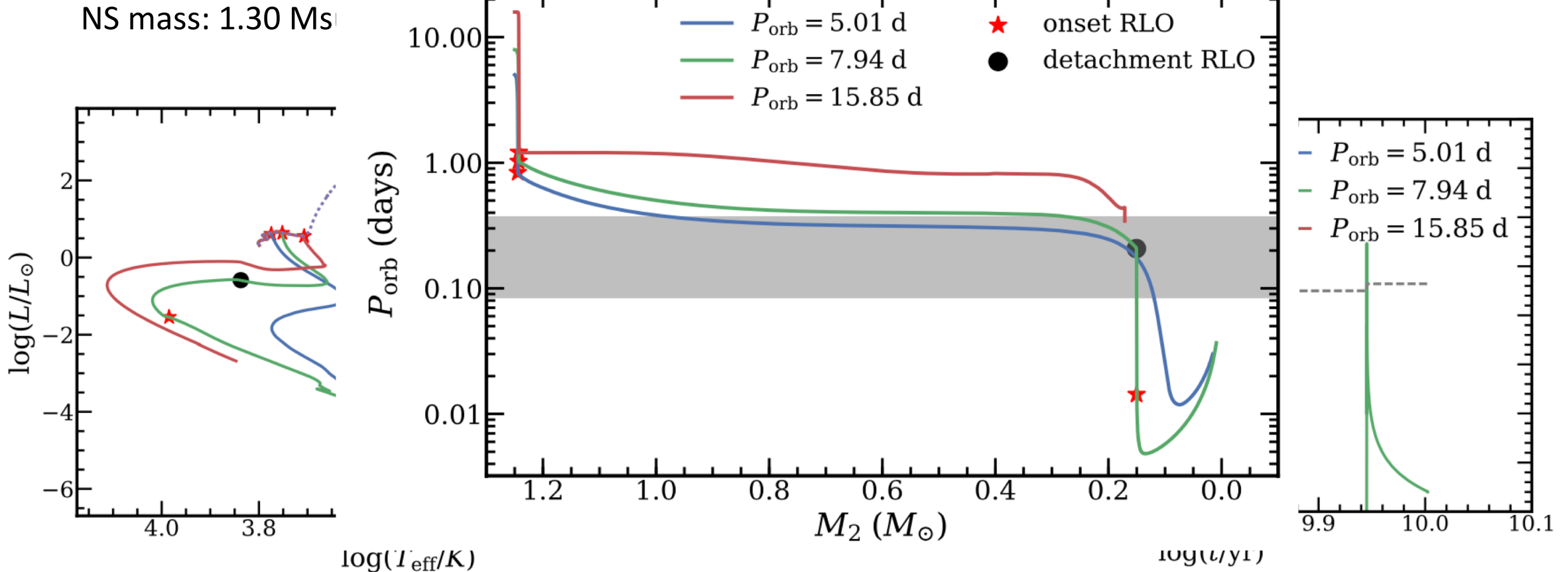
Chen et al. MNRAS, 2021

NS mass: 1.30 Msun, Donor mass: 1.25 Msun, MB: intermediate



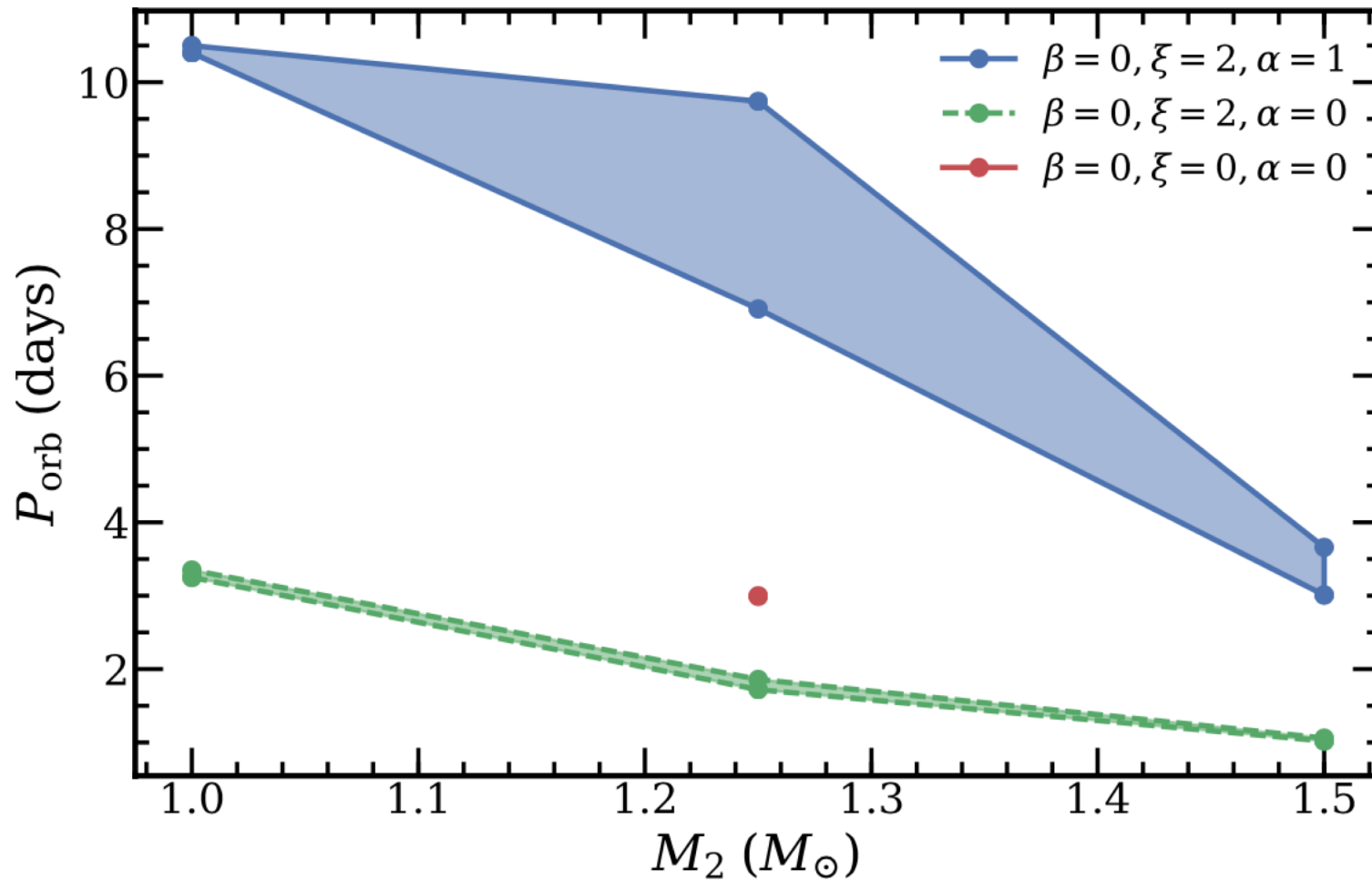
Examples of Binary evolution

Chen et al. MNRAS 2021



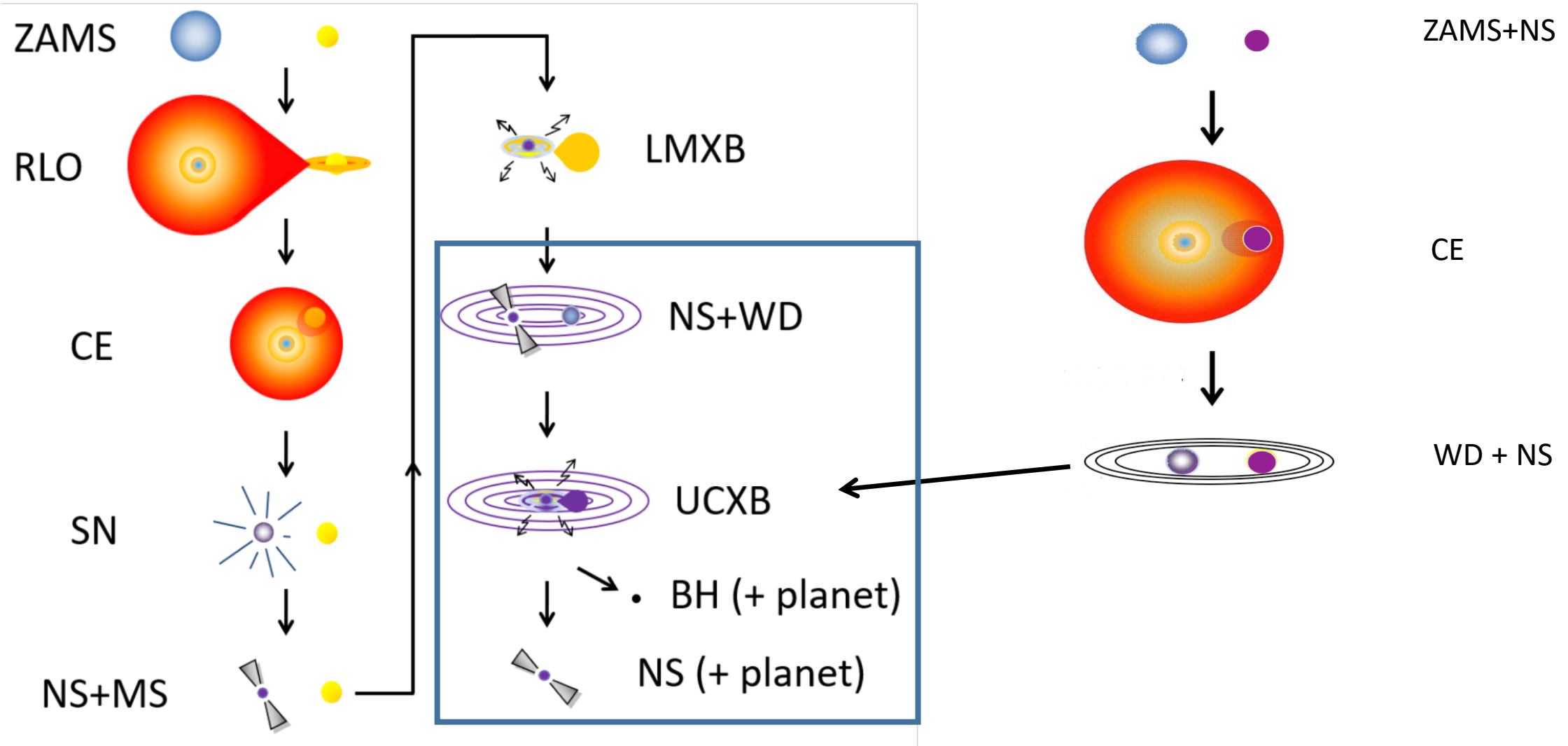
Grey Area:
Orbital period range of observed MSP+ ELM WDs

Initial Parameter Space



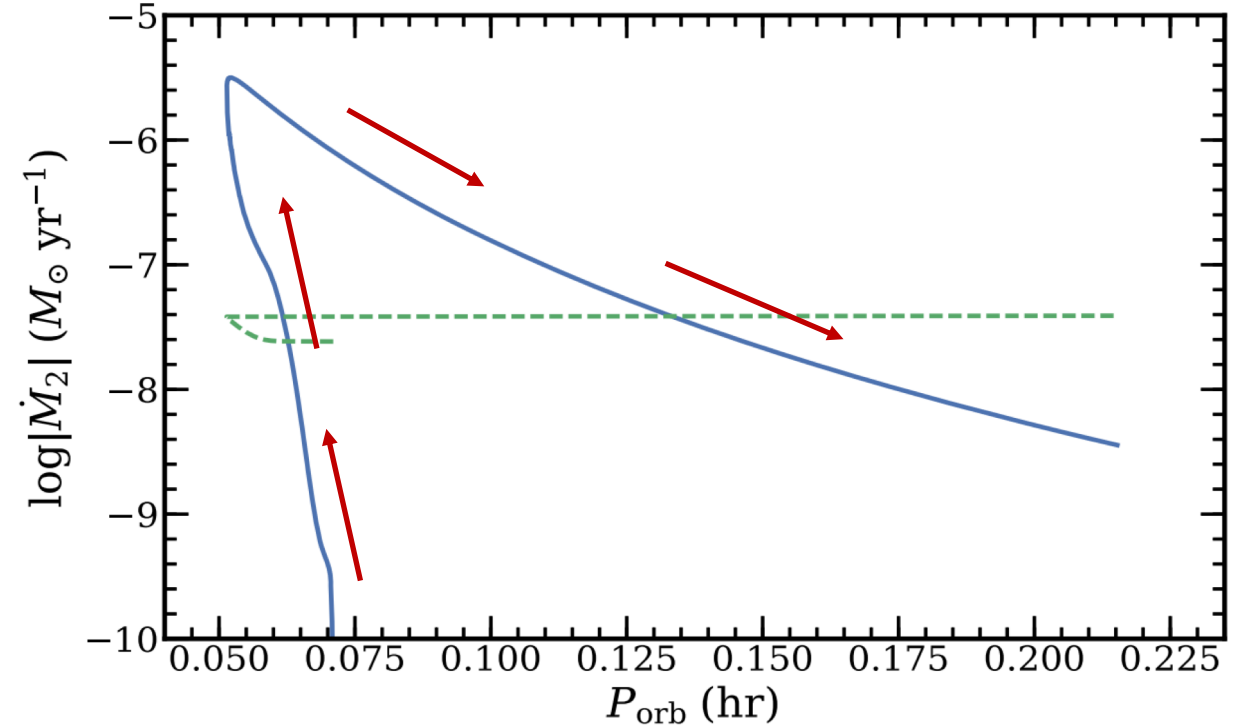
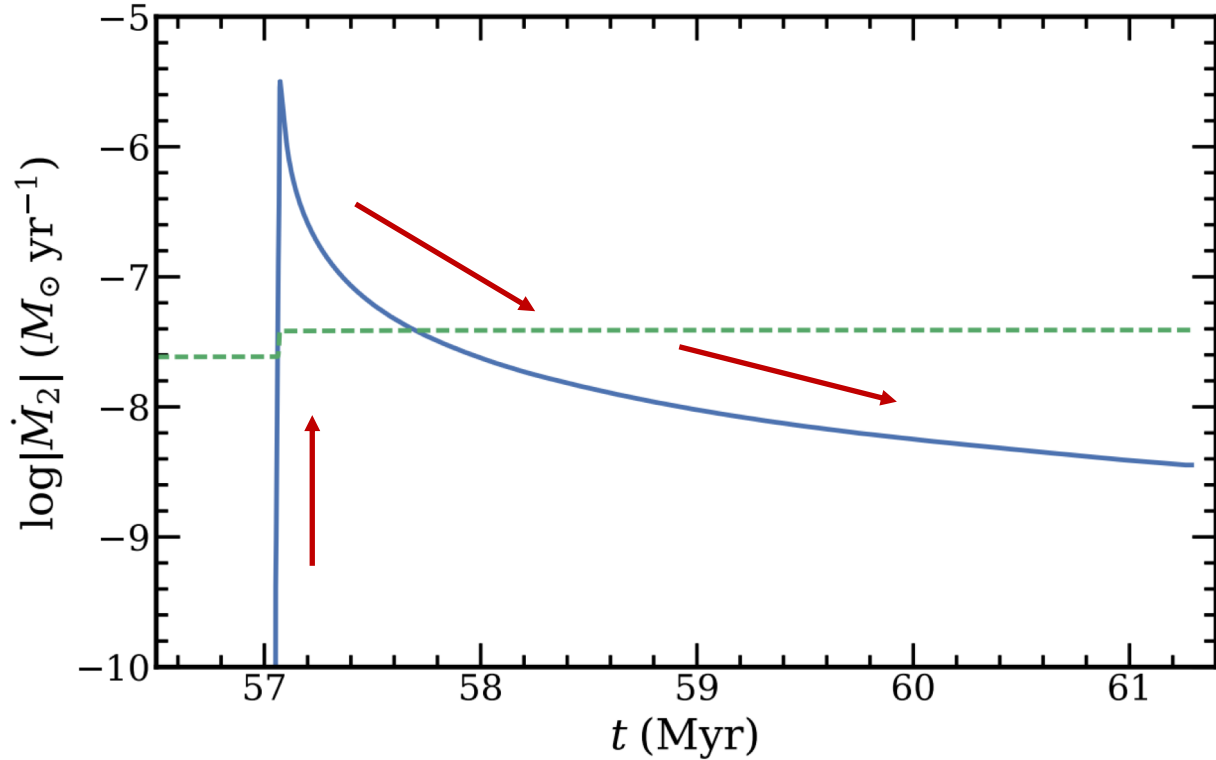
With the MB from Van et al. 2019, the parameter space for producing MSP-ELM WD and UCXBs becomes larger.

Evolution of NS+He WD binaries



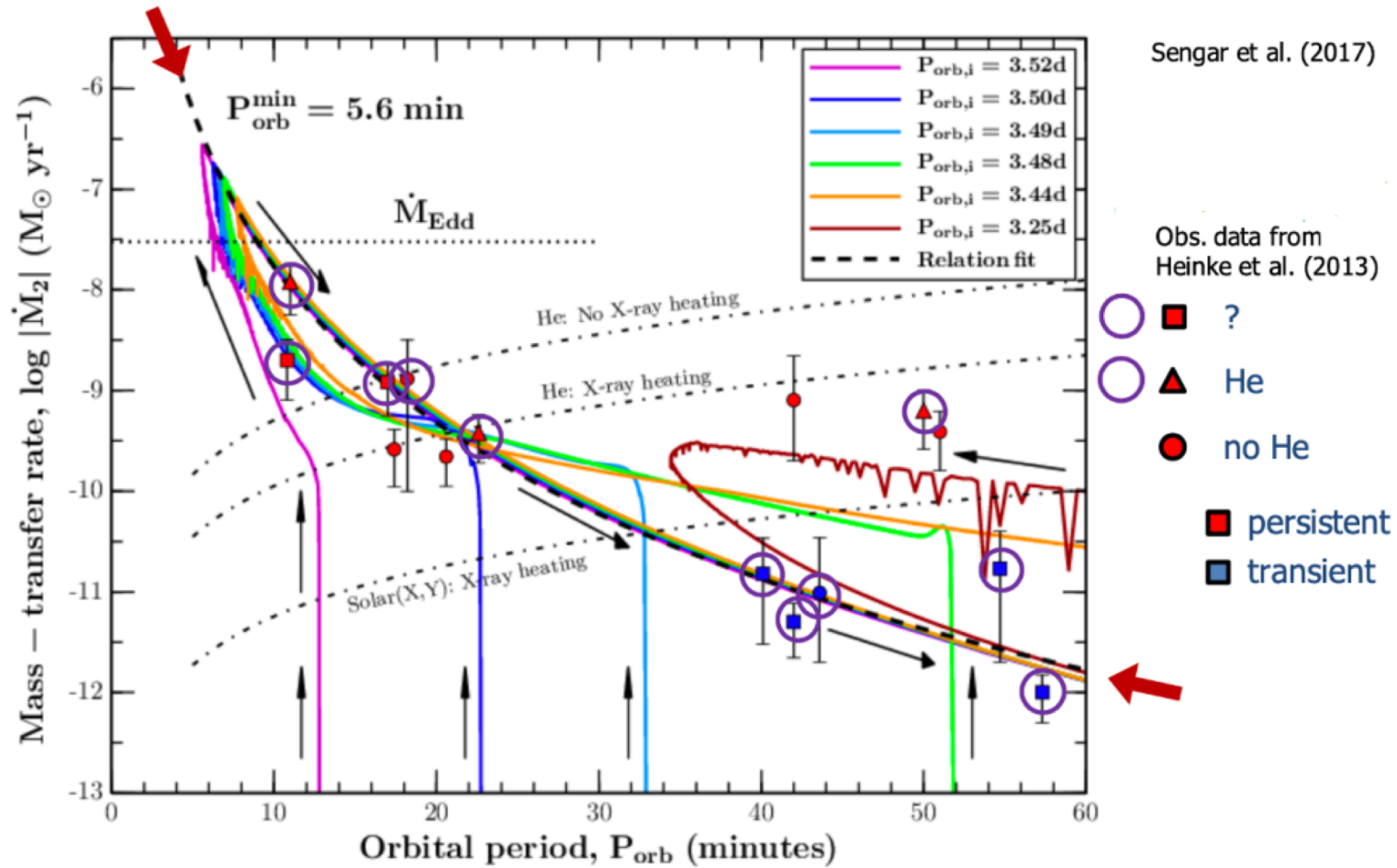
Evolution of NS+He WD binaries

Chen et al. ApJ, 2022



$M_{\text{NS}} = 1.30 M_{\text{sun}}$, $M_2 = 0.25 M_{\text{sun}}$, $P_{\text{orb}} = 0.05$ days

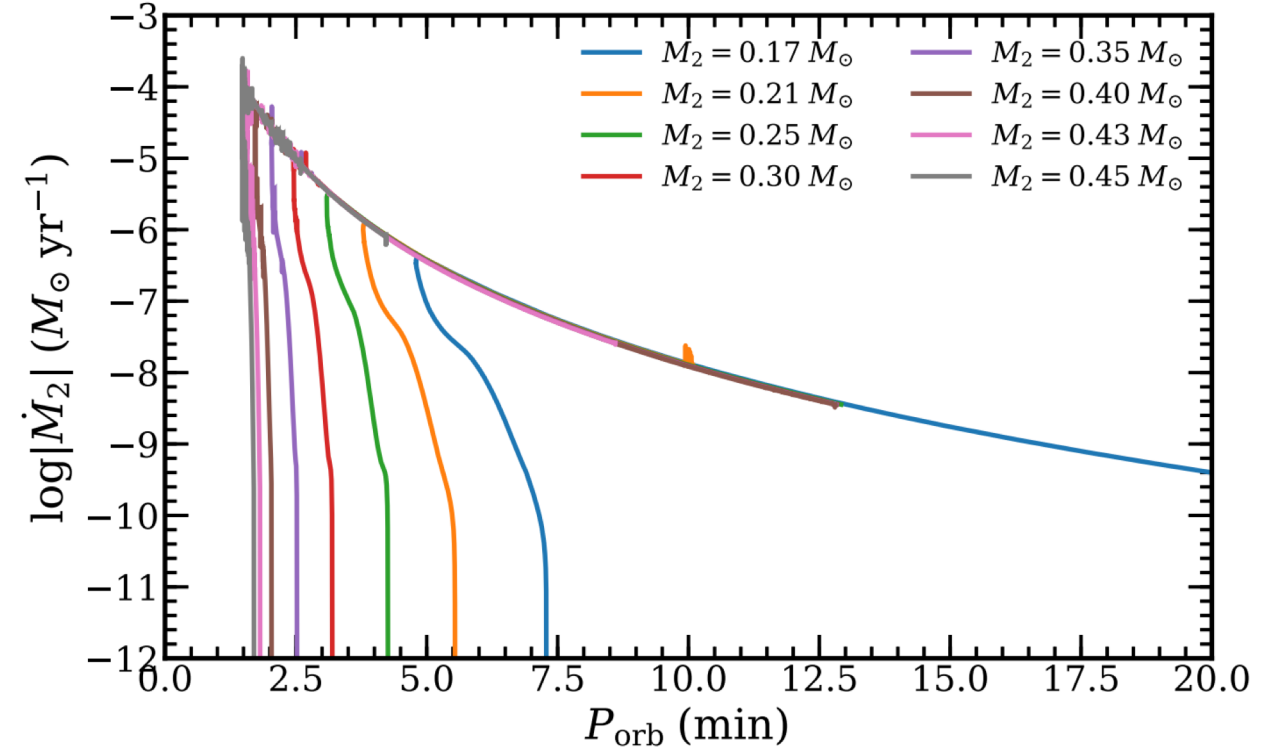
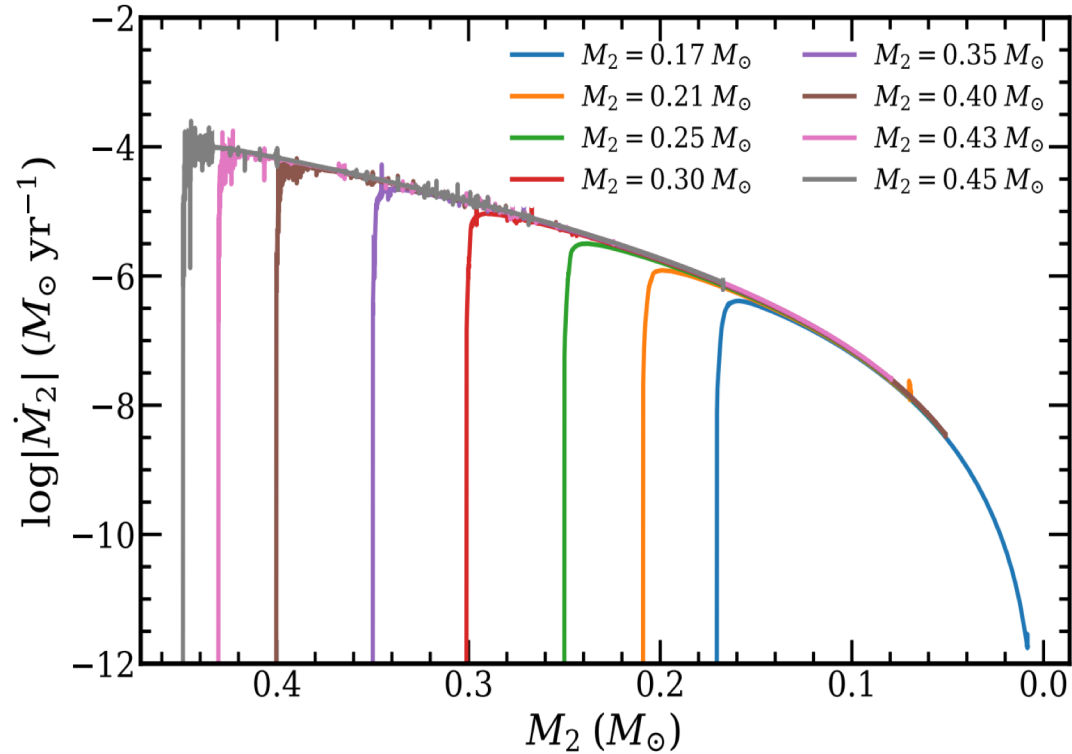
Evolution of NS+He WD binaries



Common declining branch: $\log |\dot{M}/M_\odot \text{ yr}^{-1}| = -5.15 \log(P_{orb}/\text{min}) - 2.62$

Stability of mass transfer of NS+He WDs

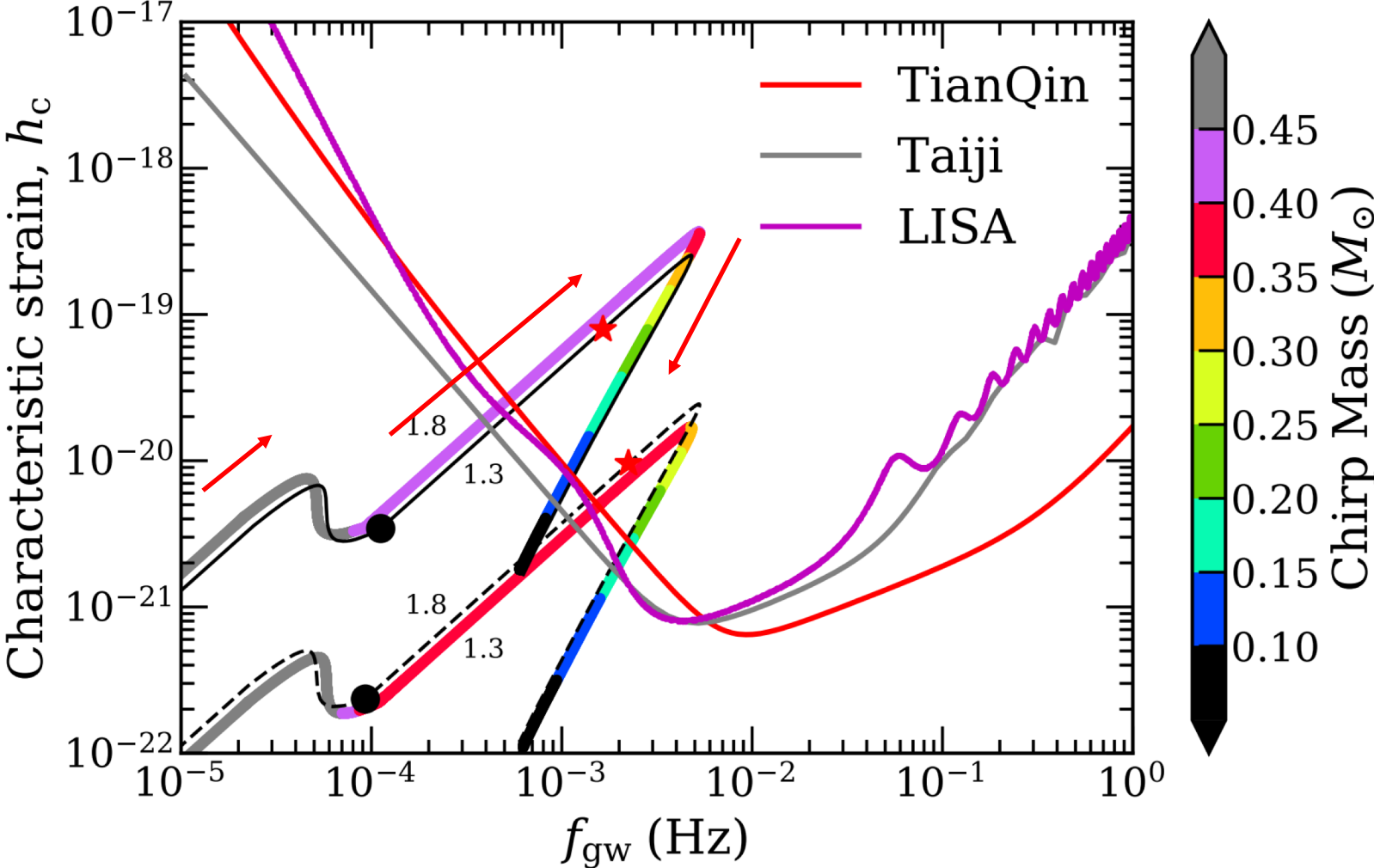
Chen et al. ApJ 2022



All NS+He WD binaries can have dynamically stable mass transfer.

minimum orbital period: 1.5-4.7min (GW frequency: 7.1-22mHz)

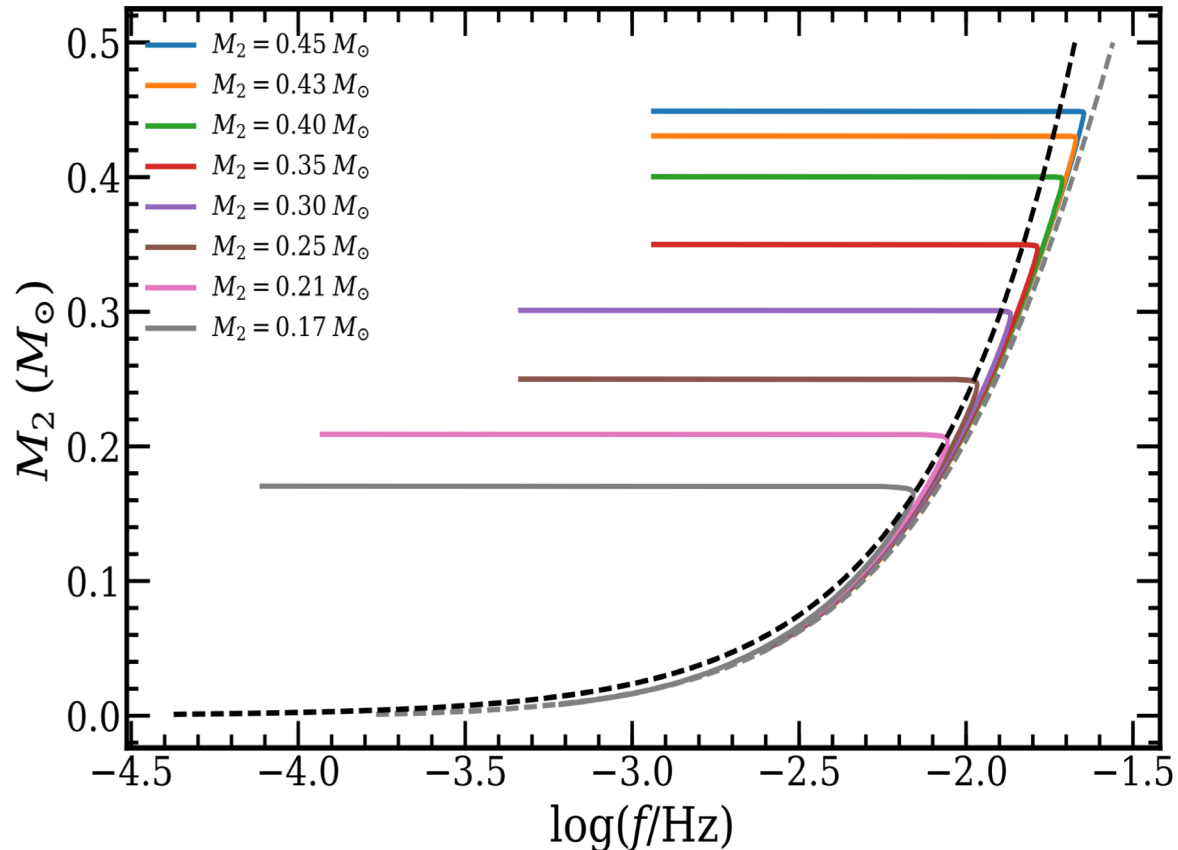
Properties as GW sources



UCXBs are important sources for TianQin, Taiji and LISA.

Properties as GW sources

Chen et al. ApJ, 2022



$$f = \frac{2}{P_{\text{orb}}} = \frac{1}{\pi} \sqrt{\frac{G(M_{\text{NS}} + M_2)}{a^3}}$$

$$\frac{R_{\text{rl}}}{a} \simeq 0.462 \left(\frac{M_2}{M_{\text{NS}} + M_2} \right)^{1/3}$$

$$R_2 = 0.013 R_\odot \left(\frac{M_2}{M_\odot} \right)^{-1/3}$$

$$R_2 = R_{\text{rl}}.$$

$$f = 4.23 \text{ mHz} \left(\frac{M_2}{0.1 M_\odot} \right)$$

The WD mass can be derived with this relation if the GW frequency is measured.

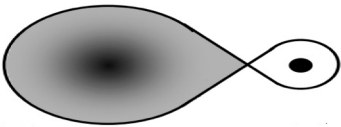
Formation of BH-UCXBs via accretion induced collapse of NSs

Chen et al. 2023, ApJ, accepted

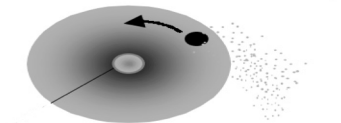
Classical formation scenario

ZAMS 

Problem:

RLO 

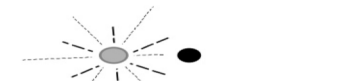
The progenitor of BH is a massive star (> 25 Msun) and the donor is a low mass star (1 – 4 Msun).

CE 

It is hard to eject the envelope.

He + MS 

CE evolution:

SN 

$$\frac{m_1 m_{1,env}}{\lambda R_1} = \alpha_{CE} \left(-\frac{Gm_1 m_2}{2a_i} + \frac{Gm_{1,c} m_2}{2a_f} \right)$$

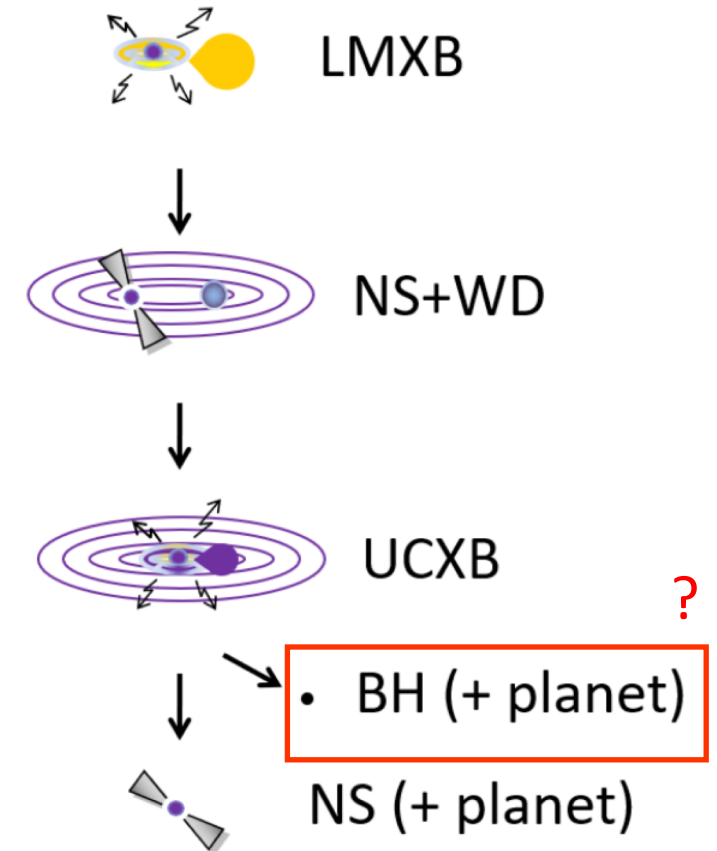
BH + MS 

Binding energy

orbital energy

BH + He/WD 

Credit: Thomas Tauris



Maximum mass of NSs

NS mass from pulsar observations

PSR J1614-2230 1.908 ± 0.016 Msun (Demorest et al. 2010, Fonseca et al. 2016)

PSR J0348+0432 2.01 ± 0.04 Msun (Antoniadis et al. 2013)

PSR 0740+6620 $2.14(+0.10/-0.09)$ Msun (Cromartie et al. 2020)

PSR 2215+5135 2.27 ± 0.17 Msun (Linares et al. 2018)

PSR 0952-0607 2.35 ± 0.17 Msun (Romani et al. 2022)

Constrain from GW 170817: Maximum mass $2.16 (+0.17/-0.15)$ Msun (Rezzolla et al. 2018)

The maximum mass of the NSs should be around 2.20 Msun.

**During the evolution phase of LMXBs, little mass is accreted by NS.
Some NSs must be born with masses close to the maximum mass.**

Kinematic effects of the AIC event

- During the collapse, we assume that an equivalent mass of 0.20 Msun is lost.
- The kick velocity of BH is assumed to be < 50 km/s.
- The binary separation after AIC

Tauris et al. 2013

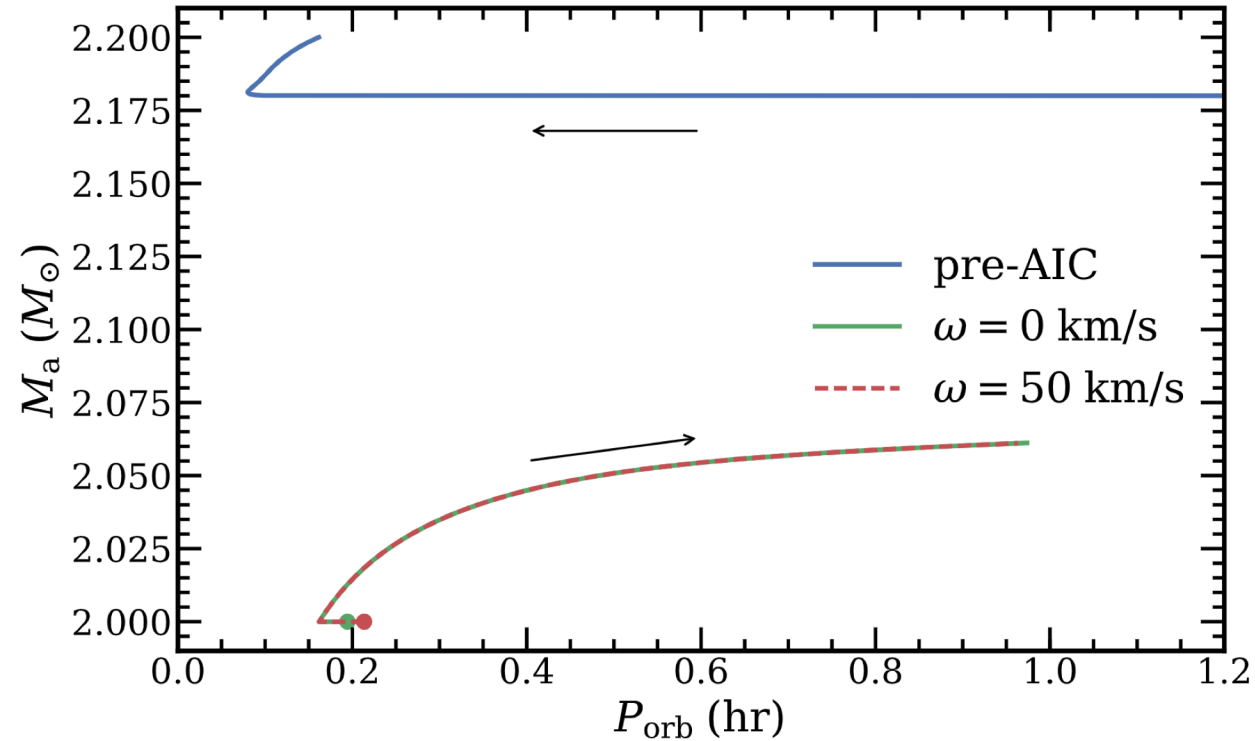
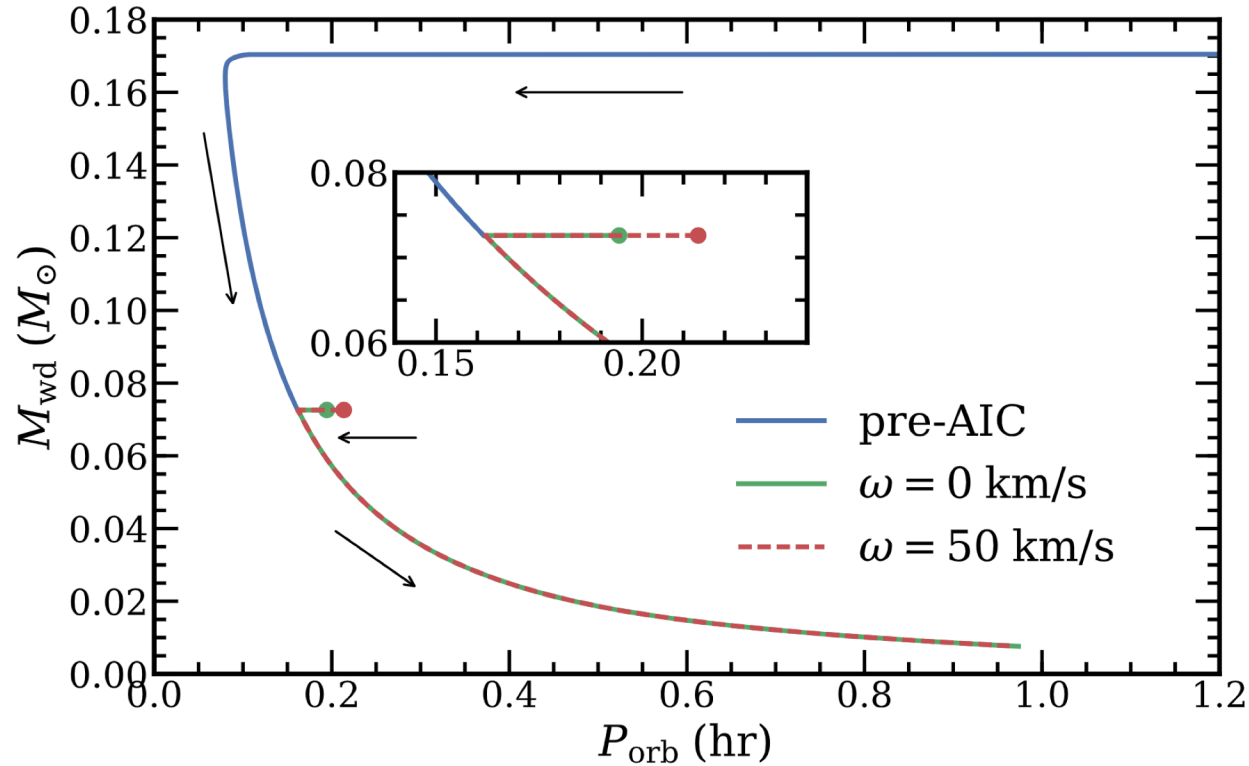
$$\frac{a}{a_0} = \left[\frac{1 - (\Delta M/M_0)}{1 - 2(\Delta M/M_0) - (w/v_c)^2 - 2 \cos \theta (w/v_c)} \right]$$

- The eccentricity after AIC

$$e = \sqrt{1 + \frac{2E_{\text{orb}} J_{\text{orb}}^2}{\mu G^2 M_{\text{NS}}^2 M_2^2}}$$

- the post-AIC orbit will be circularized due to the tidal interaction, binary separation: $a(1-e^2)$

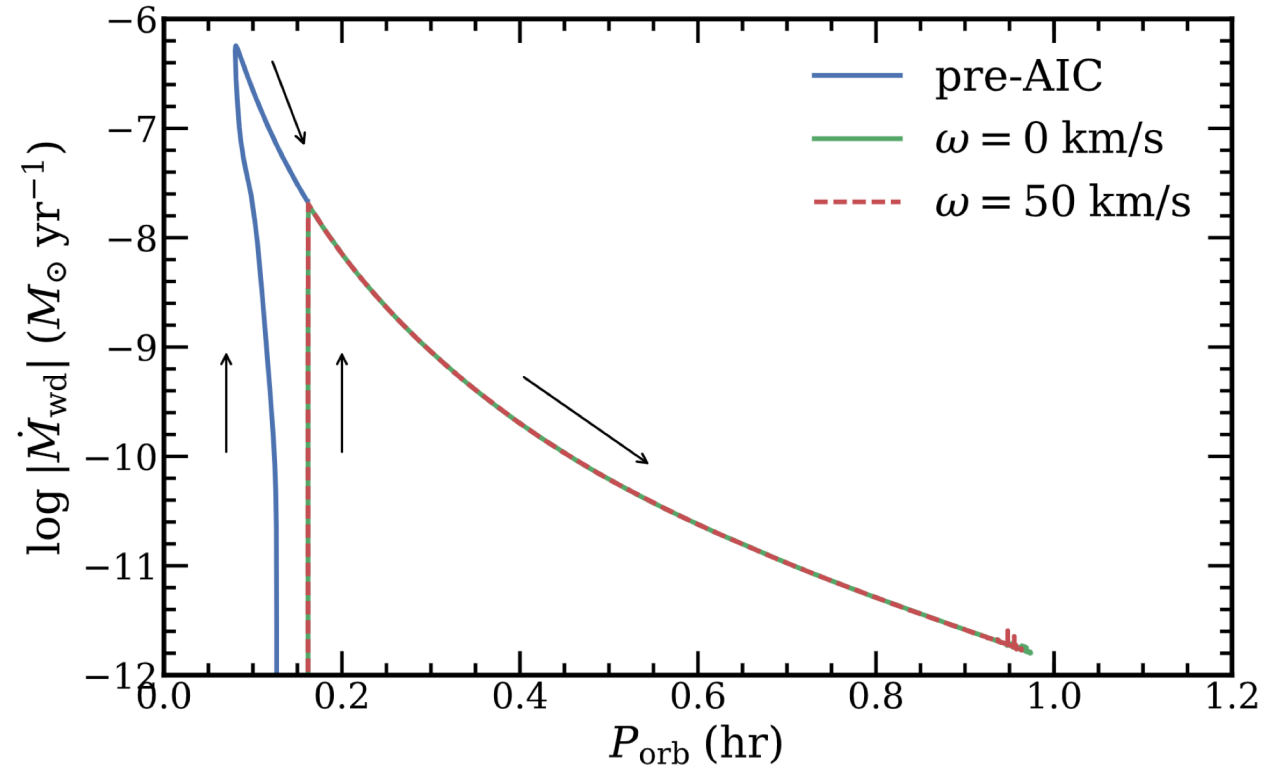
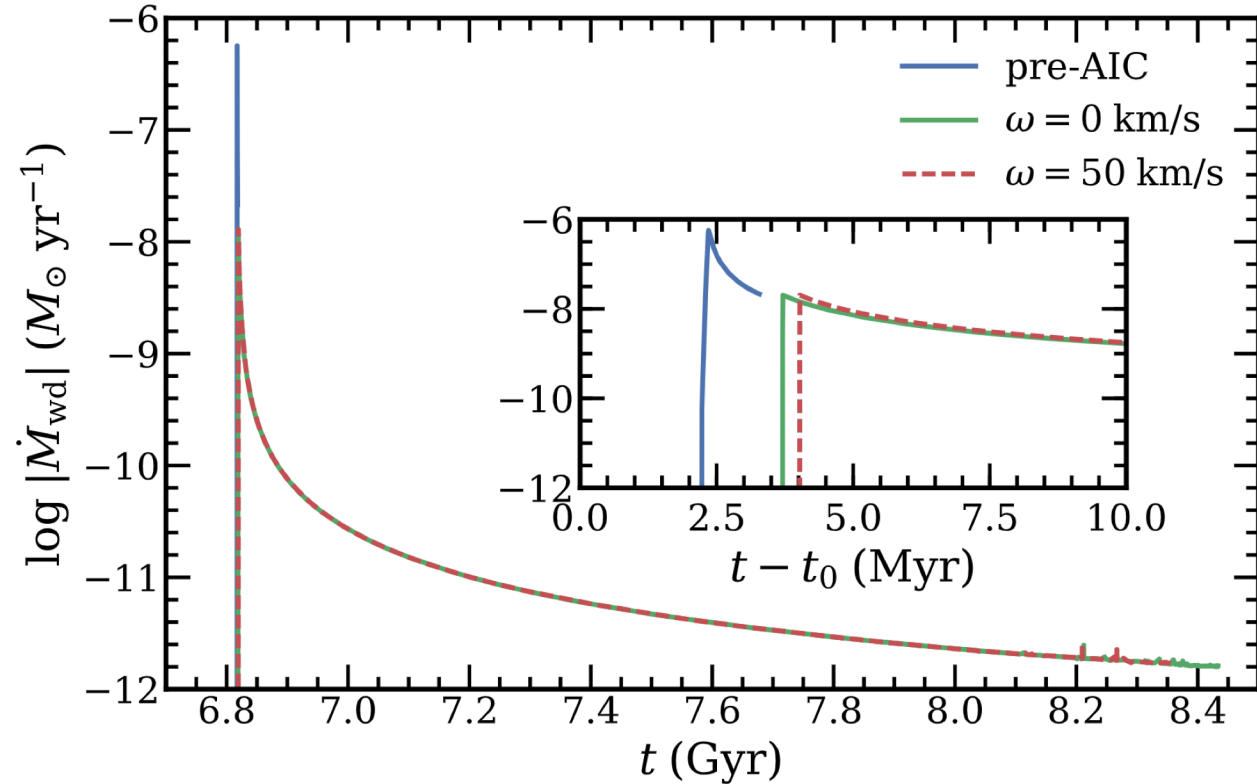
Binary evolution examples



Initial binary parameters:

$M_{\text{NS}} = 2.18 M_{\text{sun}}$, $M_{\text{wd}} = 0.17 M_{\text{sun}}$, $P_{\text{orb}} = 0.30$ days

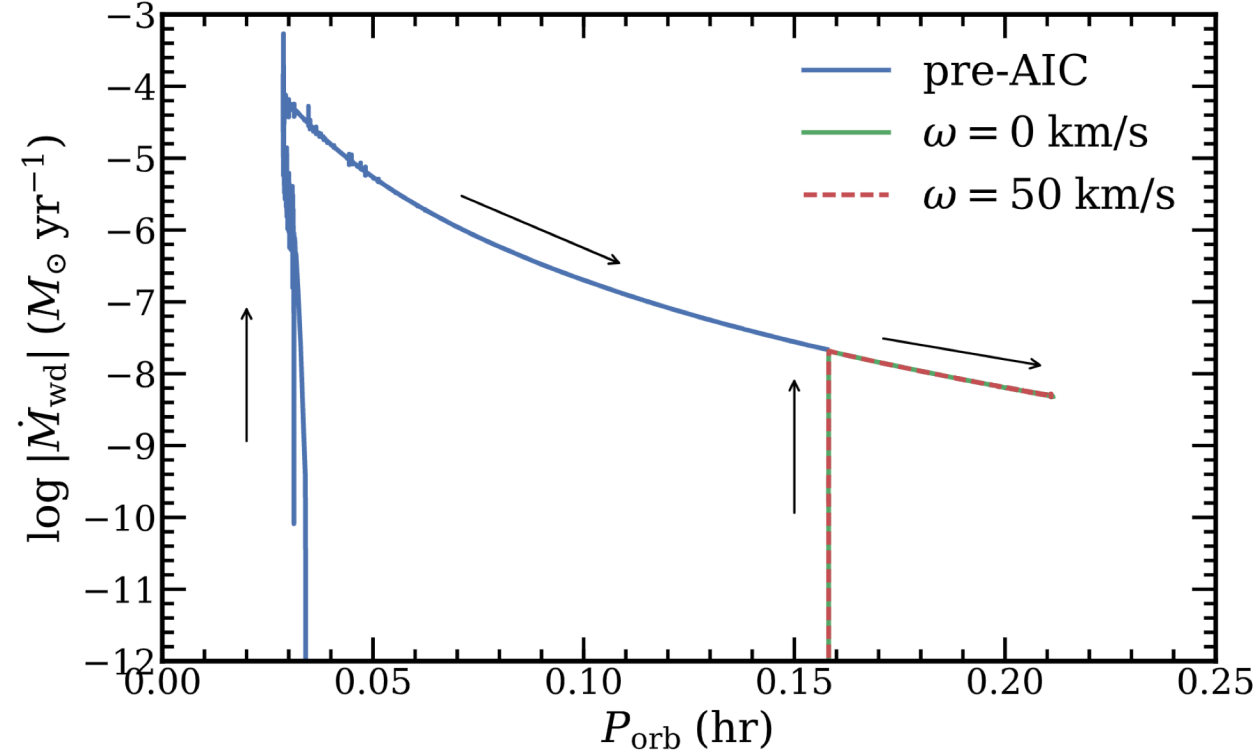
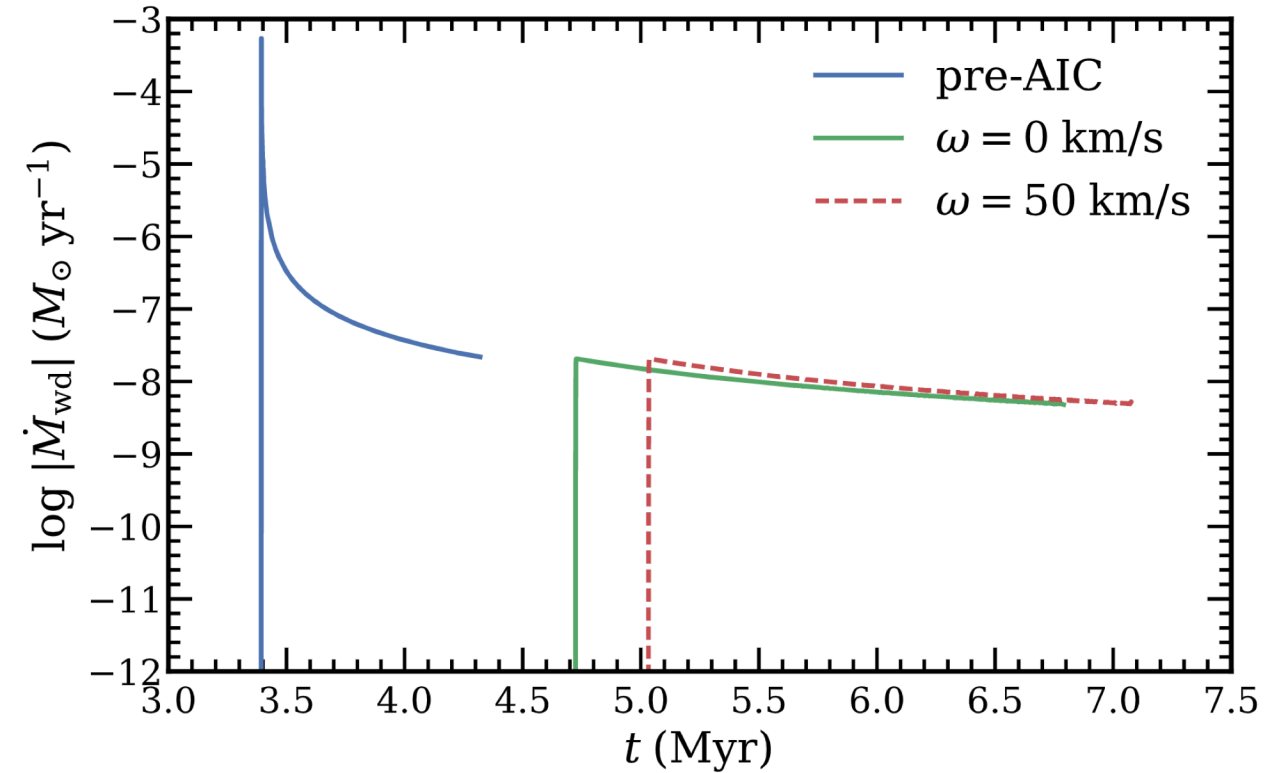
Binary evolution examples



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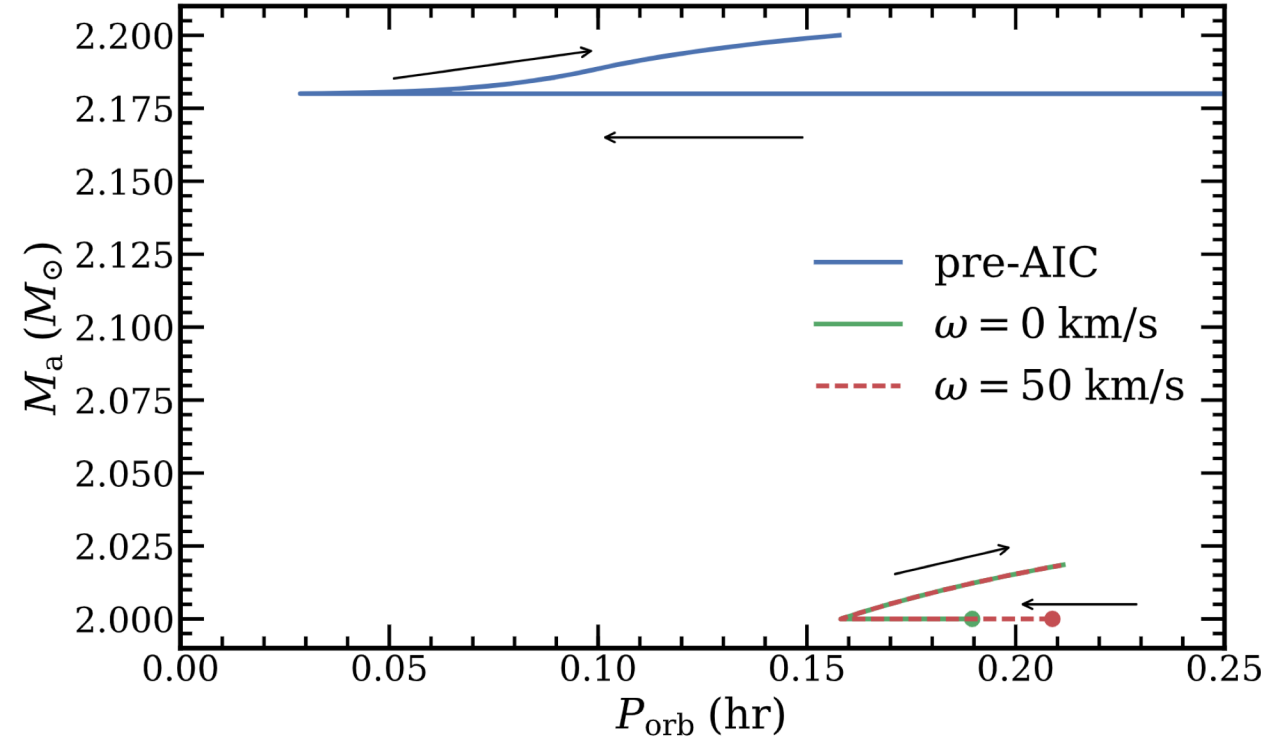
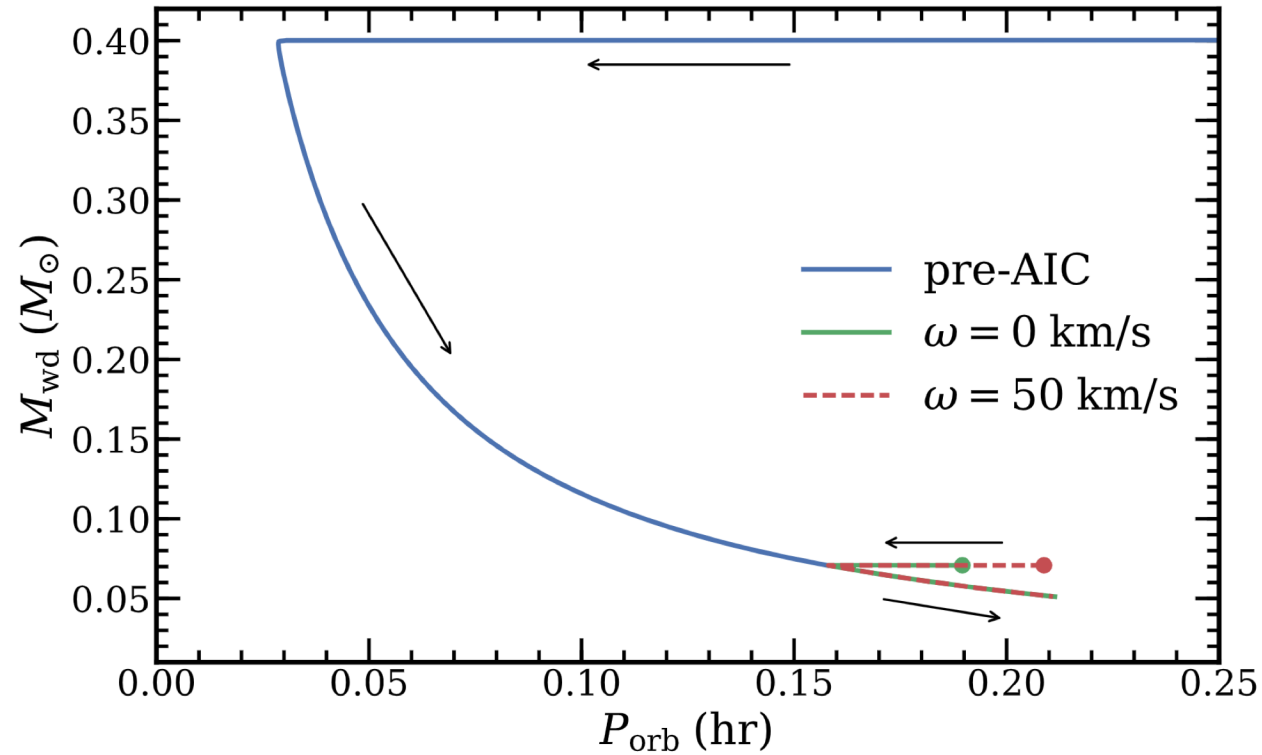
Binary evolution examples



Initial binary parameters:

$M_{\text{NS}} = 2.18 M_{\text{sun}}$, $M_{\text{wd}} = 0.40 M_{\text{sun}}$, $P_{\text{orb}} = 0.02 \text{ days}$

Binary evolution examples



Initial binary parameters:

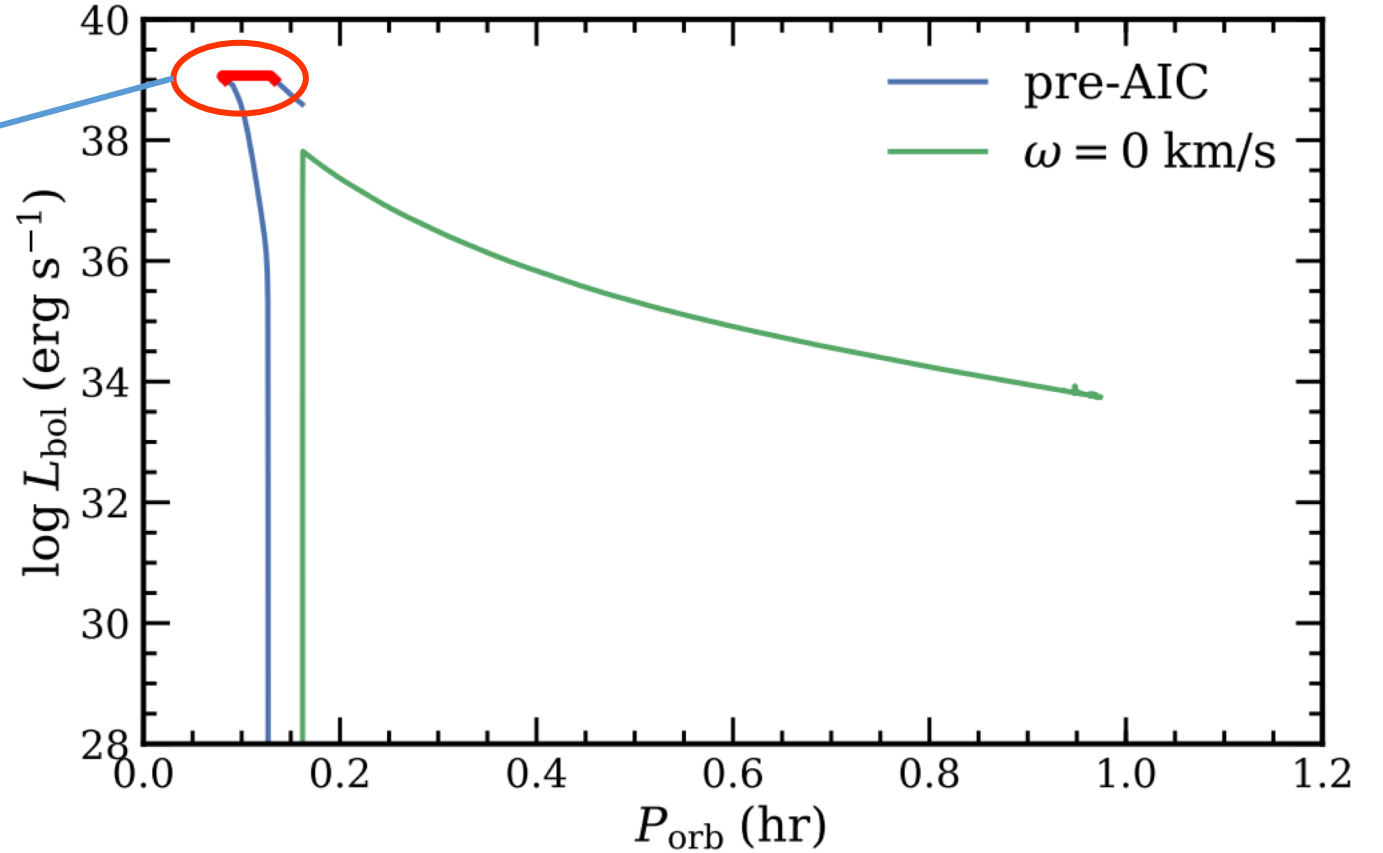
$M_{\text{NS}} = 2.18 M_{\text{sun}}$, $M_{\text{wd}} = 0.40 M_{\text{sun}}$, $P_{\text{orb}} = 0.02$ days

A new kind of ultra-luminous X-ray sources?

X-ray luminosity
 $L_x > 10^{39}$ erg/s

ultraluminous X-ray sources

This may be helpful to explain some of the ultraluminous X-ray sources observed in old stellar population environments (e.g. Irwin et al. 2004)



Summary

- With the MB prescription from Van et al. 2019, the fine-tuning problem is relieved.
- We found that all NS + He WD binaries can have dynamically stable mass transfer.
- The GW signal from NS + He WDs can be detected by LISA, TianQin and Taiji. There is a tight relation between the WD mass and GW frequency for UCXBs.
- We demonstrate that BH-UCXBs can be produced via accretion induced collapse of neutron stars.

Thank you !