

Neutron Stars in Bumblebee Theory

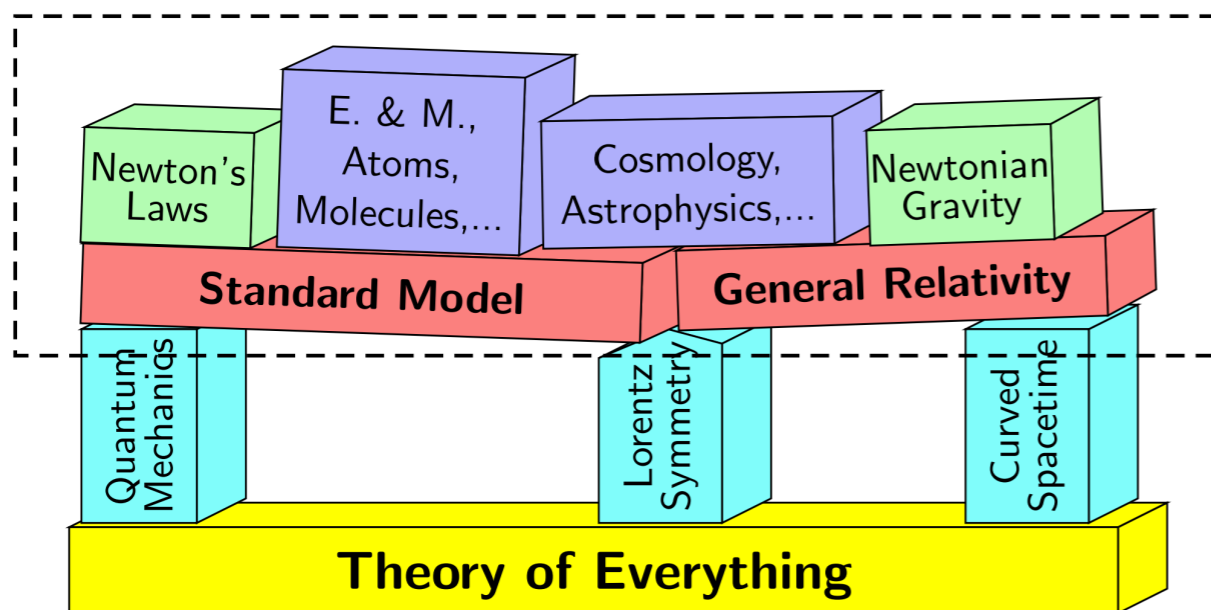
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Bumblebee Theory

$$S = \int d^4x \sqrt{-g} \left(\frac{R}{2\kappa} + \frac{\xi}{2\kappa} B^\mu B^\nu R_{\mu\nu} - \frac{1}{4} B^{\mu\nu} \underbrace{B_{\mu\nu}}_{B_{\mu\nu} = \partial_\mu B_\nu - \partial_\nu B_\mu} - \underbrace{V}_{V = V(B^\mu B_\mu - b^2)} \right) + S_m[g_{\mu\nu}, \Psi_m] \quad [\text{V. Kostelecký (2004)}]$$

SME = general description of LV and low energies



$$\mathcal{L}_{\text{SME}} = \mathcal{L}_{\text{GR}} + \mathcal{L}_{\text{SM}} + \mathcal{L}_{\text{LV}}$$

$$\mathcal{L}_{\text{LV}} = uR + s^{\mu\nu} R_{\mu\nu}^T + t^{\kappa\lambda\mu\nu} R_{\kappa\lambda\mu\nu}$$

$$u_B = \frac{1}{4} g^{\mu\nu} B_\mu B_\nu \quad s_B^{\mu\nu} = \xi B^\mu B^\nu - g^{\mu\nu} u_B$$

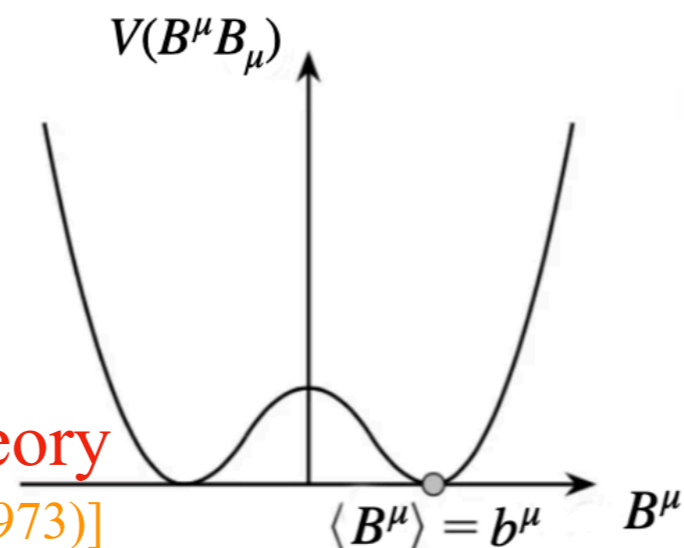
[c.f. M. Mewes]

$$V|_{B_\mu=b_\mu} = 0 \quad \left. \frac{dV}{d(B^\lambda B_\lambda)} \right|_{B_\mu=b_\mu} = 0$$

$$V = \frac{\lambda}{2} (B^\mu B_\mu - b^2)$$

Hellings-Nordtvedt theory

[R. Hellings & K. Nordtvedt (1973)]



EoM & Ansatz

Equations of motion:

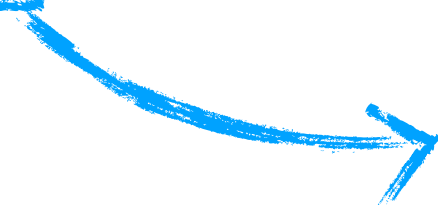
$$G_{\mu\nu} = \kappa(T_{\mu\nu}^m + T_{\mu\nu}^{\text{vec}} + T_{\mu\nu}^{\xi})$$

$$\nabla_{\nu} B^{\mu\nu} - \frac{\xi}{\kappa} R^{\mu\nu} B_{\nu} = 0$$

Static and spherical ansatz:

$$ds^2 = - e^{2\nu(r)} dt^2 + e^{2\mu(r)} dr^2 + r^2(d\theta^2 + \sin^2 \theta d\phi^2)$$

$$b_{\mu} = (b_t(r), b_r(r), 0, 0)$$


$$\left(1 - \frac{2Gm(r)}{r}\right)^{-1}$$

Two classes of solutions:

$$\text{Class I : } b_r = 0$$

$$\text{Class II : } R_{rr} = 0$$

Reduced geometric unit

$$\kappa = 1 \quad G = 1/8\pi$$

Integration

At center

$$m = \sum_{i=0}^{\infty} m_i r^i, \quad \nu = \sum_{i=0}^{\infty} \nu_i r^i, \quad b_t = \sum_{i=0}^{\infty} b_i r^i,$$

$$p = \sum_{i=0}^{\infty} p_i r^i, \quad \rho = \sum_{i=0}^{\infty} \rho_i r^i.$$

$$\nu_1 = 0 \quad b_1 = 0 \quad \rho_1 = p_1 = 0$$

Time transformation

$$t \rightarrow e^{-\Delta\nu} t$$

$$\mathbf{g} : g_{tt} \rightarrow e^{2\Delta\nu} g_{tt}$$

$$\mathbf{b} : b_t \rightarrow e^{\Delta\nu} b_b$$

System is invariant under

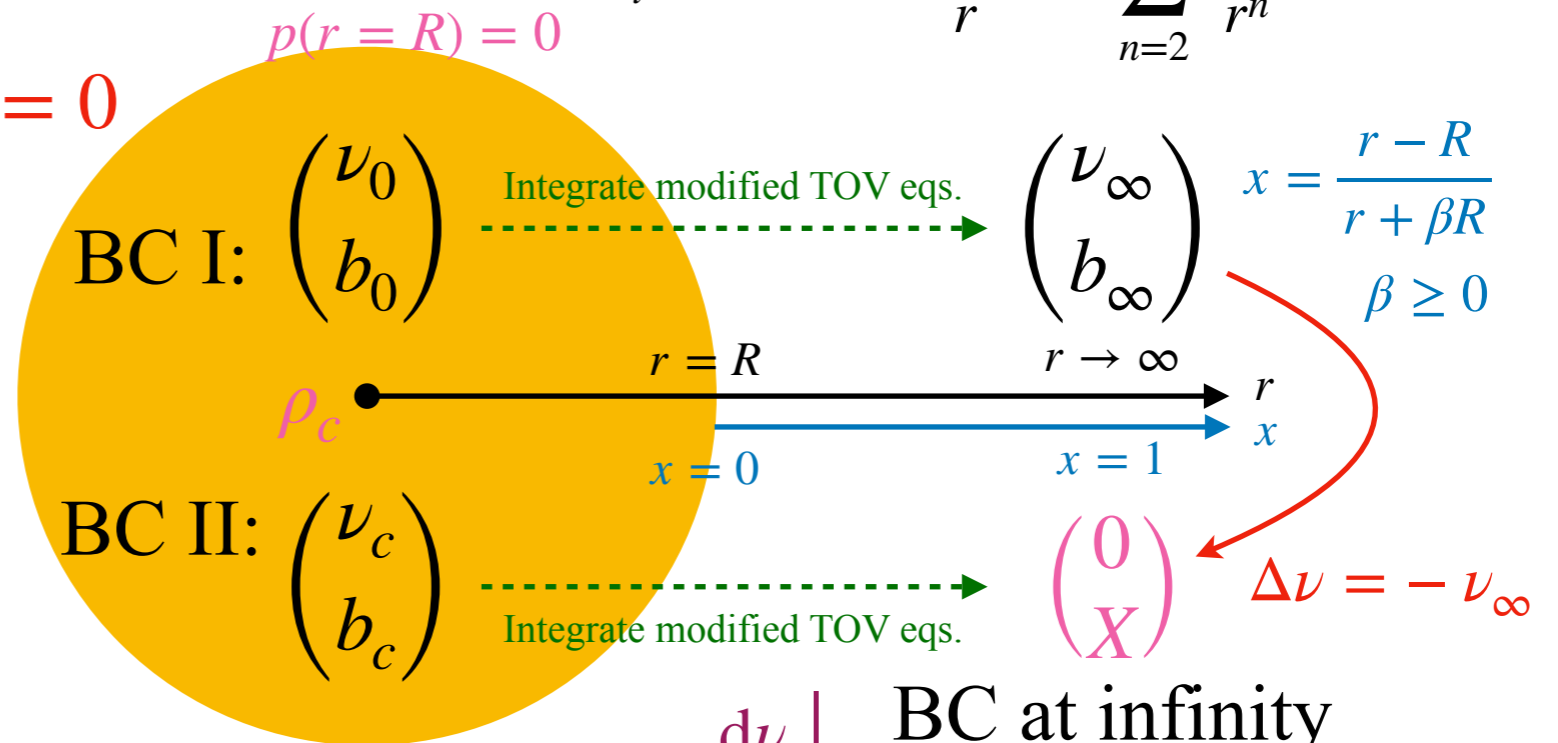
$$\nu \rightarrow \nu + \Delta\nu, \quad b_t \rightarrow e^{\Delta\nu} b_t$$

At infinity

$$\nu(r) = \sum_{n=1}^{\infty} \frac{\nu_{-n}}{r^n},$$

$$m(r) = M + \sum_{n=1}^{\infty} \frac{m_{-n}}{r^n},$$

$$b_t(r) = X + \frac{2\sqrt{\pi}Q}{r} + \sum_{n=2}^{\infty} \frac{b_{-n}}{r^n}.$$



$$M = 8\pi(1 + \beta)R \left. \frac{d\nu}{dx} \right|_{x=1}$$

$$Q = -\frac{1 + \beta}{2\sqrt{\pi}} R \left. \frac{db_t}{dx} \right|_{x=1}$$

$$M \rightarrow M$$

$$X \rightarrow e^{\Delta\nu} X$$

$$Q \rightarrow e^{\Delta\nu} Q$$

Back to GR

$$m_3 = \frac{4\pi}{3} \frac{3p_0 b_0^2 (2\kappa - \xi)\xi + 2e^{2\nu_0} \kappa \rho_0}{2e^{2\nu_0} \kappa + \xi(\xi - 2\kappa)b_0^2},$$

$$b_2 = \frac{\kappa}{6} \frac{e^{2\nu_0} (3p_0 + \rho_0)}{2e^{2\nu_0} \kappa + \xi(\xi - 2\kappa)b_0^2} \xi b_0,$$

$$\nu_2 = \frac{\kappa^2}{6} \frac{e^{2\nu_0} (3p_0 + \rho_0)}{2e^{2\nu_0} \kappa + \xi(\xi - 2\kappa)b_0^2},$$

$$p_2 = -\frac{\kappa^2}{6} \frac{e^{2\nu_0} (3p_0 + \rho_0)(p_0 + \rho_0)}{2e^{2\nu_0} \kappa + \xi(\xi - 2\kappa)b_0^2}.$$

$$b_0 = 0$$



$$\xi = 0, 2\kappa$$

$$m_3 = \frac{4\pi}{3} \rho_0$$

$$\nu_2 = \frac{2\pi}{3} G(3p_0 + \rho_0)$$

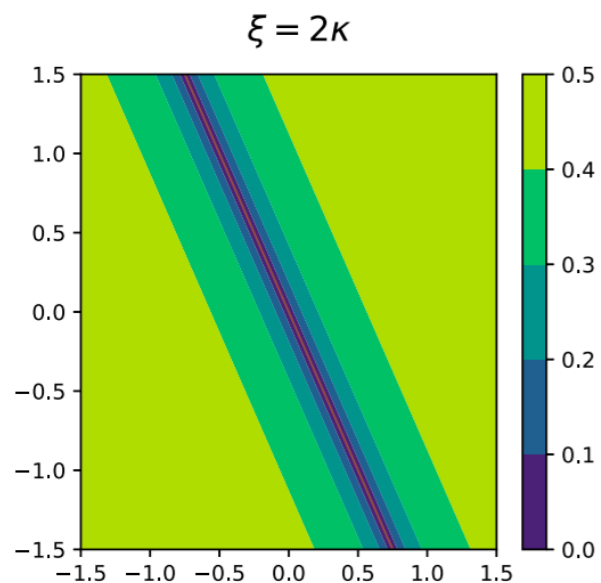
$$p_2 = -\frac{2\pi}{3} G(3p_0 + \rho_0)(p_0 + \rho_0)$$

$$b_2 = \frac{1}{6} (3p_0 + \rho_0) b_0$$

Einstein-Maxwell theory ($\xi = 0$)

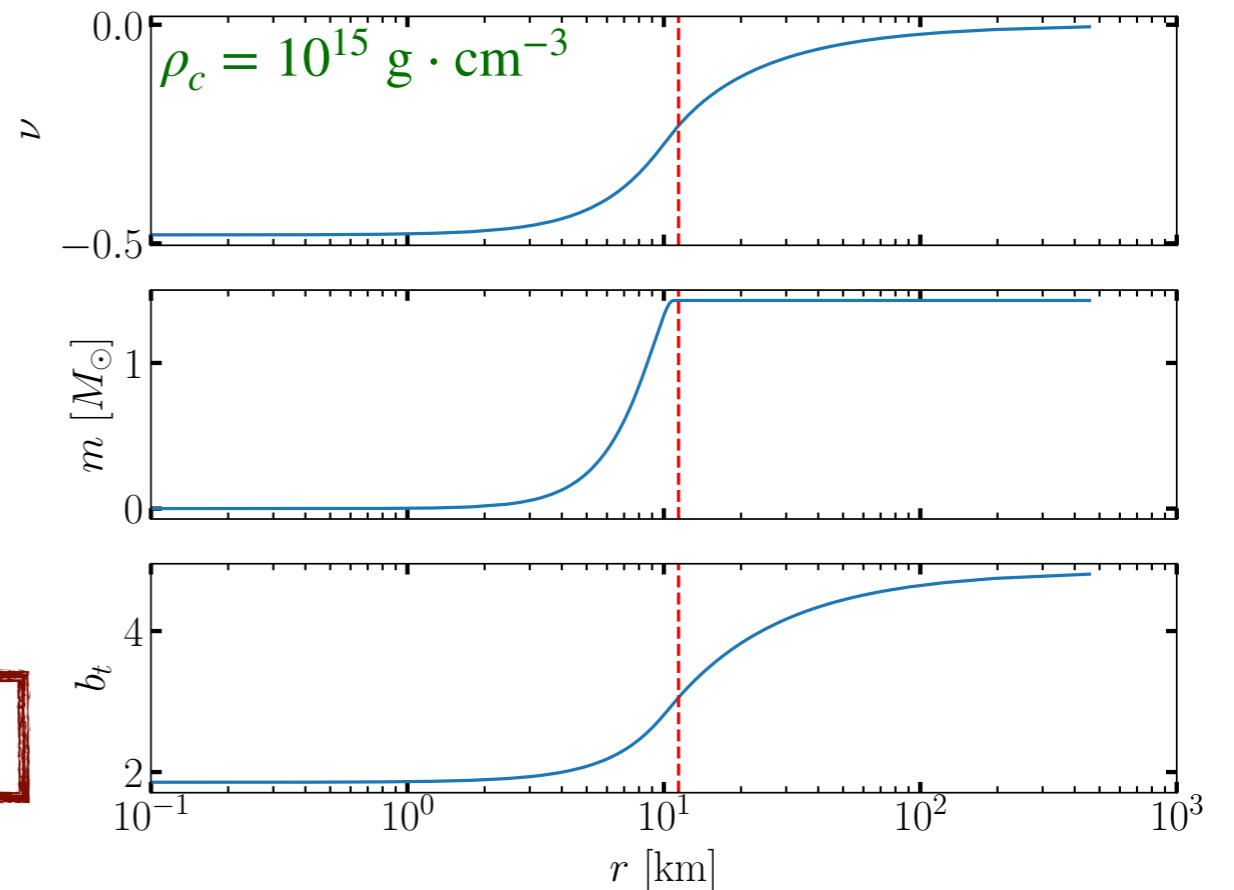
$$b_t(r) = b_0 + \int_0^r \frac{q_0}{r'^2} e^{\mu(r') + \nu(r')} dr'$$

Stealth Schwarzschild ($\xi = 2\kappa$)

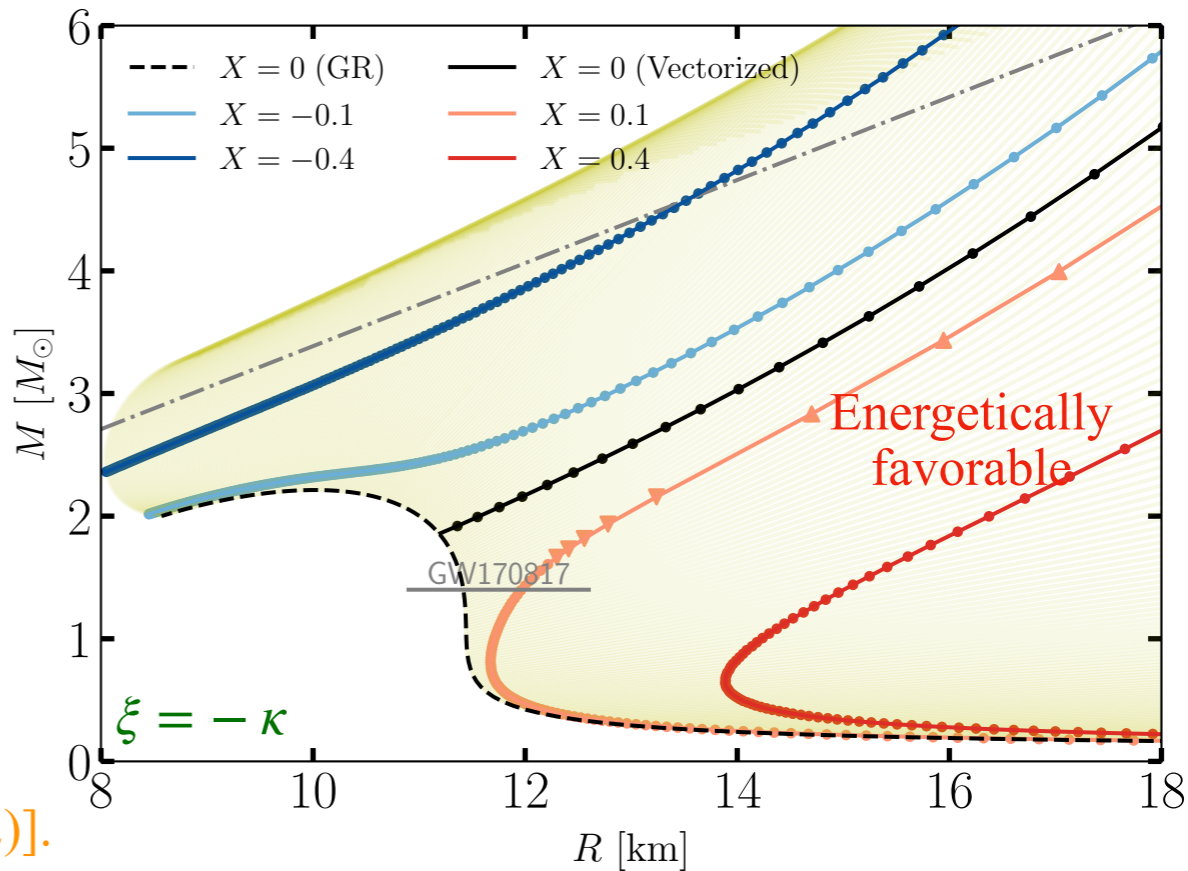
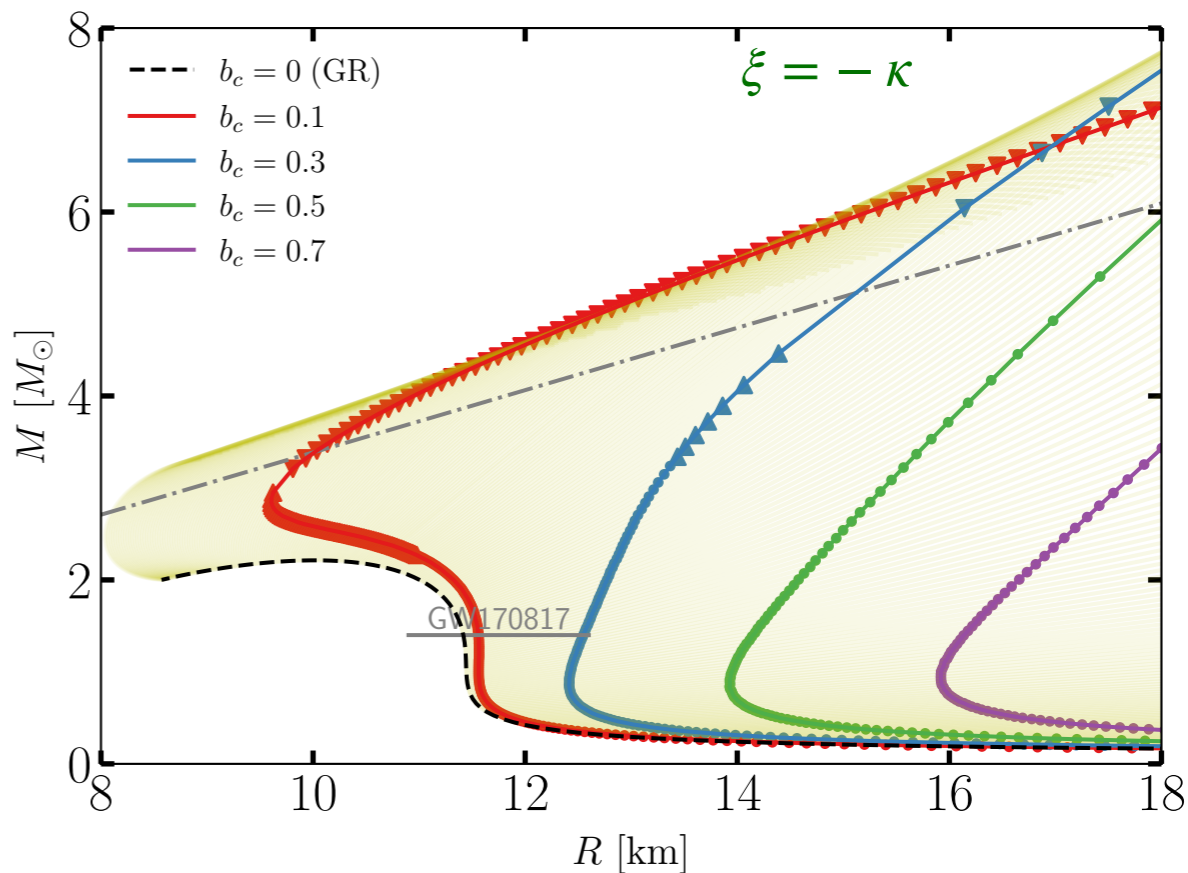
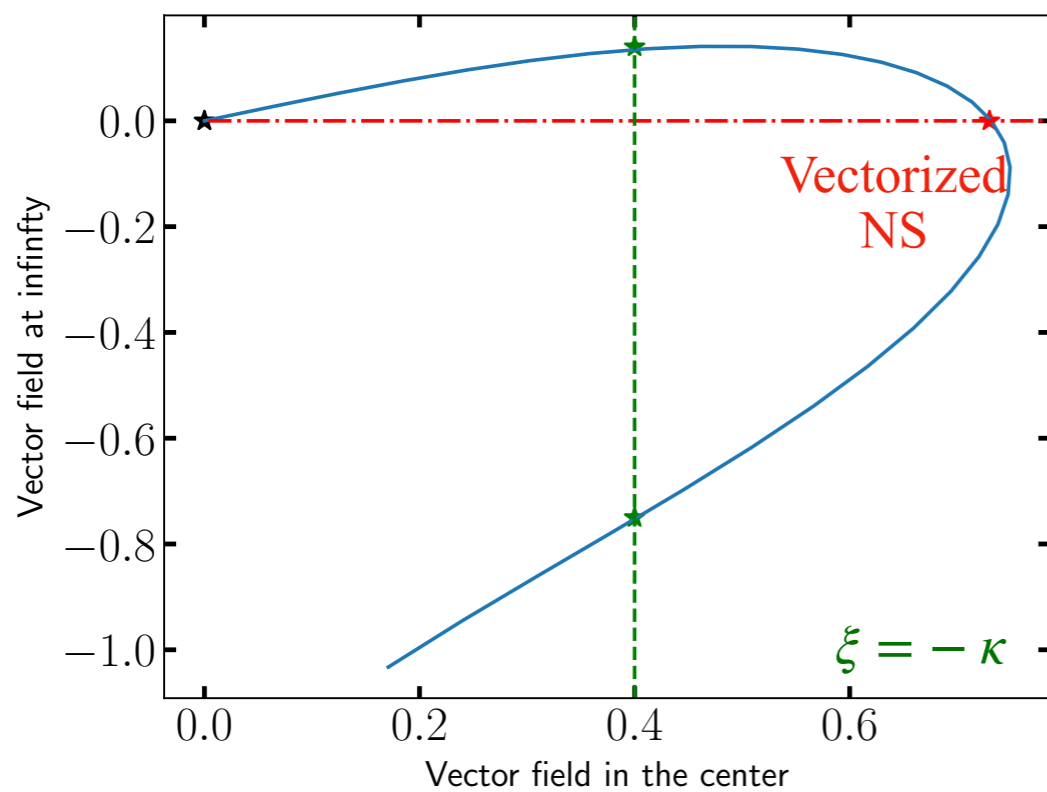
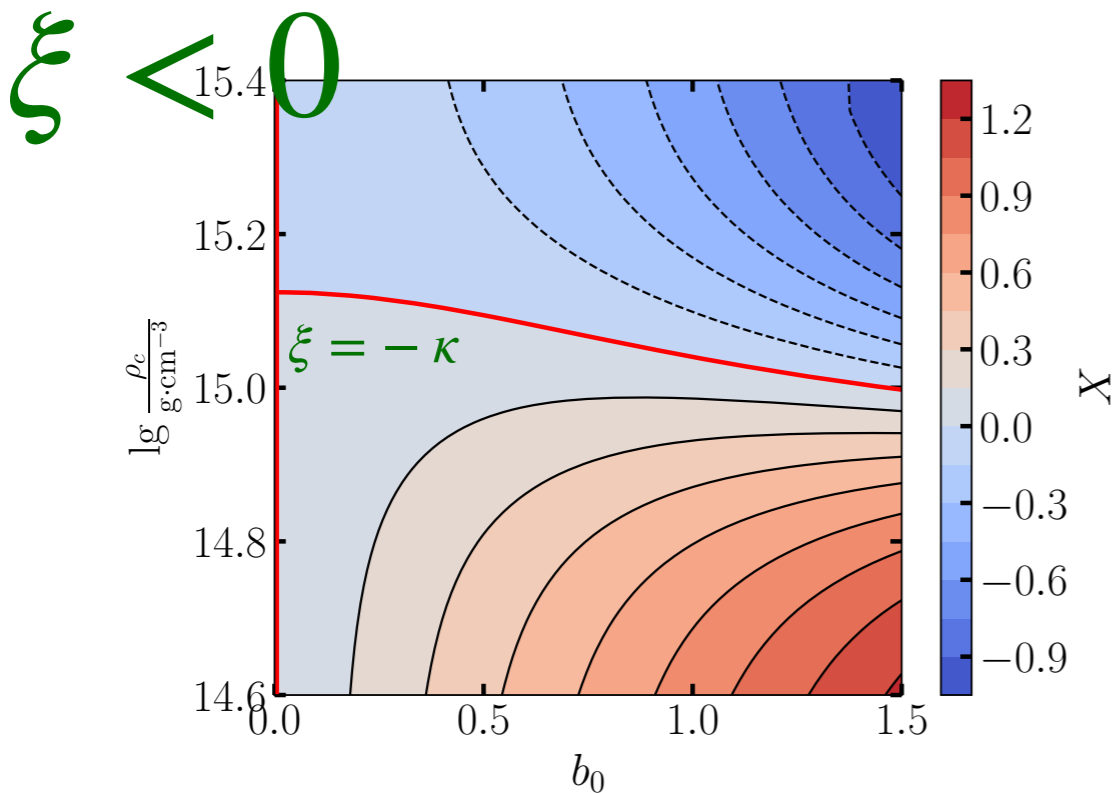


$GMX + \sqrt{\pi}Q = 0$
Only 2 of M, X, Q
are independent!

Brikhoff Theorem

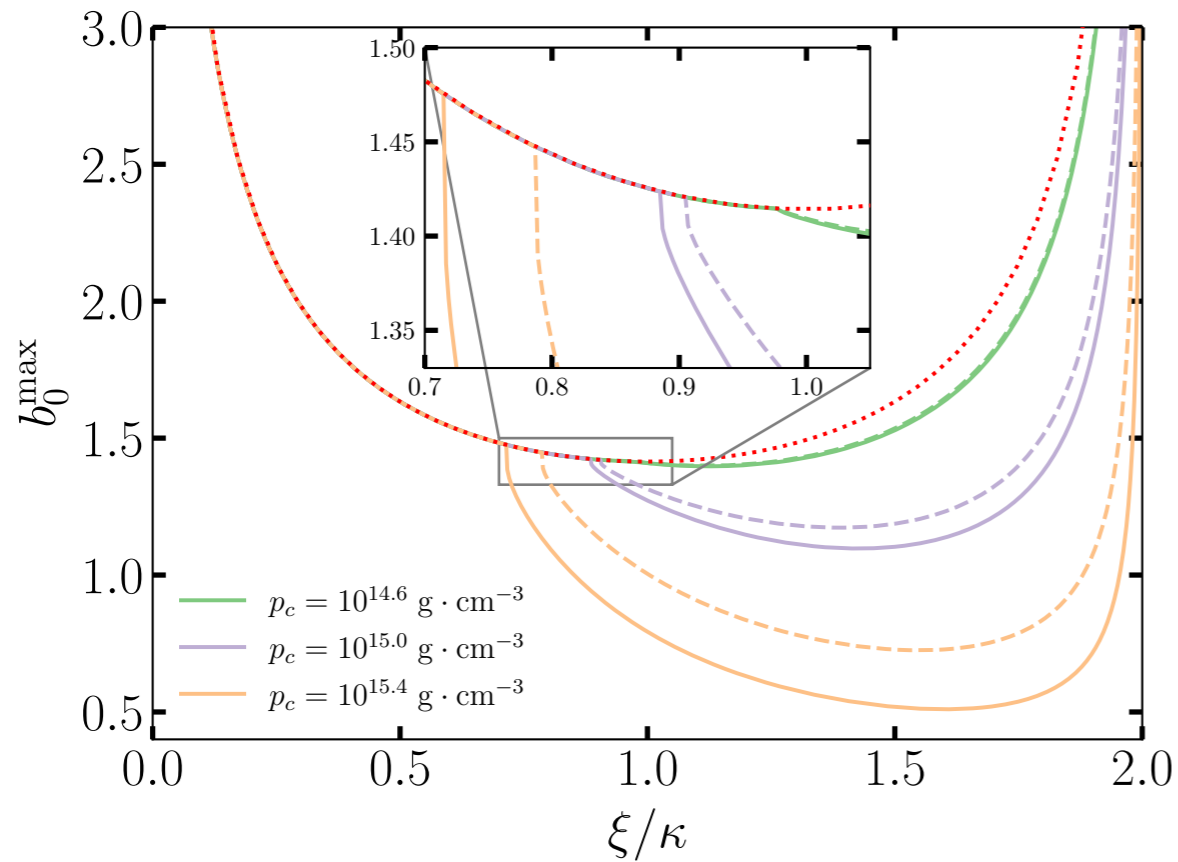


[R. Xu, D. Liang & L. Shao (2023)]



No spontaneous vectorization so far!
 See [L. Annulli *et al.* (2019)] & [H. Silva *et al.* (2022)].

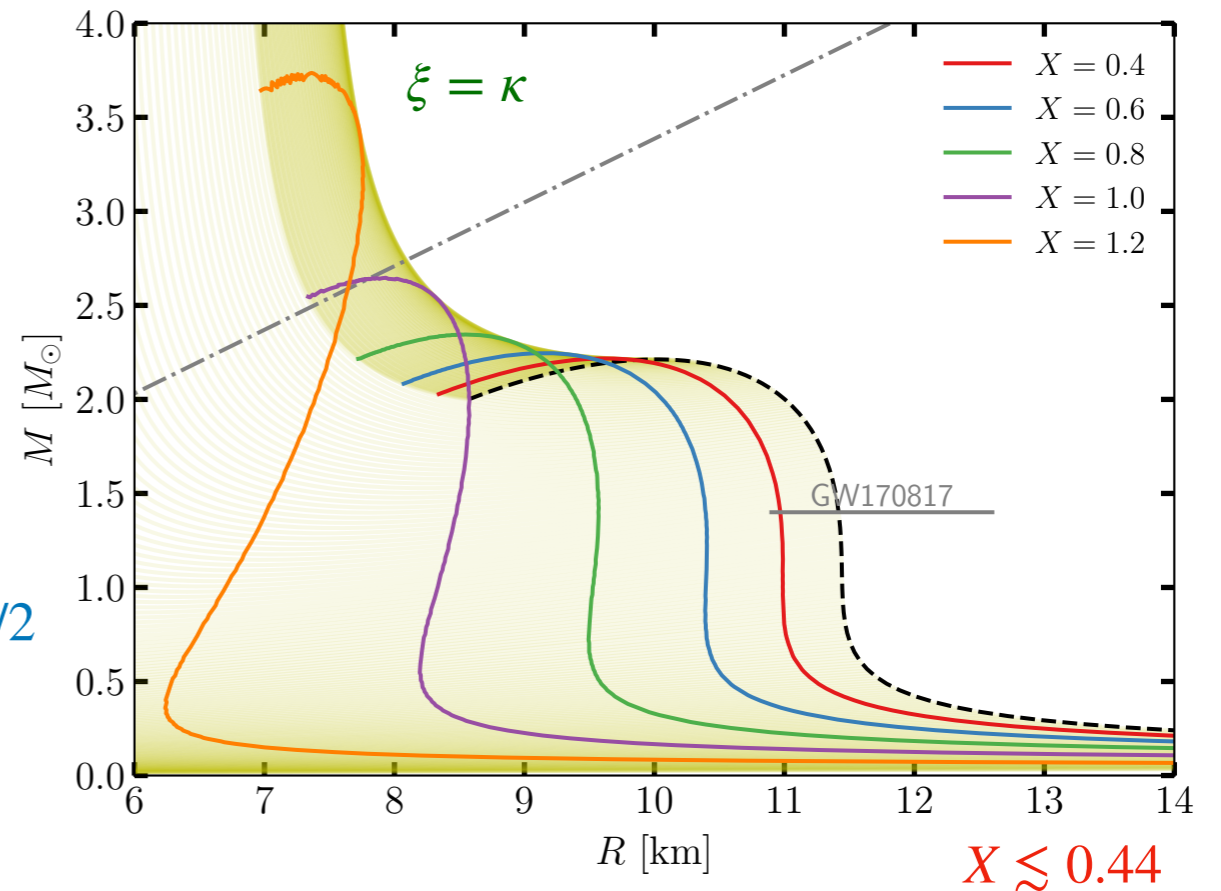
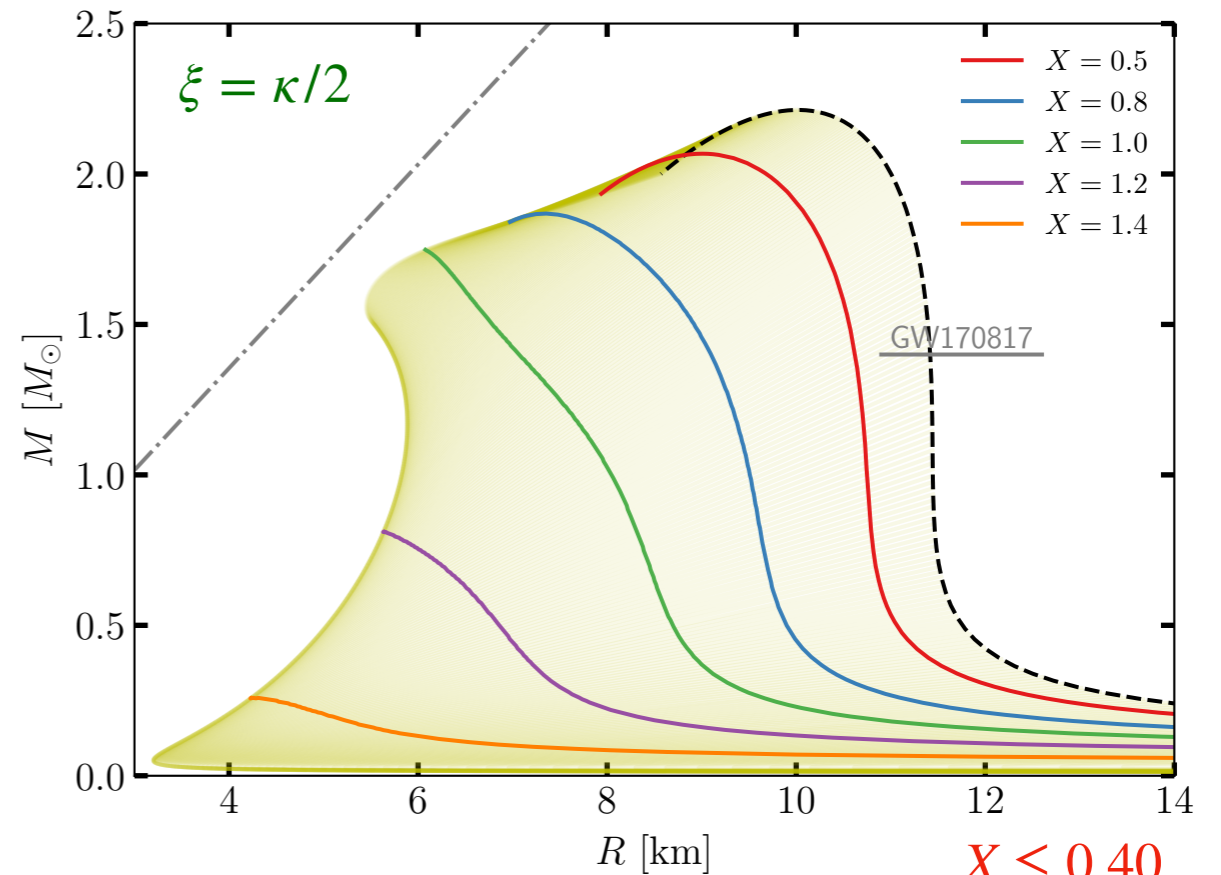
$$0 < \xi < 2\kappa$$



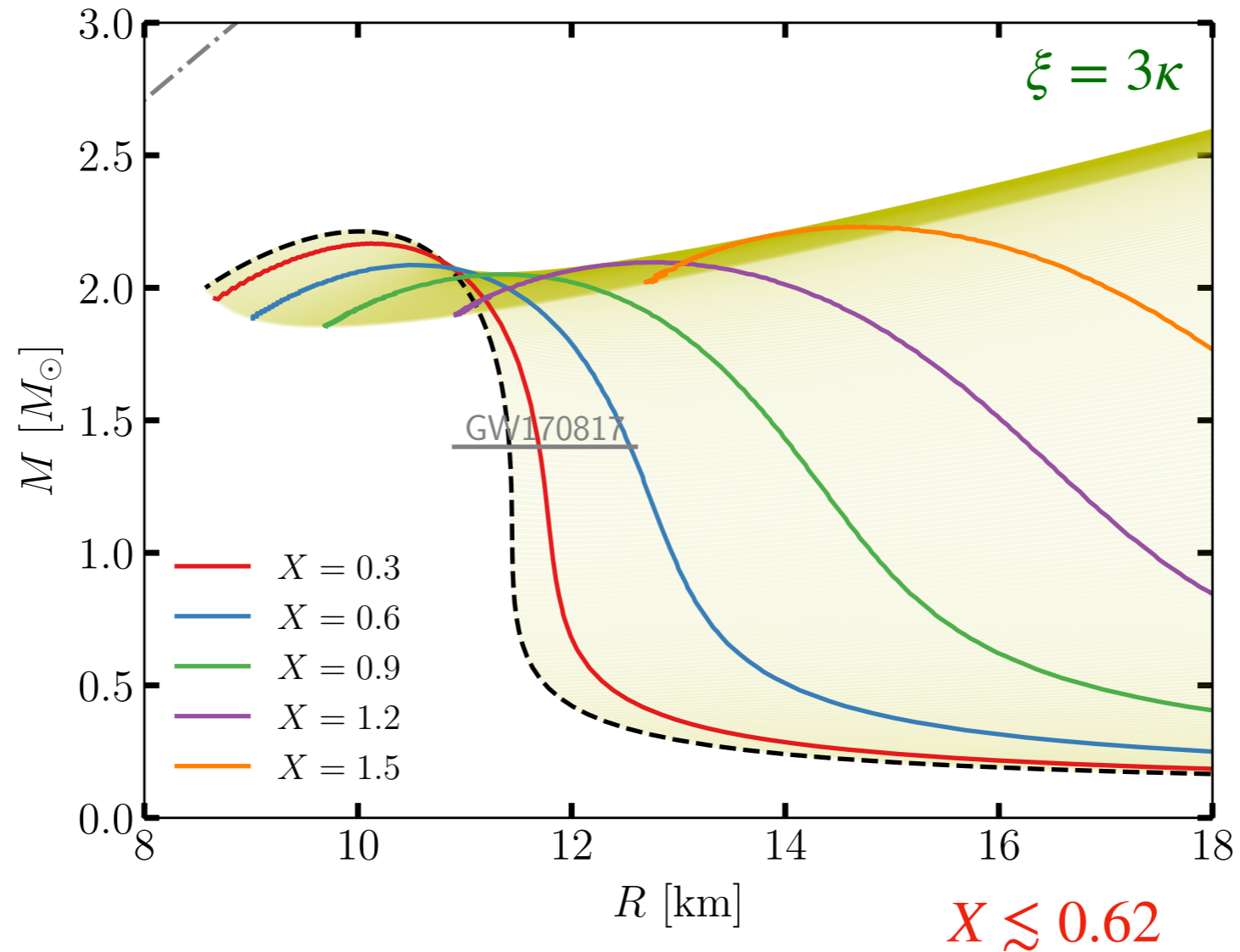
► Asymptotic behavior at center
 $\sim [2\kappa + \xi(\xi - 2\kappa)b_0^2]^{-1}$

► Asymptotic behavior at center
 $\sim [2\kappa + \xi(\xi - 2\kappa)X^2]^{-1}$

$$M, Q \rightarrow \infty \text{ as } X \rightarrow \left(\frac{2\kappa}{\xi(2\kappa - \xi)} \right)^{1/2}$$



$$\xi > 2\kappa$$



We are able to get both large radius and large mass in this range of non-minimal coupling without divergence.

Summary

- We systematically studied NSs in bumblebee theory
 - ▶ Differences between the well-studied scalar-tensor theory
- Some interesting cases
 - ▶ Large compactness ($\mathcal{C} > 0.5$) generally
 - ▶ Vectorized NSs when $\xi < 0$
 - ▶ Extremely massive (ADM mass) star with a relative small radius
 - ▶ Stealth Schwarzschild NS solutions
- Something to do
 - ▶ Perturbations on static spherical solutions
 - ✓ Tidal deformability — more constraints
 - ✓ Time-dependent perturbation — stability
 - ▶ Different possible potentials in the SME framework

Thank you!