

# EOS-dependent millihertz quasi-periodic oscillation in low-mass X-ray binary

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## ABSTRACT

We studied the frequency and critical mass accretion rate of millihertz quasi-periodic oscillations (mHz QPOs) using a one-zone X-ray burst model. The surface gravity is specified by two kinds of equation of states: neutron star (NS) and strange star (SS). The base flux,  $Q_b$ , is set in the range of 0–2 MeV nucleon<sup>-1</sup>. It is found that the frequency of mHz QPO is positively correlated to the surface gravity but negatively to the base heating. The helium mass fraction has a significant influence on the oscillation frequency and luminosity. The observed 7–9 mHz QPOs can be either explained by a heavy NS/light SS with a small base flux or a heavy SS with a large base flux. As base flux increases, the critical mass accretion rate for marginally stable burning is found to be lower. Meanwhile, the impact of metallicity on the properties of mHz QPOs was investigated using one-zone model. It shows that both the frequency and critical mass accretion rate decrease as metallicity increases. An accreted NS/SS with a higher base flux and metallicity, combined with a lower surface gravity and helium mass fraction, could be responsible for the observed critical mass accretion rate ( $\dot{m} \simeq 0.3\dot{m}_{\text{Edd}}$ ). The accreted fuel would be in stable burning if base flux is over than  $\sim 2$  MeV nucleon<sup>-1</sup>. This finding suggests that the accreting NSs/SSs in low-mass X-ray binaries showing no type I X-ray bursts possibly have a strong base heating.

**Key words:** accretion – stars:neutron star – X-rays: burst.

## 1 INTRODUCTION

It has a long history to understand dense matter with both sub-nuclear and supra-nuclear densities (Fowler 1926), and a state of neutron star (NS) matter with extremely low charge-mass-ratio was proposed by Landau (1932) before the discovery of neutron (Yakovlev et al. 2013). However, after the establishment of the standard model of particle physics, another way to make a neutral state of matter has also been focused (see e.g. Xu 2023, for a brief review), resulting in so-called strange matter with symmetry restoration of three-flavoured quarks (i.e.  $u$ ,  $d$ , and  $s$ ). It is then a general thought that strangeness would play an essential role in understanding the equation of state (EOS) of dense matter in pulsar-like compact objects. As for the basic units of supra-nuclear matter with strangeness, quarks could be itinerant there if the running coupling of the fundamental strong interaction is asymptotically free (Witten 1984; Alcock, Farhi & Olinto 1986; Haensel, Zdunik & Schaeffer 1986), but it is as well speculated that quarks could be clustered in nucleon-like units (so-called strangeons) if rich non-perturbative effects dominate at a few nuclear densities

(Xu 2003). For the sake of simplicity, we call them strange stars (SSs) in this paper, to be independent of which form the quarks take (i.e. either quarks or strangeons). Nevertheless, thermonuclear X-ray bursts on compact star’s surface are test-bed for us to identify the EOS, and this is our focus in this work.

Type I X-ray bursts are powered by unstable thermonuclear burning of accumulated hydrogen and helium on the surface of NS in a low-mass X-ray binary (LMXB; for reviews, see Lewin, van Paradijs & Taam 1993; Bildsten 1998; Strohmayer & Bildsten 2006; Galloway & Keek 2021). The mass accretion rate  $\dot{M}$  is thought to determine the different burning regimes (e.g. Fujimoto, Hanawa & Miyaji 1981; Bildsten 1998). Either theoretical arguments or observational evidence indicate that thermonuclear burning of the accreted material proceeds stably at very high  $\dot{M}$ , close to or above the Eddington mass accretion rate  $\dot{M}_{\text{Edd}}$  (Lewin et al. 1993; Bildsten 1998). Paczynski (1983) predicted an oscillatory burning regime near the transition between unstable and stable burning, which is also known as marginally stable burning. At present we know of 118 bursting sources.<sup>1</sup> A much smaller number of objects, ten

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<sup>1</sup><https://personal.sron.nl/~jeanz/bursterlist.html>

at present count, have shown millihertz quasi-periodic oscillations (mHz QPOs).

Revnivtsev et al. (2001) reported the first detection of mHz QPOs in three NS low-mass X-ray binaries: 4U 1608–52, 4U 1636–53, and Aql X-1. These mHz QPOs occur only within a narrow range of luminosity,  $L_X \approx 0.5 - 1.5 \times 10^{37}$  erg s $^{-1}$ , in the frequency range 7–9 mHz. They attributed this new phenomenon to a special mode of nuclear burning on the NS surface. Heger, Cumming & Woosley (2007) found that marginally stable nuclear burning on the NS surface could lead to the observed mHz QPOs.

There are discrepancies between the observations and theory, for instance, the disagreement in the accretion rate for the oscillation to exist. Theoretical models of hydrogen and helium burning suggest this critical mass accretion rate to be located around the Eddington limit, while observations place the marginally stable burning close to  $0.1 \dot{M}_{\text{Edd}}$ . Heger et al. (2007) proposed that the local accretion rate at the burning depth, where the mHz QPOs happen, can be higher than the entire NS surface. Piro & Bildsten (2007) and Keek, Langer & in't Zand (2009) showed a possible solution to solve this discrepancy by including mixing processes, e.g. rotation and rotationally induced magnetic fields. Furthermore, Keek, Cyburt & Heger (2014) showed that the range of accretion rates with marginally stable burning depends strongly on both the composition of the burning layer and reaction rates.

Later, these mHz QPOs were detected in seven other sources, IGR J17480–2446 (Linares et al. 2010), 4U 1323–619 (Strohmayer & Smith 2011), IGR J00291+5934 (Ferrigno et al. 2017), GS1826–238 (Strohmayer et al. 2018), EXO 0748–676 (Mancuso et al. 2019), 1RXS J180408.9–342058 (Tse et al. 2021), and 4U 1730–22 (Mancuso et al. 2023). The mHz QPOs have frequencies below  $\sim 14$  mHz. These mHz QPOs are evidently related to the thermal nuclear burning on the NS surface, as type I X-ray bursts are also present in the observations.

Some mHz QPOs appear with different characteristics. The persistent luminosity of IGR J17480–2446 when the mHz QPOs were observed was high,  $L_{2-50 \text{ keV}} \sim 10^{38}$  erg s $^{-1}$ , the frequency was always below 4.5 mHz (Linares et al. 2010). Altamirano et al. (2008) observed frequency drifts in mHz QPOs from 4U 1636–53. Lyu et al. (2015) measured the frequency drift time of the mHz QPOs in 4U 1636–536 from *XMM-Newton* and Rossi X-ray Timing Explorer and speculated that the drift time-scale was set by cooling of the deeper layers. Li, Pan & Falanga (2021) discovered a 2.7–11.3 mHz QPOs after a superburst in Aql X-1.

Understanding the observations of mHz QPOs allows us to probe the envelope composition, nuclear reaction rate, and the NS EOSs (Heger et al. 2007; Keek et al. 2014; Dohi et al. 2021; Zhen et al. 2023). No matter the one-zone model or the multi-zone model such as KEPLER for theoretical calculations of thermonuclear burning, a base flux  $Q_b$  is included as an artificial parameter, which models the heat flowing from the underlying NS crust into the envelope. Most simulations use a low value of  $Q_b = 0.1 - 0.15 \text{ MeV u}^{-1}$  (Heger et al. 2007; Meisel 2018). However, more recent work indicates that the base flux may be larger, up to  $Q_b = 2 \text{ MeV u}^{-1}$  (Haensel & Zdunik 2003; Gupta et al. 2007). Additionally, a yet unknown shallow heat source may increase the base flux (Cumming et al. 2006; Deibel et al. 2015; Fattoyev et al. 2018; Liu et al. 2021). Furthermore, a SS can also have a crust of normal baryonic matter (Alcock et al. 1986). Page & Cumming (2005) studied the superburst from SS by considering the energy release from the dripped neutrons fall into quark matter. How about the mHz QPOs from the marginally stable burning on accreting SSs? The gravity and base flux may cause the differences compared with the accreting NSs.

In this work, we investigate the dependence of the oscillation frequency and critical mass accretion rate on the gravity and base flux using one-zone X-ray burst model. The gravity will be obtained from the EOSs of typical NS/SS. In Section 2, we describe the basic equations of the one-zone model, the adopted EOSs, and inputting parameters. In Section 3, we show the results of the marginally stable burning from NS and SS. The effect of gravity, base flux, and helium mass fraction on the frequency and critical mass accretion rate are discussed. Besides, the impact of metallicity on the properties of mHz QPOs with use of one-zone model are investigated. In section 4, the general relativity corrections are discussed. Section 5 is devoted to conclusion.

## 2 ONE-ZONE MODEL

To study the marginally stable burning, we use a publicly available code of one-zone model by A. Cumming<sup>2</sup> (see also Paczynski 1983; Heger et al. 2007). We describe the basic equations and microphysics considered in the following calculations.

### 2.1 The basic equations

For the one-zone model calculations, the basic equations are as follows (Bildsten 1998; Heger et al. 2007):

$$P = g_s y, \quad (1)$$

$$c_p \frac{dT}{dt} = \varepsilon - \frac{F}{y}, \quad (2)$$

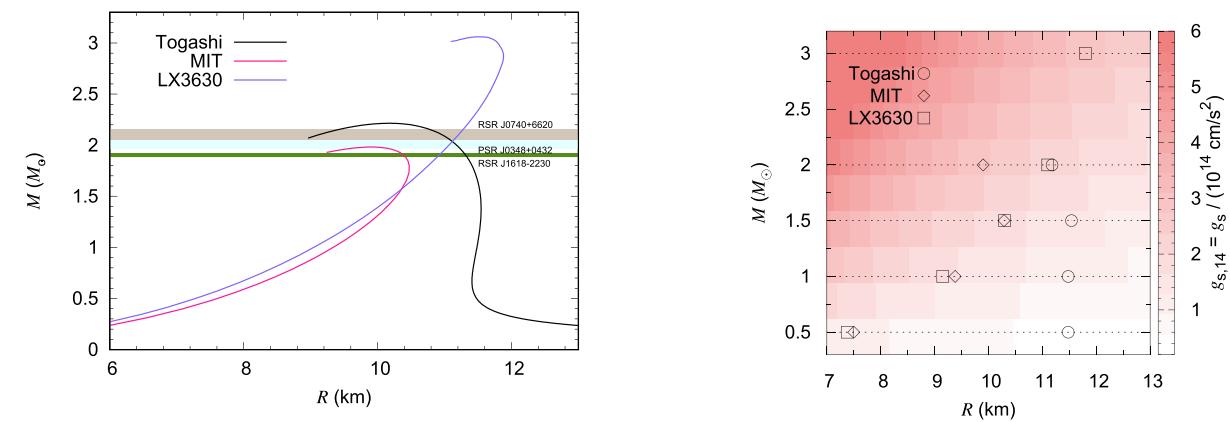
$$\frac{dy}{dt} = \dot{m} - \frac{\varepsilon}{E_*} y. \quad (3)$$

where  $g_s$  is the surface gravity,  $c_p$  is the specific heat capacity at the constant pressure,  $y$  is the column density, in units of mass per unit area, and  $\dot{m}$  is the local accretion rate.  $\varepsilon$  is the heating from the burning layer, the code considers three kinds of energy generation: the energy release from the triple- $\alpha$  reaction, hot-CNO hydrogen burning, and base heating from the bottom of the accumulating fuel layer.  $E_*$  is the energy per gram released in the burning layer,  $E_* = 5.3 \text{ MeV u}^{-1}$  is adopted in this work, where the changes of  $E_*$  with hydrogen mass fraction are not considered here (Schatz et al. 1999; Heger et al. 2007). Equation (1) is derived from hydrostatic balance, which describes the pressure  $P$  at the base of the accretion layer. Equation (2) describes the heat balance, including the energy generation rate  $\varepsilon$ , and radiative loss rate,  $-\nabla \cdot \mathbf{F}/\rho = dF/dy \approx F/y$ . The flux  $F = acT^4/3\kappa y$  (Bildsten 1998), where the opacity  $\kappa$  is calculated as described by Schatz et al. (1999). Equation (3) tracks the burning depth. The pressure  $P$  includes the non-relativistic degenerate electron pressure, ion pressure, and radiation pressure, the details of the EOS in the accreting layer can be referred to equation (24) from Paczynski (1983). In the following, we will study the marginally stable burning by using the public code onezone.

### 2.2 The EOS and heating

To date, there is no real evidence against or in favour of the SS hypothesis (Prakash et al. 2003), since the properties such as radius and moment of inertia are similar for NS/SS in the observed mass range  $1 M_\odot < M < 2 M_\odot$ . The X-ray burst phenomenon is modelled on the surface of accreting NSs in most of the works. However,

<sup>2</sup><https://github.com/andrewcumming/onezone>



**Figure 1.** Left:  $M$ – $R$  relations of NS (black), SS (pink), and strangeon star (purple) adopted in this work, and the  $1\sigma$  region of the mass measurements of PSR J1618–2230 (Demorest et al. 2010), PSR J0348+0432 (Antoniadis et al. 2013), and PSR J0740+6620 (Fonseca et al. 2021) are highlighted. Right: The surface gravity acceleration  $g_s$  on the  $M$ – $R$  plane. The symbols indicate the values of mass and radius of NS/SSs in our model calculations.

there are many works support the existence of the accreting SS (Haensel & Zdunik 1991; Page & Cumming 2005; Zhu et al. 2013; Zhang, Geng & Huang 2018). If we consider the X-ray burst from the surface of accreting SS, there will be  $\sim 10$ –100 MeV of energy that comes from dripped neutrons fall into the quark matter (Farhi & Jaffe 1984; Page & Cumming 2005). Much of this energy is carried off by neutrino, but a considerable part of this heat will be left and alter the temperature of the burning layer. Page & Cumming (2005) considered this energy and calculated the ignition condition for the superburst from SSs. The consequences of this alteration for thermal nuclear burst model have not been explored. On the other hand, the cold quark matter could be strange quark clusters because of the strong interaction between quarks (Xu 2003). A strange quark cluster is named ‘strangeon’, being coined to ‘strange nucleon’ for simplicity (Xiaoyu & Renxin 2017). Therefore, the pulsar like compact stars could be strangeon stars. How about X-ray burst on the accreting strangeon star?

In this work, we simplified the X-ray burst simulations from accreting NS and SS with use of one-zone model. The different properties of them are reified by the surface gravity which is obtained from the mass–radius ( $M$ – $R$ ) relationship and base flux. Thus, EOS and base flux are important factors to determine  $g_s$  and energy generation rate which distinguish the X-ray burst scenario. For NS, we adopt the Togashi EOS which was constructed by using realistic two-body potential and phenomenological three-body potential (Togashi et al. 2017). For SS star, we consider strange quark star and strangeon star, for the former one, the MIT bag model of quark matter with bag constant  $B = 70$  MeV fm $^{-3}$  is used (Alcock et al. 1986). For the latter one, we adopted LX3630 EOS which is based on the Lennard–Jones model with the surface baryon number density  $n_s = 0.36$  fm $^{-3}$  and the potential depth  $\epsilon = 30$  MeV (Lai & Xu 2009; Gao et al. 2022).

The  $M$ – $R$  relations of the three adopted EOSs are shown in the left panel of Fig. 1. In the right panel of Fig. 1, we show the values of the local gravity  $g_s$  obtained using the Newtonian formula  $g_s = GM/R^2$  for several  $M$  and  $R$ . However, note that general relativity corrections should be applied when relating a value of  $g_s$  to  $M$  and  $R$  from a particular EOS and when using the results for mHz QPOs (see discussion in section 4.4 of Woosley et al. 2004 and Section 4). None the less, Heger et al. (2007) noted that onezone code overpredicts the oscillation frequency, thus somewhat compensating the corrections (see footnote 7 of Heger et al. 2007). In table 1,

we show the masses and corresponding (Newtonian) surface local gravity for each EOSs, where  $g_{14}$  is the gravitational acceleration in units of  $10^{14}$  cm s $^{-2}$ .

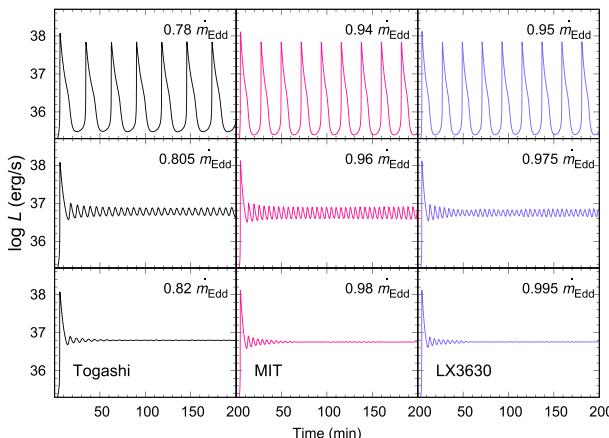
For the nuclear burning, the code considers hydrogen and helium burning (Heger et al. 2007). The nuclear energy generation rate from the hot CNO hydrogen burning is as follows (Wallace & Woosley 1981):

$$\varepsilon_{\text{hCNO}} = 5.86 \times 10^{15} Z \text{ erg g}^{-1} \text{ s}^{-1} \quad (4)$$

where  $Z$  represents the sum of the mass fraction of the metallicity inside the accreted layer. For the helium burning, we use the  $3\alpha$  reaction rate as given by Fushiki & Lamb (1987). The code includes a factor of  $E_{\text{nuc}}/E_{3\alpha}$  to enhance the  $3\alpha$  reaction rate (Heger et al. 2007), where  $E_{\text{nuc}} = 1.6 + 4.9 X_0$  MeV u $^{-1}$ ,  $X_0$  is the mass fraction of hydrogen in the accreted layer.  $E_{3\alpha} = 0.606$  MeV u $^{-1}$  is the energy release from the  $3\alpha$  reaction. Besides, the code allows for the inclusion of a base heating at the bottom of the accreted layer. The values of  $Q_b$  is set in the range of  $0$ – $2$  MeV u $^{-1}$ , the corresponding base luminosity  $L_b = Q_b \dot{M}$ . When  $Q_b > 2$  MeV u $^{-1}$ , the accreted fuel will be in stable burning which we will discuss in the next section. The total energy generation rate  $\varepsilon$  contains the nuclear burning including hydrogen and helium and the base heating. We use the public onezone code to integrate equations (2) and (3) with respect to time to determine the evolution of the burning layer. The simulations are started with an arbitrary conditions at  $T = 2 \times 10^8$  K and  $y = 2 \times 10^8$  g cm $^{-2}$  following Heger et al. (2007).

### 3 MARGINALLY STABLE BURNING FROM ACCRETING NS AND SS

The mHz QPO phenomenon is thought to occur at the boundary of unstable and stable nuclear burning. This oscillatory model of burning also called ‘marginally stable burning’. Heger et al. (2007) studied the marginally stable nuclear burning on an accreting NS with both one-zone model and multi-zone hydrodynamic KEPLER code. They derived the period of the oscillations close to 2 min, which is in good agreement with the range of oscillation frequencies reported for the mHz QPOs. Here, we consider the marginally stable burning from accreting NS and SS with one-zone model as described in section 2. For convenient, we define the local Eddington accretion rate  $\dot{m}_{\text{Edd}} \equiv 8.8 \times 10^4$  g cm $^{-2}$  s $^{-1}$  throughout this paper.



**Figure 2.** Light curves for the one-zone model at three different accretion rates, where  $M = 1.4 M_{\odot}$ ,  $Q_b = 0.15 \text{ MeV u}^{-1}$ , and  $X_0 = 0.7$ . For Togashi EoS, at  $\dot{m} = 0.78 \dot{m}_{\text{Edd}}$ , the recurrence time of bursts is 28 min. At  $\dot{m} = 0.805 \dot{m}_{\text{Edd}}$ , oscillations are seen with a period of 6.8 min. At  $\dot{m} = 0.82 \dot{m}_{\text{Edd}}$ , after a few transient oscillations, the burning evolves to a steady state. For MIT EoS, at  $\dot{m} = 0.94 \dot{m}_{\text{Edd}}$ , the recurrence time of bursts is 22 min. At  $\dot{m} = 0.96 \dot{m}_{\text{Edd}}$ , oscillations are seen with a period of 5.2 min. The burning evolves to a steady state at  $\dot{m} = 0.98 \dot{m}_{\text{Edd}}$ . For LX3630 EoS, at  $\dot{m} = 0.95 \dot{m}_{\text{Edd}}$ , the recurrence time of bursts is 22 min. At  $\dot{m} = 0.975 \dot{m}_{\text{Edd}}$ , oscillations are seen with a period of 5.2 min. The burning evolves to a steady state at  $\dot{m} = 0.995 \dot{m}_{\text{Edd}}$ .

### 3.1 Light curves for the one-zone model from accreting NS and SS

Fig. 2 shows light curves from the one-zone model at different accretion rates. We set the local gravity to be the Newtonian gravity for a  $1.4 M_{\odot}$  NS/SSs star from Togashi, MIT and LX3630 EoSs, the corresponding  $g_{14}$  equals 1.39, 1.80, and 1.82, where  $g_{14} = g_s / (10^{14} \text{ cm}^2 \text{s}^{-1})$ . We show light curves at  $\dot{m} = 0.78 \dot{m}_{\text{Edd}}$ ,  $0.805 \dot{m}_{\text{Edd}}$ ,  $0.82 \dot{m}_{\text{Edd}}$  for Togashi EoS,  $\dot{m} = 0.94 \dot{m}_{\text{Edd}}$ ,  $0.96 \dot{m}_{\text{Edd}}$ ,  $0.98 \dot{m}_{\text{Edd}}$  for MIT EoS, and  $\dot{m} = 0.95 \dot{m}_{\text{Edd}}$ ,  $0.975 \dot{m}_{\text{Edd}}$ ,  $0.995 \dot{m}_{\text{Edd}}$  for LX3630 EoS. Close to the stability boundary, the recurrence time is 28 min for Togashi EoS, 22 min for MIT, and LX3630 EoSs. At accretion rates above the boundary, the burning evolves to a steady state. At  $\dot{m} = 0.805 \dot{m}_{\text{Edd}}$ ,  $0.96 \dot{m}_{\text{Edd}}$ ,  $0.975 \dot{m}_{\text{Edd}}$ , we see oscillations with an oscillation period of 6.8, 5.2, and 5.2 min for Togashi EoS, MIT EoS, and LX3630 EoS, respectively. The more compact the star, the greater the transition mass accretion rate.

Fig. 3 shows the peak luminosity and recurrence time change with masses and base flux  $Q_b$ . The bursts are simulated at a fixed mass accretion rate  $\dot{m} = 0.3 \dot{m}_{\text{Edd}}$  and hydrogen fraction  $X_0 = 0.7$ . It shows a positive correlation for every EOS for the peak luminosity at a fixed  $Q_b$ , but an anticorrelation for every EOS for the recurrence time at a fixed  $Q_b$ . At a fixed mass for every EOS, both the peak luminosity and recurrence time are anti-correlation with  $Q_b$ . The accreted material will be in steady burning if the base flux  $Q_b$  is high, e.g. for  $0.5 M_{\odot}$  NS, the critical  $Q_b \simeq 1 \text{ MeV/u}$ , for  $2.0 M_{\odot}$  NS, the critical  $Q_b \simeq 2.1 \text{ MeV/u}$ , and for  $3.0 M_{\odot}$  strangeon star<sup>3</sup>, the critical  $Q_b \simeq 2.5 \text{ MeV/u}$ .

<sup>3</sup>The  $3 M_{\odot}$  SS is adopted to show the effect of strong gravity. However, the mass measurement of the three massive pulsars in the left panel of Fig. 1 show that the maximum mass of pulsar is about  $\sim 2 M_{\odot}$ . In order to compare with the massive mass measurement of pulsar, we thought that the accreted fuel would be in stable burning if  $Q_b \gtrsim 2 \text{ MeV u}^{-1}$  for simplicity.

### 3.2 Oscillation frequency for the one-zone model from accreting NS and SS

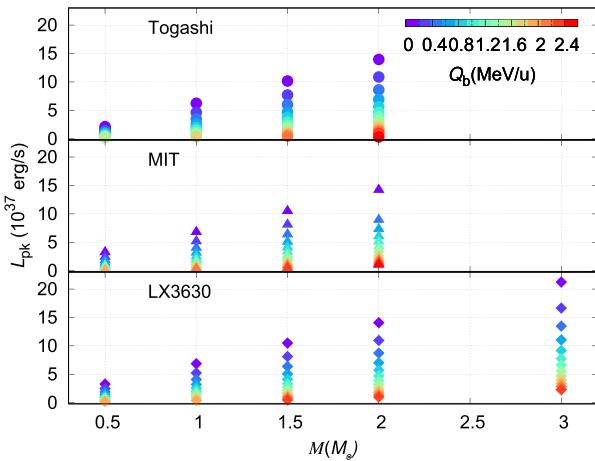
In the following, we quantify the properties of the oscillations in the light curves of all simulations with marginally stable burning by determining the oscillation period/frequency, and the maxima and minima luminosity.

Fig. 4 shows the dependence of the (uncorrected) oscillation frequency on the accretion rate for different choices of NS/SS masses. At each accretion rate, we integrate the one-zone model for  $10^6$  s, and plot the average oscillation frequency after discarding the first 100 min of data following Heger et al. (2007). We can see that the overall range of accretion rates for which oscillations are seen is very narrow,  $\Delta\dot{m} \approx 0.01\dot{m}_{\text{Edd}}$ . For  $0.5 M_{\odot}$  NS, it shows a very low-frequency mHz QPOs,  $f < 1 \text{ mHz}$ , although we have not observed such a low-frequency mHz QPOs so far. For a  $0.5 M_{\odot}$  strange (strangeon) star, the oscillation frequency is  $\sim 2 \text{ mHz}$  QPOs. For a higher-mass NS/SS, as the surface gravity increases, the oscillation frequency increases. For  $1.5 M_{\odot}$  NS and SS, the oscillation frequencies are  $\sim 3$  and  $\sim 4 \text{ mHz}$ , respectively. For  $2.0 M_{\odot}$  NS and SS, the oscillation frequencies are  $\sim 4$  and  $\sim 5 \text{ mHz}$ , respectively. For the same mass NS and SS, the oscillation frequency is higher for SS star. Our results are consistent with Fig. 4 of Heger et al. (2007).

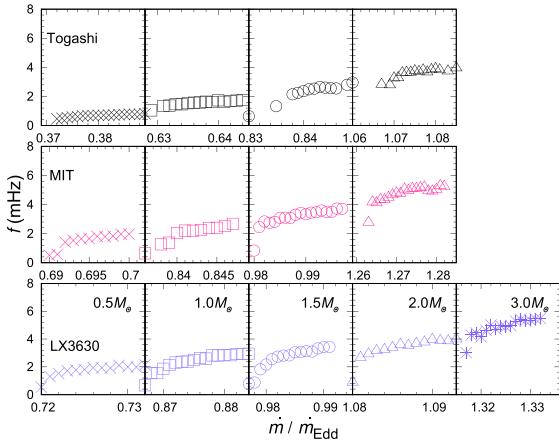
One-zone model has been used to explore how the properties of the oscillations change with parameters such as surface gravity and accretion composition (Heger et al. 2007). As the surface gravity is degeneracy between the combinations of mass and radius, the previous works do not distinguish the gravity from NS and SS. Here we adopt one-zone model, but the surface gravity is inferred from three representative EoSs. The effect of base heating  $Q_b$  on oscillation properties also will be investigated. Compared with multi-zone model coupled with the nuclear reaction network, one-zone model saves a lot of computation time. We have performed simulations for variations of helium mass fraction and base flux for a series of NS/SS masses, where the metallicity is set as  $Z = 0.01$ .

Fig. 5 shows the dependence of the oscillation frequency on the helium mass fraction of the accreted material and base flux for several NS/SS masses. When increasing  $Y_0$ , the oscillation frequency increases. Values of  $f$  range from 0.35 to  $\sim 39 \text{ mHz}$  for NS, and 0.84 to  $\sim 54 \text{ mHz}$  for SS. The oscillation frequency decreases as  $Q_b$  increases, and increases as the mass of NS/SS increases. For the same NS/SS mass, the oscillation frequency of SS is higher than NS if the base flux and hydrogen mass fraction are also the same. For NS, the 7–9 mHz QPOs can be fitted with a  $1.5 M_{\odot}$  with helium mass fraction  $Y_0$  in the range of 0.6–0.7, while for SS, the 7–9 mHz QPOs can be satisfied with a  $1.5 M_{\odot}$  with a helium mass fraction  $Y_0$  in the range of 0.5–0.6. The details of the output quantities can be found in Tables A1–A3 in Appendix.

For above series of simulations, we study the properties of oscillation frequency as well as critical mass accretion rate with different NS/SS masses, accretion compositions, and base flux. As we can get the surface gravity from the NS/SS's  $M$ - $R$  relationship (the values of the local surface gravity we adopted can be seen from Table 1), we show the distributions of the oscillation frequencies and surface gravity for different base heating and helium mass fractions scenarios in Fig. 6. The colour represents the last value of the critical mass accretion rate in its narrow range at which the burning is marginally stable. For these similar models for  $Q_b$ ,  $Y_0$ , and  $g_{14}$  variations, we find that surface gravity has a significant impact on oscillation frequency and critical mass accretion rate. At higher surface gravity, both the oscillations frequency and the critical mass accretion rate are larger. Moreover, the amplitudes of  $f$  and  $\dot{m}$



**Figure 3.** The peak luminosity and recurrence time of the bursts for a series NS/strange (strangeon) star. The different colours mark the different values of the base flux. The mass accretion rate is set as  $\dot{m} = 0.3 \dot{m}_{\text{Edd}}$  and  $X_0 = 0.7$ .



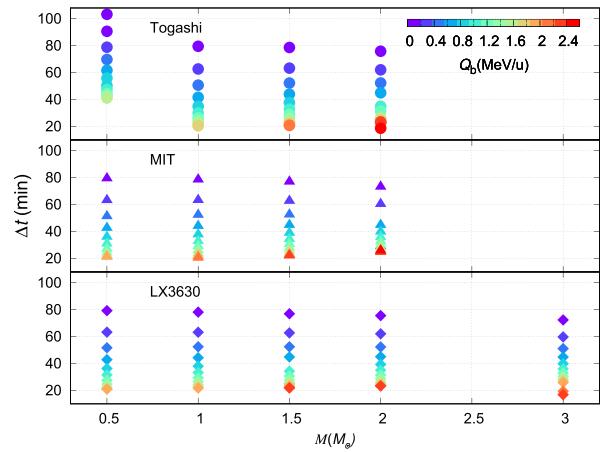
**Figure 4.** Oscillation frequency as a function of mass accretion rate for different choices of NS/strange(strangeon) star masses.  $X_0 = 0.7$  and  $Q_b = 0.15 \text{ MeV u}^{-1}$ .

**Table 1.** The values of local gravity for several values of  $M$  with different EoSs.

Mass ( $M_\odot$ )	$g_{14}(\text{Togashi})$	$g_{14}(\text{MIT})$	$g_{14}(\text{LX3630})$
0.5	0.50	1.15	1.22
1.0	1.01	1.51	1.59
1.4	1.39	1.80	1.82
1.5	1.50	1.88	1.88
2.0	2.12	2.69	2.16
3.0	—	—	2.86

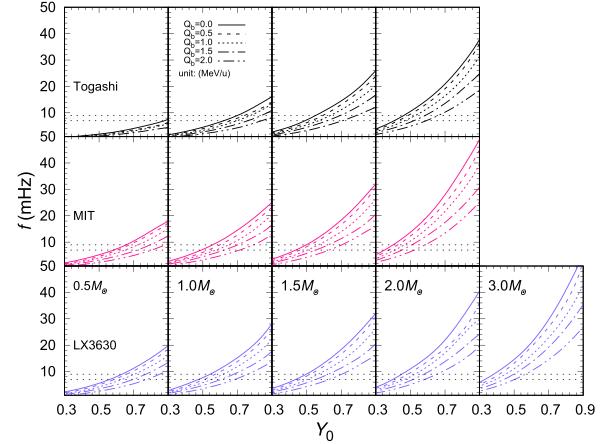
are higher at larger helium mass fraction and smaller base heating. Although our simulations are crude, the results suggest that one can infer the surface gravity if we obtain the frequency and mass accretion rate of the mHz QPOs.

We specify the amplitude of the oscillations in the light curves of all simulations with marginally stable burning by determining the maxima and minima luminosity. Fig. 7 shows the distributions of the oscillation luminosity and surface gravity for different base heating and helium mass fraction models. The colours of the symbols represent the critical mass accretion rate the same as Fig. 6. The difference between the maxima luminosity and minima luminosity



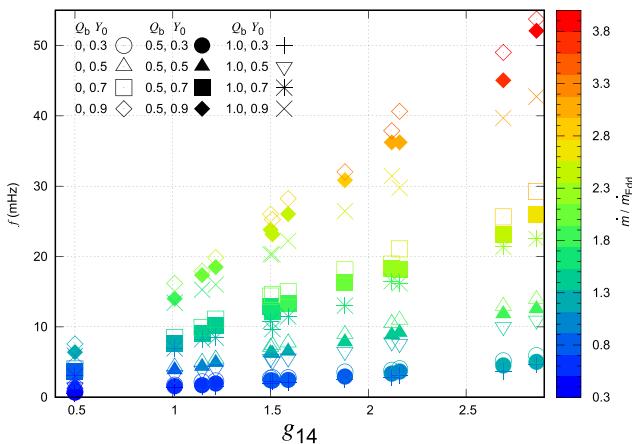
**Table 2.** The values for correcting for general relativity.

Mass ( $M_\odot$ )	$M_{\text{GR}}$	R (km)	$g_{14}$	$1 + z$
0.5	0.47	11.48	0.50	1.07
1.0	0.88	11.48	1.01	1.14
1.5	1.24	11.54	1.50	1.21
2.0	1.54	11.18	2.12	1.30
0.5	0.45	7.61	1.15	1.10
1.0	0.85	9.38	1.51	1.17
1.5	1.21	10.29	1.88	1.24
2.0	1.49	9.94	2.69	1.34
0.5	0.45	7.38	1.22	1.11
1.0	0.85	9.15	1.59	1.17
1.5	1.21	10.30	1.88	1.24
2.0	1.54	11.10	2.16	1.30
3.0	2.08	11.80	2.86	1.44



**Figure 5.** Oscillation frequency as a function of the helium mass fraction,  $Y_0$ , for different choices of NS/SS masses and base flux, where the metallicity is set as  $Z = 0.01$ . The horizontal dotted lines in each panel show the boundaries of the first observed 7–9 mHz QPOs.

increases as surface gravity increases except some dips at several gravities ( $g_{14} = 1.22, 1.39, 2.69$ ). Nevertheless, the trends in the properties are clearly: for larger  $Q_b$ , the difference between the maxima luminosity and minima luminosity are smaller, while



**Figure 6.** Oscillation frequency versus surface gravity for different base heating  $Q_b$  and helium mass fraction  $Y_0$  (marked as kinds of symbols). The colours indicate the critical mass accretion rate at which marginally stable burning occurs for each scenario.

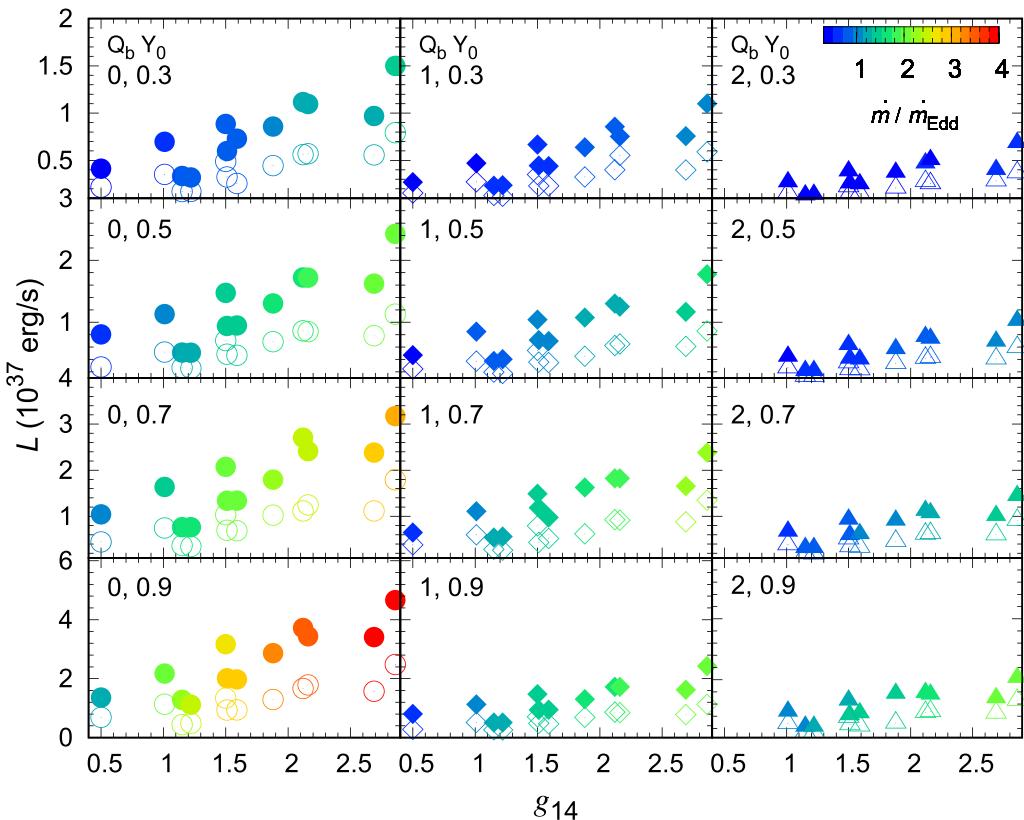
for larger  $Y_0$ , the difference between the maxima luminosity and minima luminosity are larger. Compared with the narrow luminosity observations,  $\dot{m}/\dot{m}_{\text{Edd}} \approx 0.1 - 0.3$ , we can exclude the models with low base flux and high helium mass fraction, e.g. the first panel in the fourth line of Fig. 7, for which the critical mass accretion rate is too large to fit the observations. Our models predict the critical mass accretion rate in the range  $\sim 0.27 - 4.3 \dot{m}_{\text{Edd}}$ . Although observations place the critical mass accretion rate around 0.1 to 0.3  $\dot{m}_{\text{Edd}}$  which is an order-of-magnitude lower than mass accretion rate implied by

the model calculation, it has been suggested that rotational mixing of the accreted fuel to greater depths could alleviate this apparent discrepancy (Keek et al. 2009).

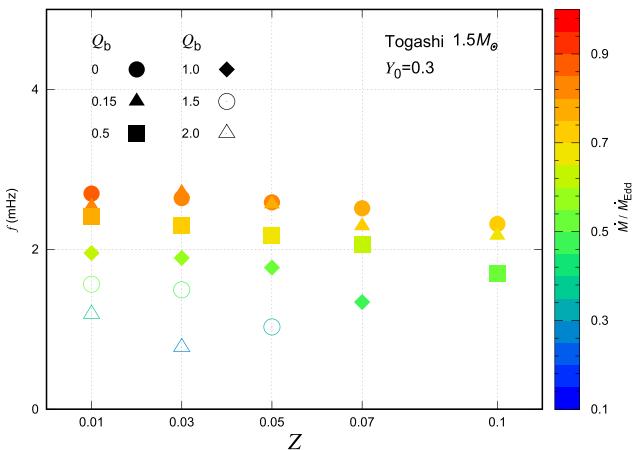
### 3.3 The impact of metallicity on the properties of mHz QPOs

From the above calculations, the oscillation properties have been studied by using one-zone model with variation in surface gravity, helium mass fraction, and base heating. Keek et al. (2014) find that  $Z$  and the nuclear reaction rates of the hot-CNO breakout reactions will change the properties of oscillation burning. However, Keek et al. (2014) adopt the multi-zone KEPLER code with a more complicated network, besides,  $Z$  is a linear function of hydrogen mass fraction  $X_0$ . So, how about the effect of  $Z$  for one-zone model? To answer this question, we study the effect of metallicity on the properties of mHz QPOs with use of one-zone model, where a typical NS ( $1.5 M_{\odot}$ ) with surface gravity  $g_{14} = 1.50$  and helium mass fraction  $Y_0 = 0.3$  is adopted,  $Z$  is set in the range of 0.01–0.1.

Fig. 8 exhibits the changes of oscillation frequency and critical mass accretion rate with the variation in metallicity. The symbols represent the different values of base heating, and the colours represent the values of critical mass accretion rate. Both the oscillation frequency and the critical mass accretion rate decrease as  $Z$  increases. For  $Q_b = 0$ , the oscillation frequency decreases by  $\approx 14$  per cent and the critical mass accretion rate decreases by  $\approx 20$  per cent, for  $Q_b = 0.5$ , the oscillation frequency decreases by  $\approx 29$  per cent and the critical mass accretion rate decreases by  $\approx 32$  per cent. For  $Q_b = 1.0 \text{ MeV u}^{-1}$ , it will be in stable burning when  $Z > 0.07$ , and for  $Q_b = 1.5 \text{ MeV u}^{-1}$  and  $2.0 \text{ Mev u}^{-1}$ , it will be in stable



**Figure 7.** Oscillation luminosity versus surface gravity for different base heating  $Q_b$  and helium mass fraction  $Y_0$ . The solid symbols in each panel indicate the maxima luminosity; the hollow symbols indicate the minima luminosity.



**Figure 8.** Oscillation frequency versus metallicity for different base heating  $Q_b$  (in unit of  $\text{MeV u}^{-1}$ ), where the helium mass fraction is set as  $Y_0 = 0.3$ ; the surface gravity  $g_{14} = 1.50$  is inferred from a  $1.5 M_\odot$  NS (Togashi EoS). The colours indicate the critical mass accretion rate at which marginally stable burning occurs for each scenario.

burning when  $Z > 0.05$  and  $Z > 0.03$ , respectively. Thus, the high metallicity is helpful to explain the observation of low critical mass accretion rate. For a fixed  $Z$ , the oscillation frequency and critical mass accretion rate decrease as  $Q_b$  increases. We show the details of the output quantities for the change of  $Z$  in Table A4 in Appendix.

### 3.4 Discussion

SS differs from NS mainly in two ways. First, the  $M-R$  relation for SS is very different from that of NS, especially, there is no minimum mass, and for  $M \lesssim 1 M_\odot$ ,  $M \propto R^3$ . In contrast, NSs have radii that decrease with increasing mass. Secondly, accretion of matter on to the SS surface will compress the matter at the base of the crust, inducing further electron captures on nuclei, followed by further neutron drip. Neutrons will be absorbed by the quark bag by the reactions  $n \rightarrow u + 2d$ ,  $d + u \rightarrow u + s$ , with heat release  $\sim 10-100$  MeV per one absorbed neutron (Farhi & Jaffe 1984; Page & Cumming 2005). However, accretion of matter on to the NS will compress the crust and induce the non-equilibrium processes including electron captures, neutron emission, and absorption and pycnonuclear fusion, with heat release  $\sim 1-2$  MeV per one accreted nucleon (Haensel & Zdunik 1990, 2003, 2008). Since the properties of X-ray burst are determined by the mass, radius, accretion composition, base flux, nuclear reaction rate etc (Meisel 2018; Johnston, Heger & Galloway 2020; Dohi et al. 2021, 2022), the X-ray burst from accreting NS should be different from accreting SS due to the above mainly differences. Furthermore, Xu (2003) conjectured astrophysical cold quark matter could be in a solid state, because the strong interaction may render quarks grouped in clusters. Thus, strange quark star could be strangeon stars, X-ray burst is likely to occur on the surface of accreting strangeon star.

It's useful to study the differences of mHz QPOs from accreting NS/SS. As the surface gravity of SS is greater than NS for the same mass, the oscillation frequency is higher for SS. On the other hand, as the heat flux at the base of the crust from accreting SS is larger than NS, it's possible that the accreted fuel is in stable burning when falling on to the surface of the stellar, and it is more likely occurs on an accreting SS with a relatively thick crust (the base of the crust is at the neutron drip density). For the simple one-zone model, when  $Q \gtrsim 2 \text{ MeV u}^{-1}$ , there will be no burst, the accreted fuel will be in stable burning. As far as we know, nearly half of all LMXBs show

type I X-ray burst (in't Zand et al. 2004), besides the LMXB which contains a black hole, SS star in LMXB is a good candidate that cannot show thermonuclear burning. However, the accreting strange (strangeon) star with a relatively thin crust (the base of the crust has not reached neutron drip density) should be excluded, because there is no neutron drip, as a result, no energy release from the transition of neutron to quark. In this case, the X-ray burst from accreting SS is similar to accreting NS.

### 4 CORRECTING FOR GENERAL RELATIVITY

To accurately model X-ray bursts, it is important to consider the effect of general relativity. Heger et al. (2007) reached the conclusion about the fact that general relativity corrections were compensated by the simplified network of the one-zone model because they compared results between KEPLER and onezone codes that they were using. The above results are calculated from the local Newton frame, it's necessary for us to correct the results for general relativity (Woosley et al. 2004). In general relativity, the gravitational acceleration at the surface of the star is  $g = GM_{\text{GR}}(1+z)/R_{\text{GR}}^2$ , where the gravitational redshift is  $z = (1 - 2GM_{\text{GR}}/(R_{\text{GR}}c^2))^{-1/2} - 1$ ,  $M_{\text{GR}}$  and  $R_{\text{GR}}$  are the general relativity mass and radius, respectively. Assume that the surface gravity in Newtonian calculations is equal to that in the general relativity

$$\frac{GM}{R^2} = \frac{GM_{\text{GR}}(1+z)}{R_{\text{GR}}^2} \quad (5)$$

We can apply our results to any combination of  $M_{\text{GR}}$  and  $R_{\text{GR}}$ . Here, we adopt the gravity in Table 1, by fixing the original radii (assume  $R = R_{\text{GR}}$ ), we can obtain  $M_{\text{GR}}$  that correspond to those values of  $g_s$  by equation (5). The details of the values for correcting for general relativity are shown in Table 2.

In the observer's frame, the oscillation period should be increased by a factor  $1+z$ ,

$$t_\infty = t(1+z) \quad (6)$$

As a result, the oscillation frequency will be divided by a factor  $1+z$ ,

$$f_\infty = \frac{f}{1+z} \quad (7)$$

We show  $f_\infty$  in Tables A1–A3 in the Appendix.

### 5 CONCLUSIONS

To study the effects of EOS on mHz QPOs, we employ a set of simulations using a simple one-zone X-ray burst model. We obtained  $g_s$  from a specified NS/SSs, where Togashi, MIT, and LX3630 EoSs are adopted in our simulations. The detailed studies show the differences of the properties of mHz QPOs due to  $g_s$ ,  $Q_b$ ,  $Y_0$ , and  $Z$ . The oscillation frequency from accreting SS is higher than that from accreting NS for the same mass, since the surface gravity of SS is greater than NS for the same mass. For a very low mass NS/SS, e.g.  $0.5 M_\odot$ , the oscillation frequency of accreting NS is less than 1 mHz when  $Y_0 = 0.3$ , while for SS, the oscillation frequency is  $\sim 2$  mHz. For a high mass NS and SS, e.g.  $1.5 M_\odot$ , the oscillation frequencies are  $\sim 3$  and  $\sim 4$  mHz, respectively. For  $2.0 M_\odot$  NS and SS, the oscillation frequencies are  $\sim 4$  and  $\sim 5$  mHz, respectively. The oscillation frequency and critical mass accretion rate decrease as base flux increases, the accreted fuel will be in stable burning when  $Q_b \gtrsim 2 \text{ MeV u}^{-1}$ , which indicates that the accreting NS/SS in LMXBs showing no type I X-ray bursts may have a strong base flux.

Meanwhile, we study the impact of accretion composition on mHz QPOs. The oscillation frequency increases as helium mass fraction increases. The oscillation frequency ranges from  $\sim 0.93\text{--}10.5$  mHz for NS,  $\sim 0.85\text{--}13.9$  mHz for SS if  $Y_0$  is in the range 0.3–0.5. The oscillation frequency ranges from  $\sim 1.5\text{--}37.9$  mHz for NS,  $\sim 3.3\text{--}53.8$  mHz for SS if  $Y_0$  is in the range 0.5–0.9. The critical mass accretion rate and oscillation amplitude increase as surface gravity and/or helium mass fraction increase, decrease as the base flux increases. Besides, the metallicity has non-negligible effect on the oscillation frequency and critical mass accretion rate. With increase of the metallicity, both the frequency and critical mass accretion rate decrease. Thus, the high metallicity and large base flux is helpful to reduce the discrepancy of critical mass accretion rate between models and observations. Since the above results are obtained from one-zone model with the local Newtonian gravity, we also correct the results for general relativity in Section 4.

So far, we have observed mHz QPOs from ten sources, for which the physical origin is considered as marginally stable burning. We have confirmed that the properties of these mHz QPOs are obviously affected by the  $M$ - $R$  relation and helium mass fraction. We also investigated the impacts of base heating and metallicity on the oscillation properties, and we found that the burst frequency and critical mass accretion rate will be decreased as the base heating and metallicity increase. We may infer the surface gravity if we know the oscillation frequency and mass accretion rate of mHz QPOs from Fig. 6. It is worth to mention that there are some high-energy mHz QPOs which are different from the low photon energy ( $\lesssim 5$  keV) mHz QPOs that are discussed in this paper, e.g. 35–95 mHz QPO is detected in the 27–120 keV energy band from the Be/X-ray binary 1A 0535+262 using Insight-HXMT data (Ma et al. 2022),  $\sim 41$  mHz QPO in the transient high-mass Be/X-ray binary pulsar 4U 0115+634 using data from the Rossi X-ray Timing Explorer (Dugair et al. 2013), the physical mechanism of these high-energy mHz QPOs is still needed to be revealed (Boroson et al. 2000; Bozzo et al. 2009).

To understand the X-ray burst phenomenon as well as marginally stable burning from accreting SS, one needs to focus the micro-physics inside SS and the nuclear reactions included rp processes in the envelope, and perform numerical calculations of the thermal evolution of SSs using a general relativistic stellar-evolution code (Dohi et al. 2020, 2021, 2022). Such a detailed study is left for our future work.

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## DATA AVAILABILITY

The data underlying this article are available in this article.

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## APPENDIX A

**Table A1.** Physical quantities of oscillation frequency  $f$ , critical mass accretion rate  $\dot{m}/\dot{m}_{\text{Edd}}$  (the last value of the critical mass accretion rate in its narrow range), maxima luminosity  $L_{36}(\text{max})$ , minima luminosity  $L_{36}(\text{min})$  with variation in surface gravity  $g_{14}$ , base heating  $Q_b$ , helium mass fraction  $Y_0$ , where the surface gravities are inferred from Togashi EoS,  $L_{36} = L/10^{36} \text{ erg s}^{-1}$ , the metallicity is set as  $Z = 0.01$ . The ‘-’ in the tables means that it is in stable burning when  $Q_b$  is high enough.  $f_\infty$  represents the general relativity correction for frequency.

$Y_0 = 0.3$ $Q_b(\text{MeV u}^{-1})$	$g_{14} = 0.50 (0.5 M_\odot)$					$g_{14} = 1.01 (1.0 M_\odot)$				
	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
0	0.75	0.70	0.402	4.12	2.16	1.71	1.50	0.667	6.97	3.52
0.15	0.70	0.65	0.380	4.53	1.79	1.63	1.43	0.641	6.59	3.44
0.5	0.68	0.64	0.341	3.45	1.87	1.57	1.38	0.577	5.87	3.16
1.0	0.55	0.51	0.273	2.69	1.55	1.34	1.18	0.482	4.71	2.74
1.5	0.39	0.36	0.202	1.92	1.18	0.93	0.82	0.382	3.67	2.20
2.0	-	-	-	-	-	0.73	0.64	0.274	2.69	1.57
$g_{14} = 1.50 (1.5 M_\odot)$										
$Q_b(\text{MeV/u})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
	0	2.70	2.23	0.879	8.85	4.89	3.85	2.96	1.125	11.17
0.15	2.53	2.09	0.85	8.76	4.66	3.80	2.92	1.081	10.80	5.37
0.5	2.41	1.99	0.760	7.98	4.14	3.36	2.58	0.975	9.90	4.81
1.0	1.95	1.61	0.642	6.68	3.50	2.85	2.19	0.820	8.56	4.00
1.5	1.56	1.29	0.516	5.20	2.91	2.53	1.95	0.67	6.32	3.58
2.0	1.19	0.98	0.385	3.85	2.22	1.80	1.38	0.508	4.68	2.77
$g_{14} = 2.12 (2.0 M_\odot)$										
$Q_b(\text{MeV/u})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
	0	1.27	1.19	0.517	5.30	2.76	2.91	2.55	0.842	9.07
0.15	1.31	1.22	0.496	4.96	2.73	2.76	2.42	0.812	8.28	4.38
0.5	1.19	1.11	0.443	4.44	2.43	2.62	2.30	0.732	7.47	3.96
1.0	0.99	0.93	0.365	3.61	2.04	2.10	1.84	0.613	6.26	3.33
1.5	0.74	0.69	0.281	2.75	1.60	1.69	1.48	0.491	4.92	2.75
2.0	-	-	-	-	-	1.38	1.21	0.368	3.55	2.14
$g_{14} = 1.50 (1.5 M_\odot)$										
$Q_b(\text{MeV/u})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
	0	4.55	3.76	1.109	11.66	5.94	6.56	5.05	1.41	13.86
0.15	4.17	3.45	1.065	11.23	5.70	6.46	4.97	1.355	13.52	6.73
0.5	3.89	3.21	0.961	10.51	5.07	5.75	4.42	1.225	12.27	6.13
1.0	3.49	2.88	0.815	8.39	4.55	4.83	3.72	1.025	11.13	4.83
1.5	2.80	2.31	0.651	7.03	3.48	4.03	3.10	0.830	8.94	4.01
2.0	2.03	1.68	0.500	4.84	2.92	3.06	2.35	0.647	6.03	3.51
$g_{14} = 2.12 (2.0 M_\odot)$										
$Q_b(\text{MeV/u})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
	0	4.55	3.76	1.109	11.66	5.94	6.56	5.05	1.41	13.86
0.15	4.17	3.45	1.065	11.23	5.70	6.46	4.97	1.355	13.52	6.73
0.5	3.89	3.21	0.961	10.51	5.07	5.75	4.42	1.225	12.27	6.13
1.0	3.49	2.88	0.815	8.39	4.55	4.83	3.72	1.025	11.13	4.83
1.5	2.80	2.31	0.651	7.03	3.48	4.03	3.10	0.830	8.94	4.01
2.0	2.03	1.68	0.500	4.84	2.92	3.06	2.35	0.647	6.03	3.51
$g_{14} = 0.50 (0.5 M_\odot)$										
$Q_b(\text{MeV u}^{-1})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
	0	1.93	1.80	0.639	6.57	3.42	4.35	3.82	1.033	10.82
0.15	1.87	1.75	0.612	6.39	3.24	4.19	3.68	0.994	10.21	5.33
0.5	1.70	1.59	0.550	5.58	3.01	3.86	3.39	0.891	9.36	4.76
1.0	1.46	1.36	0.456	4.69	2.46	3.19	2.80	0.745	8.49	3.71
1.5	1.12	1.05	0.360	3.54	2.03	2.70	2.37	0.605	6.35	3.25
2.0	-	-	-	-	-	1.96	1.72	0.460	4.45	2.64
$g_{14} = 1.50 (1.5 M_\odot)$										
$Q_b(\text{MeV/u})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
	0	7.09	5.86	1.353	14.74	7.05	10.48	8.06	1.720	17.24
0.15	6.56	5.42	1.302	13.93	6.81	9.36	7.20	1.651	16.78	8.06
0.5	6.27	5.18	1.175	12.77	6.16	8.82	6.78	1.498	15.01	7.42

**Table A1** – *continued*

$Y_0 = 0.3$	$g_{14} = 0.50 (0.5 M_\odot)$					$g_{14} = 1.01 (1.0 M_\odot)$				
$Q_b (\text{MeV u}^{-1})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
1.0	5.50	5.55	0.996	10.42	5.45	7.86	6.05	1.262	13.01	6.23
1.5	4.52	3.74	0.81	8.12	4.58	6.27	4.82	1.035	10.02	5.40
2.0	3.10	2.56	0.615	6.35	3.45	4.75	3.65	0.795	7.64	4.20
$Y_0 = 0.6$	$g_{14} = 0.50 (0.5 M_\odot)$					$g_{14} = 1.01 (1.0 M_\odot)$				
$Q_b (\text{MeV u}^{-1})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
0	2.73	2.55	0.770	8.25	3.96	6.80	5.96	1.245	12.80	6.67
0.15	2.69	2.51	0.74	7.81	3.89	6.15	5.39	1.196	12.50	6.33
0.5	2.43	2.27	0.665	6.99	3.50	5.79	5.08	1.081	11.28	5.75
1.0	2.21	2.07	0.555	5.80	2.94	4.82	4.23	0.915	9.49	4.85
1.5	1.81	1.69	0.440	4.65	2.33	3.70	3.25	0.740	7.32	4.16
2.0	–	–	–	–	–	2.93	2.57	0.560	5.69	3.11
$Y_0 = 0.6$	$g_{14} = 1.50 (1.5 M_\odot)$					$g_{14} = 2.12 (2.0 M_\odot)$				
$Q_b (\text{MeV u}^{-1})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
0	10.35	8.55	1.630	17.13	8.68	15.29	11.76	2.067	20.47	10.35
0.15	9.75	8.06	1.572	15.89	8.75	14.75	11.35	1.990	19.86	9.93
0.5	8.46	6.99	1.42	15.03	7.68	13.23	10.18	1.800	18.05	8.87
1.0	7.47	6.17	1.200	12.61	6.45	11.34	8.72	1.531	15.06	7.87
1.5	6.31	5.21	0.980	9.90	5.53	9.16	7.05	1.25	11.37	6.88
2.0	4.71	3.89	0.755	7.23	4.52	7.03	5.41	0.970	8.77	5.43
$Y_0 = 0.7$	$g_{14} = 0.50 (0.5 M_\odot)$					$g_{14} = 1.01 (1.0 M_\odot)$				
$Q_b (\text{MeV u}^{-1})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
0	4.02	3.76	0.920	10.41	4.44	8.82	7.74	1.486	15.66	7.76
0.15	3.70	3.46	0.885	9.17	4.68	8.50	7.46	1.428	15.02	7.54
0.5	3.61	3.37	0.800	8.12	4.37	7.65	6.71	1.295	13.17	7.01
1.0	3.03	2.83	0.671	6.49	3.83	6.94	6.09	1.095	11.09	6.01
1.5	2.36	2.01	0.535	5.60	2.88	5.75	5.04	0.89	8.94	4.94
2.0	–	–	–	–	–	4.14	3.63	0.68	6.68	3.85
$Y_0 = 0.7$	$g_{14} = 1.50 (1.5 M_\odot)$					$g_{14} = 2.12 (2.0 M_\odot)$				
$Q_b (\text{MeV u}^{-1})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
0	14.37	11.88	1.945	20.74	10.34	20.08	15.45	2.466	25.62	11.76
0.15	13.44	11.11	1.871	19.85	9.91	19.16	14.74	2.375	24.26	11.61
0.5	12.92	10.68	1.700	17.79	9.21	18.32	14.09	2.16	20.94	11.19
1.0	10.82	8.94	1.442	14.88	7.94	16.50	12.69	1.83	18.23	9.26
1.5	8.77	7.25	1.175	11.34	6.90	13.02	10.02	1.500	13.72	8.24
2.0	6.41	5.30	0.905	9.29	5.12	10.22	7.86	1.160	11.20	6.12
$Y_0 = 0.8$	$g_{14} = 0.50 (0.5 M_\odot)$					$g_{14} = 1.01 (1.0 M_\odot)$				
$Q_b (\text{MeV u}^{-1})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
0	5.52	5.16	1.095	11.25	5.80	12.82	11.25	1.768	18.10	9.44
0.15	4.98	4.65	1.051	11.91	5.08	11.19	9.82	1.700	17.67	9.02
0.5	4.55	4.25	0.950	10.15	4.98	10.62	9.32	1.540	16.21	8.06
1.0	4.40	4.11	0.802	8.27	4.34	9.42	8.26	1.305	12.76	7.45
1.5	3.26	2.86	0.650	6.41	3.68	7.75	6.80	1.065	10.26	6.16
2.0	–	–	–	–	–	5.81	5.10	0.82	8.00	4.65
$Y_0 = 0.8$	$g_{14} = 1.50 (1.5 M_\odot)$					$g_{14} = 2.12 (2.0 M_\odot)$				
$Q_b (\text{MeV u}^{-1})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
0	18.73	15.48	2.315	26.90	11.09	28.25	21.73	2.945	28.44	15.08
0.15	18.52	15.31	2.232	23.77	11.66	27.32	21.02	2.835	27.85	14.37
0.5	17.54	14.50	2.020	22.00	10.52	25.64	19.72	2.570	25.93	12.72
1.0	15.43	12.75	1.720	18.24	9.28	21.93	16.87	2.178	21.28	11.29
1.5	12.53	10.36	1.410	14.58	7.80	19.61	15.08	1.805	17.06	9.50
2.0	9.21	7.61	1.090	11.01	6.19	13.66	10.51	1.400	12.92	7.55
$Y_0 = 0.9$	$g_{14} = 0.50 (0.5 M_\odot)$					$g_{14} = 1.01 (1.0 M_\odot)$				
$Q_b (\text{MeV u}^{-1})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
0	7.58	7.08	1.300	13.57	6.90	16.18	14.19	2.105	21.70	11.33
0.15	6.69	6.25	1.250	12.99	6.62	16.13	14.15	2.025	21.33	10.60
0.5	6.41	5.99	1.130	11.66	6.04	14.01	12.29	1.835	19.60	9.49
1.0	5.89	5.50	0.960	9.82	5.25	13.44	11.79	1.565	16.32	8.40
1.5	4.47	4.18	0.780	7.87	4.28	10.96	9.61	1.280	13.19	7.00
2.0	–	–	–	–	–	8.09	7.10	0.99	9.57	5.65

**Table A1** – *continued*

$Y_0 = 0.3$ $Q_b(\text{MeV u}^{-1})$	$f(\text{mHz})$	$g_{14} = 0.50 (0.5 M_\odot)$			$L_{36}(\text{min})$	$f(\text{mHz})$	$g_{14} = 1.01 (1.0 M_\odot)$			$L_{36}(\text{min})$
		$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$			$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	
$g_{14} = 1.50 (1.5 M_\odot)$										
0	26.04	21.52	2.765	31.71	13.44	38.76	29.82	3.519	36.43	16.87
0.15	25.64	21.19	2.664	27.65	14.41	37.04	28.49	3.391	33.10	17.25
0.5	23.81	19.68	2.420	25.10	13.18	36.23	27.87	3.090	27.34	17.22
1.0	20.33	16.80	2.066	19.95	12.11	31.45	24.19	2.630	25.96	13.40
1.5	16.84	13.92	1.69	17.83	9.21	24.88	19.14	2.165	19.72	11.96
2.0	12.08	9.98	1.325	13.18	7.69	18.52	14.25	1.699	15.71	9.23

**Table A2.** The same as Table A1 but the surface gravities are inferred from MIT EoS.

$Y_0 = 0.3$ $Q_b(\text{MeV u}^{-1})$	$f(\text{mHz})$	$g_{14} = 1.15 (0.5 M_\odot)$			$L_{36}(\text{min})$	$f(\text{mHz})$	$g_{14} = 1.51 (1.0 M_\odot)$			$L_{36}(\text{min})$
		$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$			$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	
0	2.20	2.00	0.732	3.35	1.71	2.80	2.39	0.889	5.99	3.24
0.15	1.96	1.78	0.701	3.29	1.60	2.61	2.23	0.851	5.95	2.99
0.5	1.73	1.57	0.635	2.82	1.54	2.36	2.02	0.765	5.54	2.62
1.0	1.49	1.35	0.530	2.32	1.30	2.20	1.88	0.645	4.45	2.30
1.5	1.17	1.06	0.415	1.98	0.96	1.70	1.45	0.520	3.42	1.96
2.0	0.84	0.76	0.305	1.31	0.76	1.21	1.03	0.39	2.58	1.50
$g_{14} = 1.88 (1.5 M_\odot)$										
0	3.59	2.90	1.035	8.58	4.44	5.63	4.20	1.329	10.47	5.23
0.15	3.23	2.60	0.995	8.28	4.31	5.43	4.05	1.278	10.09	5.06
0.5	2.94	2.37	0.900	7.42	3.95	4.52	3.37	1.16	8.91	4.69
1.0	2.52	2.03	0.755	6.39	3.24	3.71	2.77	0.980	7.56	3.98
1.5	2.13	1.72	0.61	5.05	2.67	3.09	2.31	0.797	5.90	3.33
2.0	1.51	1.22	0.460	3.71	2.08	2.36	1.76	0.605	4.00	2.84
$g_{14} = 1.15 (0.5 M_\odot)$										
$Y_0 = 0.4$ $Q_b(\text{MeV u}^{-1})$	$f(\text{mHz})$	$g_{14} = 1.15 (0.5 M_\odot)$			$L_{36}(\text{min})$	$f(\text{mHz})$	$g_{14} = 1.51 (1.0 M_\odot)$			$L_{36}(\text{min})$
		$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$			$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	
0	3.46	3.15	0.923	4.24	2.15	4.49	3.84	1.112	8.06	3.77
0.15	3.31	3.01	0.888	4.02	2.09	4.23	3.62	1.070	7.56	3.77
0.5	2.83	2.57	0.798	3.73	1.84	3.82	3.26	0.960	7.73	2.99
1.0	2.59	2.35	0.673	2.99	1.61	3.35	2.86	0.810	6.14	2.70
1.5	2.01	1.83	0.540	2.43	1.31	2.65	2.26	0.66	4.50	2.42
2.0	1.53	1.39	0.405	1.77	1.01	2.03	1.74	0.5	3.36	1.88
$g_{14} = 1.88 (1.5 M_\odot)$										
0	5.93	4.78	1.295	11.03	5.42	8.77	6.54	1.660	13.27	6.45
0.15	5.69	4.59	1.247	10.29	5.40	8.38	6.25	1.598	12.65	6.25
0.5	5.18	4.18	1.13	9.36	4.93	7.41	5.53	1.445	11.80	5.62
1.0	4.41	3.56	0.955	7.81	4.18	6.41	4.78	1.225	9.70	4.89
1.5	3.60	2.90	0.770	5.74	3.67	5.38	4.01	1.000	7.28	4.31
2.0	2.65	2.14	0.590	4.74	2.68	3.98	2.97	0.766	5.87	3.19
$g_{14} = 1.15 (0.5 M_\odot)$										
$Y_0 = 0.5$ $Q_b(\text{MeV u}^{-1})$	$f(\text{mHz})$	$g_{14} = 1.15 (0.5 M_\odot)$			$L_{36}(\text{min})$	$f(\text{mHz})$	$g_{14} = 1.51 (1.0 M_\odot)$			$L_{36}(\text{min})$
		$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$			$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	
0	4.98	4.53	1.131	5.11	2.66	7.34	6.27	1.361	9.40	4.84
0.15	4.99	4.54	1.086	5.02	2.54	7.22	6.17	1.310	9.12	4.62
0.5	4.33	3.94	0.980	4.57	2.28	6.17	5.27	1.185	8.19	4.26
1.0	3.85	3.5	0.825	3.76	1.94	5.18	4.43	1.000	7.12	3.56
1.5	3.26	2.96	0.670	2.96	1.66	4.36	3.73	0.815	5.43	3.03
2.0	2.24	2.04	0.510	2.14	1.30	3.19	2.73	0.620	4.13	2.35
$g_{14} = 1.88 (1.5 M_\odot)$										
0	8.91	7.19	1.585	13.04	6.87	12.92	9.64	2.025	16.24	7.83
0.15	8.77	7.07	1.523	12.59	6.58	12.17	9.08	1.945	16.01	7.36
0.5	7.82	6.31	1.370	11.26	5.98	11.82	8.82	1.765	13.09	7.39

**Table A2 – continued**

$Y_0 = 0.3$		$g_{14} = 1.15 (0.5 M_\odot)$				$g_{14} = 1.51 (1.0 M_\odot)$				
$Q_b (\text{MeV u}^{-1})$	$f (\text{mHz})$	$f_\infty (\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f (\text{mHz})$	$f_\infty (\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
1.0	6.51	5.25	1.150	10.77	4.52	10.04	7.49	1.501	11.68	6.10
1.5	5.65	4.56	0.950	7.83	4.22	8.21	6.13	1.230	9.28	5.15
2.0	4.11	3.31	0.730	5.68	3.35	6.20	4.63	0.950	6.90	4.08
$Y_0 = 0.6$		$g_{14} = 1.15 (0.5 M_\odot)$				$g_{14} = 1.51 (1.0 M_\odot)$				
$Q_b (\text{MeV u}^{-1})$	$f (\text{mHz})$	$f_\infty (\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f (\text{mHz})$	$f_\infty (\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
0	7.03	6.39	1.360	6.17	3.20	9.47	8.09	1.635	11.69	5.61
0.15	7.41	6.74	1.308	5.96	3.03	9.36	8.00	1.575	10.99	5.53
0.5	6.56	5.96	1.18	4.92	3.01	8.91	7.62	1.420	9.56	5.17
1.0	5.63	5.12	1.000	4.50	2.40	7.22	6.17	1.210	8.16	4.44
1.5	4.47	4.06	0.810	3.67	1.93	6.29	5.38	0.980	6.65	3.60
2.0	3.39	3.08	0.620	2.63	1.56	4.68	4.00	0.760	4.95	2.93
$Y_0 = 0.7$		$g_{14} = 1.88 (1.5 M_\odot)$				$g_{14} = 2.69 (2.0 M_\odot)$				
$Q_b (\text{MeV u}^{-1})$	$f (\text{mHz})$	$f_\infty (\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f (\text{mHz})$	$f_\infty (\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
0	12.72	10.26	1.900	16.62	7.72	19.38	14.46	2.439	18.89	9.81
0.15	12.17	9.81	1.827	15.91	7.46	18.73	13.98	2.351	17.85	9.70
0.5	11.90	9.60	1.660	13.77	7.22	17.54	13.09	2.130	16.40	8.59
1.0	10.16	8.19	1.405	11.75	6.04	14.37	10.72	1.811	14.06	7.33
1.5	7.75	6.25	1.125	10.58	4.50	11.99	8.95	1.48	10.62	6.45
2.0	6.13	4.94	0.89	6.94	4.11	9.16	6.84	1.15	8.20	5.10
$Y_0 = 0.8$		$g_{14} = 1.15 (0.5 M_\odot)$				$g_{14} = 1.51 (1.0 M_\odot)$				
$Q_b (\text{MeV u}^{-1})$	$f (\text{mHz})$	$f_\infty (\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f (\text{mHz})$	$f_\infty (\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
0	9.92	9.02	1.621	7.63	3.68	14.62	12.50	1.955	13.38	6.98
0.15	9.98	9.07	1.561	7.11	3.63	13.44	11.49	1.880	13.10	6.63
0.5	9.11	8.28	1.41	6.72	3.18	12.25	10.47	1.700	12.59	5.76
1.0	8.25	7.50	1.195	5.41	2.83	9.69	8.28	1.420	11.87	4.33
1.5	6.06	5.51	0.975	4.35	2.37	8.17	6.98	1.160	9.02	3.87
2.0	4.85	4.41	0.750	3.16	1.90	6.61	5.65	0.91	5.97	3.43
$Y_0 = 0.9$		$g_{14} = 1.88 (1.5 M_\odot)$				$g_{14} = 2.69 (2.0 M_\odot)$				
$Q_b (\text{MeV u}^{-1})$	$f (\text{mHz})$	$f_\infty (\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f (\text{mHz})$	$f_\infty (\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
0	24.88	20.06	2.705	24.65	10.63	37.88	28.27	3.475	27.21	13.86
0.15	23.47	18.93	2.608	21.89	11.23	37.88	28.27	3.345	25.75	13.43
0.5	22.22	17.92	2.360	18.83	10.56	32.68	24.39	3.040	23.96	12.14
1.0	18.94	15.27	2.000	16.65	8.70	29.76	22.21	2.59	20.14	10.43
1.5	14.49	11.69	1.62	15.18	6.46	23.15	17.28	2.11	17.24	8.22
2.0	11.99	9.67	1.290	9.84	6.03	18.52	13.82	1.670	11.74	7.49

**Table A2** – *continued*

$Y_0 = 0.3$ $Q_b(\text{MeV u}^{-1})$	$f(\text{mHz})$	$g_{14} = 1.15 (0.5 M_\odot)$			$L_{36}(\text{min})$	$f(\text{mHz})$	$g_{14} = 1.51 (1.0 M_\odot)$			$L_{36}(\text{min})$
		$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$			$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	
$g_{14} = 1.88 (1.5 M_\odot)$										
0	32.05	25.85	3.235	28.64	12.93	49.02	36.58	4.155	34.11	15.72
0.15	32.05	25.85	3.119	26.07	13.19	47.62	35.54	4.01	30.52	16.41
0.5	30.86	24.89	2.837	23.76	12.18	45.05	33.62	3.652	28.58	14.38
1.0	26.46	21.34	2.410	19.29	10.47	39.68	29.61	3.100	23.16	13.06
1.5	20.58	15.36	1.920	20.03	7.00	30.86	23.03	2.500	22.15	9.08
2.0	15.87	12.80	1.485	15.40	5.78	25.25	18.84	2.010	14.00	9.02

**Table A3.** The same as Table A1 but the surface gravities are inferred from LX3630 EoS. As the surface gravity of  $1.5M_\odot$  is the same as that from MIT EoS, we leave it out here.

$Y_0 = 0.3$ $Q_b(\text{MeV u}^{-1})$	$f(\text{mHz})$	$g_{14} = 1.22 (0.5 M_\odot)$			$L_{36}(\text{min})$	$f(\text{mHz})$	$g_{14} = 1.59 (1.0 M_\odot)$			$L_{36}(\text{min})$
		$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$			$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	
0	2.32	2.09	0.765	3.19	1.72	2.79	2.38	0.915	7.29	2.56
0.15	2.12	1.91	0.730	3.21	1.58	2.71	2.32	0.885	5.86	3.01
0.5	1.96	1.77	0.660	2.84	1.47	2.47	2.11	0.800	5.12	2.81
1.0	1.59	1.43	0.551	2.38	1.21	2.16	1.85	0.670	4.40	2.30
1.5	1.30	1.17	0.440	1.83	1.02	1.79	1.53	0.540	3.41	1.92
2.0	0.85	0.77	0.32	1.37	0.73	1.28	1.09	0.405	2.50	1.49
$g_{14} = 2.16 (2.0 M_\odot)$										
0	4.07	3.13	1.140	10.97	5.70	5.93	4.12	1.390	15.00	7.92
0.15	3.76	2.89	1.090	11.91	4.88	5.36	3.72	1.330	13.78	7.93
0.5	3.66	2.82	0.987	10.05	4.77	4.76	3.31	1.210	13.13	6.92
1.0	3.10	2.38	0.838	7.53	5.52	4.59	3.19	1.025	10.99	5.90
1.5	2.57	1.98	0.680	6.24	3.59	3.45	2.40	0.830	8.49	5.06
2.0	1.82	1.40	0.505	5.05	2.55	2.60	1.81	0.630	6.85	3.69
$g_{14} = 1.22 (0.5 M_\odot)$										
$Y_0 = 0.4$ $Q_b(\text{MeV u}^{-1})$	$f(\text{mHz})$	$g_{14} = 1.22 (0.5 M_\odot)$			$L_{36}(\text{min})$	$f(\text{mHz})$	$g_{14} = 1.59 (1.0 M_\odot)$			$L_{36}(\text{min})$
		$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$			$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	
0	3.65	3.29	0.960	4.25	2.04	4.89	4.18	1.155	7.60	3.91
0.15	3.20	2.88	0.920	4.27	1.90	4.72	4.03	1.11	7.46	3.74
0.5	3.07	2.77	0.830	3.76	1.74	4.42	3.78	1.000	6.92	3.25
1.0	2.66	2.40	0.695	3.16	1.46	3.62	3.09	0.845	5.64	2.87
1.5	2.10	1.89	0.545	2.81	1.07	2.87	2.45	0.680	4.56	2.32
2.0	1.54	1.39	0.420	1.81	0.95	2.16	1.85	0.520	2.91	2.07
$g_{14} = 2.16 (2.0 M_\odot)$										
0	7.09	5.45	1.430	13.53	7.29	9.16	6.36	1.735	19.23	9.65
0.15	6.17	4.75	1.370	13.81	6.64	9.01	6.26	1.673	18.24	9.47
0.5	5.95	4.58	1.240	11.48	6.45	7.79	5.41	1.51	17.00	8.38
1.0	5.22	4.02	1.050	10.32	5.22	6.97	4.84	1.285	13.80	7.44
1.5	3.95	3.04	0.830	9.56	3.72	5.65	3.92	1.035	12.10	5.71
2.0	3.11	2.39	0.655	5.99	3.47	4.11	2.85	0.800	8.68	4.68
$g_{14} = 1.22 (0.5 M_\odot)$										
$Y_0 = 0.5$ $Q_b(\text{MeV u}^{-1})$	$f(\text{mHz})$	$g_{14} = 1.22 (0.5 M_\odot)$			$L_{36}(\text{min})$	$f(\text{mHz})$	$g_{14} = 1.59 (1.0 M_\odot)$			$L_{36}(\text{min})$
		$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$			$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	
0	5.41	4.87	1.175	5.09	2.56	7.72	6.60	1.410	9.47	4.66
0.15	5.34	4.81	1.13	4.90	2.46	7.03	6.01	1.355	9.01	4.56
0.5	4.89	4.41	1.020	4.48	2.22	6.49	5.55	1.22	7.90	4.22
1.0	4.02	3.62	0.850	4.04	1.74	6.04	5.16	1.035	6.97	3.49
1.5	3.27	2.95	0.675	3.50	1.33	4.68	4.00	0.825	6.10	2.59
2.0	2.46	2.22	0.530	2.14	1.25	3.37	2.88	0.645	4.01	2.33
$g_{14} = 2.16 (2.0 M_\odot)$										
0	10.96	8.43	1.740	17.20	8.48	13.89	9.65	2.110	25.95	10.71
0.15	9.75	7.50	1.675	16.22	8.35	13.55	9.41	2.025	21.98	11.59

**Table A3** – *continued*

$Y_0 = 0.3$ $Q_b(\text{MeV u}^{-1})$	$g_{14} = 1.22 (0.5 M_\odot)$					$g_{14} = 1.59 (1.0 M_\odot)$				
	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
0.5	9.16	7.38	1.51	14.28	7.76	12.53	8.70	1.840	20.44	10.26
1.0	7.65	5.88	1.285	12.51	6.49	10.96	7.61	1.558	17.76	8.58
1.5	6.29	4.84	1.030	11.09	4.86	8.82	6.13	1.28	12.99	7.85
2.0	4.80	3.69	0.810	7.36	4.33	6.59	4.58	0.990	10.24	5.93
$Y_0 = 0.6$	$g_{14} = 1.22 (0.5 M_\odot)$					$g_{14} = 1.59 (1.0 M_\odot)$				
$Q_b(\text{MeV u}^{-1})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
0	7.65	6.89	1.415	6.09	3.08	10.35	8.85	1.695	11.45	5.58
0.15	7.34	6.61	1.363	5.83	2.98	9.98	8.53	1.632	10.94	5.47
0.5	6.86	6.18	1.230	5.35	2.67	9.80	8.38	1.480	9.94	5.01
1.0	5.93	5.34	1.040	4.50	2.31	7.65	6.54	1.255	7.98	4.39
1.5	4.90	4.41	0.840	3.65	1.86	6.56	5.61	1.020	6.47	3.61
2.0	3.65	3.29	0.650	2.54	1.57	4.93	4.21	0.790	4.48	3.13
$Y_0 = 0.7$	$g_{14} = 2.16 (2.0 M_\odot)$					$g_{14} = 2.86 (3.0 M_\odot)$				
$Q_b(\text{MeV u}^{-1})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
0	15.15	11.65	2.090	22.67	9.28	19.61	13.62	2.540	28.92	13.72
0.15	13.77	10.59	2.015	19.52	10.04	19.16	13.31	2.450	26.52	13.92
0.5	13.89	10.68	1.830	17.76	9.23	19.61	13.62	2.220	24.69	12.41
1.0	11.74	9.03	1.543	15.72	7.55	15.72	10.92	1.891	19.03	11.59
1.5	9.36	7.2	1.265	11.57	6.77	12.92	8.97	1.545	15.75	9.45
2.0	7.61	5.85	0.985	8.95	5.27	9.86	6.85	1.201	12.22	7.44
$Y_0 = 0.8$	$g_{14} = 1.22 (0.5 M_\odot)$					$g_{14} = 1.59 (1.0 M_\odot)$				
$Q_b(\text{MeV u}^{-1})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
0	11.11	10.01	1.688	7.66	3.47	15.02	12.84	2.026	13.42	6.86
0.15	10.48	9.44	1.626	6.92	3.60	14.75	12.61	1.950	12.71	6.60
0.5	10.16	9.15	1.475	6.33	3.26	13.33	11.39	1.76	11.16	6.17
1.0	8.46	7.62	1.240	5.61	2.65	11.49	9.82	1.5	9.74	5.20
1.5	7.03	6.33	1.010	4.48	2.20	9.69	8.28	1.225	7.90	4.28
2.0	5.18	4.67	0.780	3.15	1.83	7.12	6.09	0.940	6.18	3.29
$Y_0 = 0.8$	$g_{14} = 2.16 (2.0 M_\odot)$					$g_{14} = 2.86 (3.0 M_\odot)$				
$Q_b(\text{MeV u}^{-1})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
0	21.10	16.23	2.501	24.13	12.46	29.24	20.31	3.034	31.73	17.92
0.15	19.84	15.26	2.410	23.17	12.14	27.32	18.97	2.925	32.37	16.43
0.5	18.12	13.94	2.180	21.80	10.63	26.04	18.08	2.660	28.82	15.26
1.0	16.18	12.45	1.851	18.22	9.17	22.52	15.64	2.26	23.81	13.49
1.5	13.33	10.25	1.525	14.76	7.76	18.73	13.01	1.861	19.77	10.92
2.0	9.98	7.68	1.18	10.71	6.29	14.12	9.81	1.450	14.44	9.12
$Y_0 = 0.8$	$g_{14} = 1.22 (0.5 M_\odot)$					$g_{14} = 1.59 (1.0 M_\odot)$				
$Q_b(\text{MeV u}^{-1})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
0	14.88	13.41	2.011	8.69	4.38	21.65	18.50	2.413	15.77	8.14
0.15	14.62	13.17	1.938	8.15	4.32	20.33	17.38	2.322	15.52	7.77
0.5	13.66	12.31	1.755	7.56	3.85	18.94	16.19	2.110	14.03	7.14
1.0	11.66	10.50	1.490	6.40	3.29	16.03	13.70	1.795	11.59	6.25
1.5	9.58	8.63	1.220	4.97	2.83	12.72	10.87	1.470	9.50	5.12
2.0	7.25	6.53	0.940	3.80	2.20	9.98	8.53	1.135	7.09	4.13
$Y_0 = 0.8$	$g_{14} = 2.16 (2.0 M_\odot)$					$g_{14} = 2.86 (3.0 M_\odot)$				
$Q_b(\text{MeV u}^{-1})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$f_\infty(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
0	31.45	24.19	2.980	26.79	15.88	40.65	28.23	3.624	40.27	20.13
0.15	26.88	20.68	2.875	27.70	14.48	39.68	27.56	3.495	38.18	19.69
0.5	24.15	18.58	2.593	30.67	10.78	35.46	24.63	3.160	38.88	16.02
1.0	21.93	16.87	2.215	21.02	11.30	29.76	20.67	2.710	29.32	15.62
1.5	18.12	13.94	1.82	18.07	8.94	26.46	18.38	2.220	24.66	12.63
2.0	14.01	10.78	1.420	12.87	7.57	19.38	13.46	1.74	16.23	11.56

**Table A3** – *continued*

$Y_0 = 0.3$	$g_{14} = 1.22 (0.5 M_\odot)$				$g_{14} = 1.59 (1.0 M_\odot)$					
$Q_b (\text{MeV u}^{-1})$	$f(\text{mHz})$	$f_\infty (\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$f_\infty (\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
$Y_0 = 0.9$										
		$g_{14} = 1.22 (0.5 M_\odot)$				$g_{14} = 1.59 (1.0 M_\odot)$				
$Q_b (\text{MeV u}^{-1})$	$f(\text{mHz})$	$f_\infty (\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$f_\infty (\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
0	19.84	17.87	2.395	11.16	4.84	28.25	24.15	2.878	19.73	9.43
0.15	19.61	17.67	2.310	9.69	5.15	27.78	23.74	2.776	18.53	9.45
0.5	18.52	16.68	2.090	8.42	4.89	26.04	22.26	2.52	16.54	8.59
1.0	16.03	14.44	1.788	7.04	4.29	22.22	18.99	2.15	14.27	7.41
1.5	13.44	12.11	1.466	6.45	3.28	17.73	15.15	1.760	11.99	5.76
2.0	10.22	9.21	1.140	4.60	2.70	13.77	11.77	1.370	9.03	4.85
		$g_{14} = 2.16 (2.0 M_\odot)$				$g_{14} = 2.86 (3.0 M_\odot)$				
$Q_b (\text{MeV u}^{-1})$	$f(\text{mHz})$	$f_\infty (\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$f_\infty (\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
0	40.65	31.27	3.565	38.97	15.70	53.76	37.33	4.339	60.29	19.29
0.15	39.68	30.52	3.438	32.95	17.31	50.51	35.08	4.184	50.75	21.48
0.5	36.23	27.87	3.12	30.84	15.40	52.08	36.17	3.812	41.30	21.93
1.0	29.76	22.89	2.670	25.39	13.73	42.74	29.68	3.260	32.52	20.23
1.5	24.15	18.58	2.160	22.86	10.17	35.46	24.63	2.680	27.24	16.48
2.0	19.84	15.26	1.720	15.14	9.58	26.88	18.67	2.100	20.65	13.25

**Table A4.** Physical quantities of oscillation frequency  $f$ , critical mass accretion rate  $\dot{m}/\dot{m}_{\text{Edd}}$  (the last value of the critical mass accretion rate in its narrow range), maxima luminosity  $L_{36}(\text{max})$ , and minima luminosity  $L_{36}(\text{min})$  with variation in metallicity  $Z$ . The gravity is set as  $g_{14} = 1.50$ , which corresponds to a  $1.5 M_\odot$  NS from Togashi EoS.

$Y_0 = 0.3$	$Q_b = 0 (\text{MeV u}^{-1})$			$Q_b = 0.15 (\text{MeV u}^{-1})$				
$Z$	$f(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
0.01	2.70	0.879	8.85	4.89	2.53	0.850	8.76	4.66
0.03	2.64	0.850	8.72	4.66	2.71	0.812	8.82	4.29
0.05	2.59	0.812	8.52	4.37	2.56	0.777	7.92	4.32
0.07	2.51	0.773	8.25	4.13	2.29	0.731	7.79	3.89
0.1	2.32	0.703	7.33	3.87	2.17	0.653	6.95	3.50
		$Q_b = 0.5 (\text{MeV u}^{-1})$			$Q_b = 1.0 (\text{MeV u}^{-1})$			
$Q_b (\text{MeV u}^{-1})$	$f(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
0.01	2.41	0.760	7.98	4.14	1.95	0.642	6.68	3.50
0.03	2.30	0.726	7.55	3.97	1.89	0.595	5.99	3.40
0.05	2.17	0.680	7.25	3.65	1.77	0.535	5.38	3.02
0.07	2.07	0.625	6.73	3.33	1.34	0.450	4.38	2.63
0.1	1.70	0.515	5.32	2.85	–	–	–	–
		$Q_b = 1.5 (\text{MeV u}^{-1})$			$Q_b = 2.0 (\text{MeV u}^{-1})$			
$Q_b (\text{MeV u}^{-1})$	$f(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$	$f(\text{mHz})$	$\dot{m}/\dot{m}_{\text{Edd}}$	$L_{36}(\text{max})$	$L_{36}(\text{min})$
0.01	1.56	0.516	5.20	2.91	1.19	0.385	3.85	2.22
0.03	1.43	0.455	4.10	2.84	0.77	0.265	2.52	1.61
0.05	1.03	0.335	3.25	1.98	–	–	–	–
0.07	–	–	–	–	–	–	–	–
0.1	–	–	–	–	–	–	–	–

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