**Ref: Y. Dong, Z. Hu, R. Xu and L. Shao, PRD 108,104039 (2023)**







#### **Moment of inertia for axisymmetric neutron stars in the standard model extension**

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**Supervisor: Lijing Shao**

- **• Moments of inertia (MOI) of neutron star (NS)** 
	- Build relations between **angular momentum** and **angular velocity**

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#### Credited: ScienceNews



- **• Moments of inertia (MOI) of neutron star (NS)** 
	- Build relations between **angular momentum** and **angular velocity**
- **• For Theories** 
	- **• Structure of NS:** MOI differs under different equation of state (EOS)
	- **• Universal relations (I-Love-Q)**: Insensitivity to EOS to help us learn gravity



J. B. Hartle, 1967, ApJ J. B. Hartle & K. S. Thorne, 1968 ApJ Credited: K. Yagi & N. Yunes, 2013, PRD



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- **• For observations**
	- **• High precision pulsar timing: direct detection** of MOI from Lense-Thirring effect in pulsar timing H. Hu et al., 2020, MNRAS



Credited: M. Kramer et al., 2022, PRX



- **• For observations**
	- **• High precision pulsar timing: direct detection** of MOI from Lense-Thirring effect in pulsar timing H. Hu et al., 2020, MNRAS
	- **• Glitch**: timing irregularities in pulsar timing observation
		- **•** Changes in MOI will result in variations in angular velocity



Credited: C. M. Espinoza et al., 2010, MNRAS



N. Andersson et al., 2003, PRL X. Y. Lai et al., 2023, MNRAS

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### **Calculation of moment of inertia**

- **• Moments of inertia of spherically symmetric neutron stars in GR** 
	- Solve the perturbed **Einstein's field equation**
	- For same density distribution, Newton or GR give **different results**

J. B. Hartle, 1967, ApJ J. B. Hartle & K. S. Thorne, 1968 ApJ





## **Calculation of moment of inertia**

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**A consistent approach to calculate**

**MOI of axisymmetric NS numerically**

## **Standard-Model Extension (SME)**

• Effective field theory that involve operators for Lorentz violation

• The neutron stars have a preferred direction



 $\frac{1}{16\pi G} \left( -uR + s^{\mu\nu}R_{\mu\nu}^T + t^{\alpha\beta\gamma\delta}C_{\alpha\beta\gamma\delta} \right)$ 

$$
\mathscr{L}_{LV}^{(4)} = \frac{1}{16\pi G} \left( -\right.
$$

Q. G. Bailey & V. A. Kostelecky , 2006, PRD





Credited: R. Xu et al., 2020, PLB



# Calculations

# MOI for spherical NSs in GR

Results

# MOI for axisymmetric NSs in SME

#### **Moment of inertia in GR**

• Stationary and axially symmetric metric,

 $ds^2 = -H^2dt^2 + Q^2dr^2 + r^2K^2 \left[d\theta^2 + \sin^2\theta(d\varphi - Ldt)^2\right]$ 

• Field equation,

#### 4-velocity Slowly rotation,  $u^r = u^\theta = 0$ ,  $u^\varphi = \Omega u^t$  $L(r, \theta) = \omega(r, \theta) + \mathcal{O}(\Omega^3)$

$$
R^t_\rho=8\pi T^t_\rho
$$

 $\rfloor$ 

J. B. Hartle, 1967, ApJ J. B. Hartle & K. S. Thorne, 1968 ApJ





#### **Moment of inertia in GR**

• Stationary and axially symmetric metric,

• Field equation,

• Further (partial differential equation, PDE),

$$
ds^2 = -H^2dt^2 + Q^2dr^2 + r^2K^2\left[d\theta^2 + \sin^2\theta(d\varphi - Ldt)^2\right]
$$

#### 4-velocity Slowly rotation, Frame dragging,  $u^r = u^\theta = 0$ ,  $u^\varphi = \Omega u^t$  $L(r, \theta) = \omega(r, \theta) + \mathcal{O}(\Omega^3)$  $\bar{\omega} = \Omega - \omega(r, \theta)$

 $\partial$  $\frac{\partial}{\partial \theta}$  (sin<sup>3</sup> $\theta$ ∂*ω*¯  $\left(\frac{\partial}{\partial \theta}\right) = 0$ 

 $\rfloor$ 

$$
R^t_\varphi=8\pi T^t_\varphi
$$

$$
\frac{1}{r^4}\frac{\partial}{\partial r}\left(r^4j\frac{\partial\bar{\omega}}{\partial r}\right) + \frac{4}{r}\frac{dj}{dr}\bar{\omega} + \frac{e^{(\lambda-\nu)/2}}{r^2}\frac{1}{\sin^3\theta}\frac{\partial}{\partial r}
$$

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• Stationary and axially symmetric metric,

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$$

$$
R^t_\varphi=8\pi T^t_\varphi
$$

$$
\frac{1}{r^4} \frac{\partial}{\partial r} \left( r^4 j \frac{\partial \bar{\omega}}{\partial r} \right) + \frac{4}{r} \frac{dj}{dr} \bar{\omega} + \frac{e^{(\lambda - \nu)/2}}{r^2} \frac{1}{\sin^3 \theta} \frac{\partial}{\partial \theta} \left( \sin^3 \theta \right)
$$

• For spherical NS in GR,

 $\rfloor$ ∂*ω*¯  $\left(\frac{\partial}{\partial \theta}\right) = 0$ 4-velocity Slowly rotation, Frame dragging, Spherical NS in GR,  $u^r = u^\theta = 0$ ,  $u^\varphi = \Omega u^t$  $L(r, \theta) = \omega(r, \theta) + \mathcal{O}(\Omega^3)$  $\bar{\omega} = \Omega - \omega(r, \theta)$  $\bar{\omega}(r) = \Omega - \frac{2J}{r^2}$ *r*3 *I* = *J*/Ω

$$
\frac{1}{r^4}\frac{d}{dr}\left(r^4j\frac{d\bar{\omega}}{dr}\right) + \frac{4}{r}\frac{dj}{dr}\bar{\omega} = 0
$$

J. B. Hartle, 1967, ApJ J. B. Hartle & K. S. Thorne, 1968 ApJ





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#### **Moment of inertia in SME**

- We still want to calculation this field equation, but there's some difference
	- **• Difference 1: field equation has another term** 
		- We should solve  $R^t_\varphi = 8\pi T^t_\varphi V^t_\varphi$ , where  $V^\mu_\nu = \bar{s}^{\alpha\beta} g^{\mu\delta}$

*QCCCBailey et al., 2014, PRD* 



R. Xu et al., 2020, PLB





#### **Moment of inertia in SME**

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	- **• Difference 2: axisymmetric deformation of neutron stars**

$$
\int \rho^{(1)}(r,\theta) = -\frac{1}{6}\bar{s}^{zz}(3\cos^2\theta - 1)r\rho^{(0)}(r)_{,r}
$$

$$
P^{(1)}(r,\theta) = -\frac{1}{6} \bar{s}^{zz} (3 \cos^2 \theta - 1) r P^{(0)}(r)_{,r}
$$

Q. G. Bailey et al., 2014, PRD



R. Xu et al., 2020, PLB





#### **Moment of inertia in SME**

- We still want to calculation this field equation, but there's some difference
	- **• Difference 1: field equation has another term** 
		- We should solve  $R^t_\varphi = 8\pi T^t_\varphi V^t_\varphi$ , where  $V^\mu_\nu = \bar{s}^{\alpha\beta} g^{\mu\delta} G_{\delta\alpha\beta\nu}$
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$$

$$
\bullet \ P^{(1)}(r,\theta) = -\frac{1}{6} \bar{s}^{zz} (3 \cos^2 \theta - 1) r P^{(0)}(r)_{,r}
$$

Q. G. Bailey et al., 2014, PRD



R. Xu et al., 2020, PLB

$$
\frac{1}{r^4} \frac{\partial}{\partial r} \left( r^4 j \frac{\partial \bar{\omega}}{\partial r} \right) + \frac{4}{r} \frac{dj}{dr} \bar{\omega} + \frac{e^{(\lambda - \nu)/2}}{r^2} \frac{1}{\sin^3 \theta} \frac{\partial}{\partial \theta} \left( \sin^3 \theta \frac{\partial \bar{\omega}}{\partial \theta} \right) = S'_1(r, \theta) + S'_2(r, \theta)
$$





- **•** Far field behaviors of *ω*¯
- **•** Solve PDE for *r* ∈ (0,∞)
- **• Finite element method** 
	- Numerically, integrating r to infinity is impossible

**•** Change of variables  $(r, \theta) \rightarrow \left(x \equiv$  $r\cos\theta$ 

**• How to get MOI?** 

#### **Solve PDE numerically** 1  $\partial$ ∂*ω*¯ 4 d*j*

**•** Improve accuracy greatly

*r*4

 $\frac{\partial}{\partial r}$  (*r*<sup>4</sup>*j* 

<sup>∂</sup>*<sup>r</sup>* ) <sup>+</sup>

*r*

d*r*

 $\bar{\omega}$  +

$$
\frac{e^{(\lambda-\nu)/2}}{r^2} \frac{1}{\sin^3\theta} \frac{\partial}{\partial\theta} \left(\sin^3\theta \frac{\partial\bar{\omega}}{\partial\theta}\right) = S'_1(r,\theta) + S'_2(r,\theta)
$$



#### **Solve PDE numerically**

- **• Process of numerical solving** 
	- Derive the modified PDE of MOI
	- **•** Change of variables

$$
(r, \theta) \to \left(x \equiv \frac{r \cos \theta}{r + R}, y \equiv \frac{r \sin \theta}{r + R}\right)
$$

- Solve  $(x, y)$  PDE to get numerical solution
- **•** Fitting with polynomial to get MOI with far field behaviors of  $\bar{\omega}$

$$
\bar{\omega}(r,\theta) = A + \frac{B(\theta)}{r^3} + \frac{C(\theta)}{r^4} + \dots
$$



#### **Results**

- **•** Corrections to MOIs **caused by axisymmetric deformations**
	- $k = \delta I_1/I$ , for chosen deformations it is 1/3
	- For deformation caused by Lorentz violation,
		- **Ratio**  $k = \delta I_1/I$  from GR differs from **Newtonian one**
		- For a neutron star with  $1.4 M_{\odot}$ , it is  $8 \%$





#### **Summary**

- We develop a **consistent approach** to calculate moment of inertia for **axisymmetric** neutron stars in the Standard-Model Extension
	- **• Finite element method**
	- **Ratio**  $k = \delta I_1/I$  from GR differs from Newtonian **one** 
		- For a neutron star with  $1.4 M_{\odot}$ , it is  $8 \%$





